

1 An in-silico study to determine whether changes to mitochondria  
2 organization through engineered mitochondrial dynamics can  
3 enhance bioenergetics in cardiomyocytes

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11

12 **SUMMARY**

13 Mitochondria are the powerhouse of the cell and owing to their unique energetic demands,  
14 heart muscles contain a high density of mitochondria. In conditions such as heart failure and  
15 diabetes-induced heart disease, changes in the organization of cardiac mitochondria are  
16 common. While recent studies have also shown that cardiac mitochondria split and fuse  
17 throughout the cell, a mechanistic understanding of how mitochondrial dynamics may affect  
18 energy output is lacking. Using a mathematical model that has been fitted to experimental  
19 data, we test if briefly altering fission or fusion rates improves ATP production and supply in  
20 cardiomyocytes. Unexpectedly, we found that cardiac bioenergetics, e.g., the ADP/ATP ratio,  
21 were robust to changes in fusion and fission rates and consequently mitochondria  
22 organization. Our study highlights complex nonlinear feedback loops that are at play in the  
23 cross-talk between mitochondrial dynamics and bioenergetics. The study motivate further in-  
24 silico and experimental investigations to determine the mechanistic basis for new therapies  
25 that target mitochondrial dynamics.

26 **INTRODUCTION**

27 Mitochondria meet cellular energy demands by converting nutrients into chemical energy in the  
28 form of adenosine triphosphate (ATP). In cardiomyocytes, mitochondria have a particularly  
29 challenging job, as the heart must pump up to six liters of blood per minute, continuously from  
30 birth (Cattermole et al., 2017). Far from being static, mitochondria are continually splitting and  
31 fusing in response to energy demands and stressors (Youle and van der Bliek, 2012). While  
32 there is an emerging body of experimental work on fission and fusion, there are no quantitative  
33 frameworks that investigate if changes to mitochondria arrangement through fission and fusion,  
34 can help with bioenergetics.

35

36 Cardiac mitochondria are organized into networks of clusters that are interspersed between  
37 contractile protein bundles (Ghosh et al., 2018; Glancy et al., 2017). Furthermore, recent  
38 studies have shown that cardiac mitochondria dynamically change their organization (Ong et  
39 al., 2015; Glancy et al., 2017) and that it is disrupted in diseased cells (Cao and Zheng, 2019;  
40 Galloway and Yoon, 2015). For example, Chen et al. (2012) observed that deficiencies in  
41 mitochondrial fusion proteins Mfn1 and Mfn2 cause mild cardiomyopathy. Additionally, defects  
42 in the mitochondrial fusion protein OPA1 may cause left ventricular hypertrophy (Piquereau et  
43 al., 2012), leading to an increased risk of arrhythmias and heart failure (Frey et al., 2004; Frey  
44 and Olson, 2003). Studies by Jarosz et al. (2017) and others have also shown altered  
45 organization of mitochondria in diabetes-induced heart disease.

46

47 A key idea is that mitochondrial fission and fusion ensure mitochondrial quality control. For  
48 example, mitochondrial fusion mediates internal machinery sharing, such as mitochondrial  
49 respiration and the equilibration of mitochondrial membrane potential (Chen et al., 2005; Eisner  
50 et al., 2014; Eisner et al., 2018; Twig et al., 2008). Thus, by diluting mitochondrial damage,  
51 fusion facilitates quality control. Drp1 governs mitochondrial fission (Eisner et al., 2018; Ong et  
52 al., 2015), which facilitates quality control by promoting the fragmentation of highly damaged  
53 mitochondria and limits the propagation of mitochondrial dysfunction (Eisner et al., 2018;  
54 Glancy et al., 2017; Ong et al., 2015). Fission decreases during periods of bioenergetic stress

55 (Gomes et al., 2011; Liesa and Shirihai, 2013), but may increase locally to minimize the spread  
56 of mitochondrial dysfunction.

57

58 Additionally, mitochondria may increase their mass via biogenesis in response to nutrient  
59 deprivation or biophysical stress (Jager et al., 2007; Mihaylova and Shaw, 2011; Scarpulla,  
60 2011). Mitochondrial biogenesis is mediated by the protein PGC-1 $\alpha$  and facilitates quality  
61 control by promoting mitochondrial homeostasis (Boland et al., 2013; Dalmasso et al., 2017).  
62 Despite tantalizing evidence that mitochondrial dynamics could be a possible therapeutic target  
63 to improve mitochondrial function and hence energy production (Ong et al., 2015; Ong et al.,  
64 2010), the precise mechanistic and causal relationship between mitochondrial dynamics and  
65 bioenergetics is still to be explored and would help identify effective drug targets that are based  
66 on the underlying mechanism.

67

68 We have recently shown that changes to mitochondria morphology and their spatial distribution  
69 can influence the distribution of energy metabolites and consequently contractile force under  
70 hypoxic, and high-workload conditions (Jarosz et al., 2017; Ghosh et al., 2018). These findings  
71 were based on a biophysics-based computational model of mitochondrial function and spatially  
72 detailed geometric models of cardiomyocyte architecture derived from electron microscopy  
73 images. The works of Eisner et al. (2017) and Glancy et al. (2017) suggest that mitochondrial  
74 fusion and fission rates increase at high workloads. These experimental and computational  
75 studies suggest that dynamic alterations to mitochondria density and organization, via altered  
76 fusion and fission rates, might help meet high energy demands at high workloads.

77

78 In this study, we integrated the current understanding, that is outlined above, of the role of  
79 fusion/fission dynamics, biogenesis, and mitochondria organization on cardiac bioenergetics  
80 into a semi-quantitative mechanistic modelling framework. We have considered three plausible  
81 mechanisms by which mitochondrial dynamics and biogenesis can regulate mitochondrial  
82 bioenergetics: (i) increased mitochondrial connectivity can enhance the OXPHOS capacity of  
83 individual mitochondria; (ii) increased mitochondrial volume can increase the total OXPHOS  
84 capacity of the cell; and (iii) mitochondrial network reorganization can favorably alter the

85 diffusion distances between mitochondria and myofibrils for a rapid and steady supply of ATP.  
86 We used this model to test whether increasing or decreasing fusion or fission rates would affect  
87 bioenergetics. Specifically, we sought to: (i) investigate whether mitochondria network  
88 morphological changes stemming from alterations in network connectivity affect bioenergetics  
89 in physiological and pathological conditions; (ii) investigate how altered bioenergetics could  
90 affect fusion/fission dynamics; and (iii) determine key parameters that need to be measured to  
91 robustly validate or negate our model predictions and thus formalize a mechanistic model of  
92 the link between mitochondrial dynamics and bioenergetics. As a simplifying assumption, we  
93 limit our study to investigate cross-talk on an acute scale of two minutes, which suffices to  
94 observe dynamic changes in mitochondria network morphology experimentally (Glancy et al.,  
95 2017).

96

97 Our computational model is a hybrid agent-based- and partial differential equation model. The  
98 agent-based model (ABM) simulates changes in mitochondrial connectivity and mitochondrial  
99 mass such as fission, fusion, and biogenesis, while the partial differential equation (PDE)  
100 system models various bioenergetic interactions such as oxidative-phosphorylation  
101 (OXPHOS), ATP hydrolysis, and the breakdown of reactive oxygen species. In particular, we  
102 assumed that mitochondrial connectivity, via fission and fusion, directly feeds forward into  
103 OXPHOS and electron transport chain (ETC) kinetics, which then feeds back into fission and  
104 fusion dynamics. We then calibrated the model against existing experimental data on cardiac  
105 mitochondrial bioenergetics and dynamics in the literature.

106

107 Remarkably, our simulations show that bioenergetics are robust to varied fission and fusion  
108 rates in the short term under physiological conditions. However, fusion and fission may  
109 enhance bioenergetics when mitochondrial function is compromised. Since these findings  
110 largely depend on ETC enzyme kinetic rates, they highlight the need for experimental  
111 measurements of how ETC enzyme kinetics change during mitochondrial fission and fusion.  
112 Moreover, we predict that high workloads may increase mitochondrial volume fractions, which  
113 may enhance energetics to meet these high workload demands. Indeed, this study reveals a

114 challenging inverse problem, if the ADP/ATP ratio is robust to changes in fission and fusion,  
115 when and under what circumstances do fission and fusion impact bioenergetics?

116

117 **RESULTS**

118 **A hybrid agent-based model of mitochondrial dynamics and continuum reaction-  
119 diffusion model of bioenergetics**

120 Details of the mathematical model equations that were coupled to create our computational  
121 model of mitochondrial dynamics and bioenergetics are provided in the [Methods](#) section. [Figure](#)  
122 [1A](#) illustrates the initial geometry used by the model and is inspired by longitudinal views of  
123 cardiac cell architecture under the electron microscope. [Figure 1B](#) outlines the basic factors  
124 that change during interactions. In brief, mitochondrial dynamics and biogenesis are modulated  
125 by energetic stress – a catch-all term that encompasses the ratio of ADP-to-ATP, mitochondrial  
126 connectivity, and mitochondrial damage. Consequently, changes in energetic stress can alter  
127 the mitochondrial network architecture which further leads to alterations in the mitochondrial  
128 ATP synthesis rate and the resulting ADP-to-ATP ratio across the cell. The mitochondrial ADP-  
129 to-ATP ratio then governs the energetic stress, establishing a feedback loop between  
130 mitochondrial dynamics and bioenergetics. Further details are provided in the [Methods](#) section,  
131 specifically in the subsection titled “Agent based model” .

132

133 [Figure 2](#) demonstrates that the model captures the key bioenergetic and mitochondrial  
134 dynamics features found in current experimental data in the literature. In particular, the model  
135 reproduces intracellular ADP (see [Figure 2A](#)) and O<sub>2</sub> distributions (see [Figure 2B](#)) similar to  
136 those reported by [Vendelin et al. \(2000\)](#) and [Takahashi et al. \(1998\)](#). To simulate the different  
137 workloads depicted in [Figure 2A](#), we individually vary  $X_{ATPase}$ , a parameter that quantifies ATP  
138 consumption (see [Methods](#), specifically the subsection “ATP consumption”, for specific  
139 details). To emulate the experimental set up of [Takahashi et al. \(1998\)](#) in [Figure 2B](#), we  
140 change the boundary value of O<sub>2</sub> from 47.25 μM to 21 μM, which results in a parabolic O<sub>2</sub>  
141 distribution *qualitatively* similar to that reported by [Takahashi et al. \(1998\)](#).

142

143

144 **The model predicts that rate changes in fission and fusion rates do not impact ADP/ATP  
145 ratios in healthy cardiomyocytes**

146 To determine the impact of fission and fusion rates on the average ADP/ATP ratio in healthy  
147 cells, we simultaneously vary our characteristic fission and fusion rates,  $\lambda_{\text{split}}$  and  $\lambda_{\text{fuse}}$ , over a  
148 range of -80% to 200% while holding all other parameters constant at their base values.

149 These characteristic fission and fusion rates are linked to the probability of a fission or fusion  
150 event occurring via the following formulas:

151 
$$p_{\text{fuse}} = 1 - \exp \left[ - \left( \lambda_{\text{fuse}} + \frac{E_s}{I_{E_0} + E_s} \right) \Delta t \right], \text{ and } p_{\text{split}} = \max \left[ 1 - \exp(-\lambda_{\text{split}} M_X \Delta t), \frac{d}{d_c + d} \right].$$

152 In these formulas,  $\Delta t$  is the size of each ABM time step,  $E_s$  describes energetic stress (defined  
153 in Equation 44),  $M_X$  denotes the mitochondrial mass of a given mitochondrial matrix, and  $d$   
154 describes mitochondrial damage. Further details are given in the [Methods](#) section, specifically  
155 the subsection “Agent based model”. Given that *large* changes to either fission or fusion are  
156 likely to be highly deleterious ([Ong et al., 2015](#)), we assume our range of -80% to 200% is  
157 physiologically reasonable. For each fusion or fission rate, we consider the average ADP/ATP  
158 ratio from 5 model runs for a simulated duration of 2 minutes. We find that at basal levels, the  
159 average ADP/ATP ratio is  $9.17 \times 10^{-3}$  (see [Movie S1](#) for a visualization of the ADP/ATP dynamics  
160 predicted by the model). Moreover, we find that varying our fission and fusion rates result in  
161 minimal deviations from our basal ADP/ATP ratio despite inducing changes in mitochondria  
162 network morphology, specifically, the median mitochondrial cluster size (see [Figure 3](#)). This  
163 suggests that ADP/ATP ratios in healthy cells are robust to variations in fission and fusion over  
164 short timeframes. The same pattern of robustness is also observed in the average PCr/ATP  
165 ratio, another bioenergetic parameter that is used to assess cardiac performance.

166

167 To identify the mechanisms that help in maintaining the robustness of ADP/ATP ratios, we  
168 analyzed the average mitochondrial membrane potential, the average concentration of  
169 inorganic phosphate, and the average ATP hydrolysis rate. We found that the net myofibrillar  
170 ATP hydrolysis rate,  $v_{\text{ATPase}} = \frac{X_{\text{ATPase}}}{1 + R \frac{[\text{ATP}][\text{Pi}]}{[\text{ADP}]}}$  (see [Methods](#), specifically the subsection “ATP

171 consumption”, for more details), which is equivalent to mitochondrial ATP synthesis rate in a  
172 steady state, does not vary substantially despite the variation in median mitochondrial cluster  
173 size. Similarly, the average membrane potential was also robust to changes in  $\lambda_{\text{split}}$  and  $\lambda_{\text{fuse}}$ .

174

175 This bioenergetic stability can be attributed to several mechanisms. In our model, energetic  
176 stress modulates the frequency of fission and fusion events. These events alter mitochondria  
177 network morphology, which affects OXPHOS activity. As a result, the ADP/ATP ratio changes  
178 and with it energetic stress. This change is then integrated into our fission and fusion rates,  
179 thus establishing a feedback loop. These feedback controls result in a stable state of  
180 mitochondrial dynamics, whereby bioenergetic parameters do not vary substantially, despite  
181 changes to the characteristic fission and fusion rates (see [Figure 3](#); also see [Movie S1](#)).

182

183 In addition to these control mechanisms, robustness is additionally maintained via  
184 intracellular shuttling, specifically, adenylate kinase shuttling ([Dzeja and Terzic, 2009](#)) and  
185 creatine kinase phosphate shuttling ([Bessman and Geiger, 1981](#); [Meyer et al., 1984](#)). These  
186 shuttles impart an additional layer of robustness to ADP/ATP ratios and ATP hydrolysis rates  
187 for varied characteristic fission and fusion rates.

188 **Model predicts high workloads increase dynamism while hypoxia causes mitochondrial  
189 clustering**

190 *High workloads*

191 To determine how mitochondrial dynamics are acutely affected by an increased workload, we  
192 track the number of fission and fusion events that occur for  $\text{VO}_2$  values ranging from  $80 \text{ }\mu\text{mol}$   
193  $\text{min}^{-1} \text{ g dw}^{-1}$  to  $140 \text{ }\mu\text{mol min}^{-1} \text{ g dw}^{-1}$ . These particular values are motivated by calculations by  
194 [Vendelin et al. \(2000\)](#), who estimate the largest physiological  $\text{VO}_2$  in adult rat hearts to be  
195  $160 \text{ }\mu\text{mol min}^{-1} \text{ g dw}^{-1}$ . As such,  $\text{VO}_2$  values ranging from  $80 \text{ }\mu\text{mol min}^{-1} \text{ g dw}^{-1}$  (50% of the  
196 largest physiological value) to  $140 \text{ }\mu\text{mol min}^{-1} \text{ g dw}^{-1}$  (87.5% of the largest physiological value),  
197 describe high workload conditions. We found that the number of fusion events increased with  
198 workload (see [Figure 4A](#)) for the entire range of  $\text{VO}_2$ . This is a consequence of the gradual rise  
199 in ADP/ATP ratio (see [Figure 2A](#)) which contributes to an elevation of energetic stress, which  
200 in turn increases the rates of fusion and biogenesis. However, higher energetic stress also

201 leads to an increased likelihood of mitochondrial damage (see Equation 48), which would result  
202 in a slight increase in fission (see Equation 47) and membrane depolarization (see [Figure 4B](#);  
203 see also Equation 31). The net effects of higher fission and fusion rates are larger median sizes  
204 of mitochondrial clusters (see [Figure 4B](#)) which is consistent with experimental studies ([Picard](#)  
205 [et al., 2013](#); [Yoo et al., 2019](#)). It is important to note that these studies track changes on a scale  
206 of hours. By contrast, our simulations track changes on a scale of minutes. As such, increases  
207 in the median cluster size may not represent true biogenesis (which occurs on a scale of ~ 23  
208 minutes at basal conditions, [Dalmasso et al. \(2017\)](#)), but rather aggregation as a result of  
209 increases in mitochondrial outer membrane connectivity which can take place within a shorter  
210 time frame (scale of seconds) ([Glancy et al., 2017](#)). Nevertheless, these findings highlight how,  
211 by modulating the frequency of mitochondrial dynamics, mitochondria effectively perform  
212 network maintenance ensuring consistent energetics even in high workload conditions.

213

214 *Hypoxia*

215 To determine how short-term mitochondrial dynamics differs under hypoxic conditions, we  
216 simulate hypoxia and track network fragmentation. Hypoxia is simulated by imposing a constant  
217 concentration of 5  $\mu$ M on the boundary. We found that under hypoxic conditions, mitochondrial  
218 membrane potentials were rapidly depolarized (visualized in [Movie S2](#)), resulting in an increase  
219 in the ADP/ATP ratio (also visualized in [Movie S2](#)). Under our modelling assumptions, this  
220 increases energetic stress (see Equation 44). Consequently, under hypoxic conditions,  
221 mitochondrial networks in our model rapidly increased fusion (see [Figure 4C](#); see also [Movie](#)  
222 [S2](#)). Hypoxic conditions result in certain mitochondrial subnetworks becoming damaged,  
223 resulting in an average increase in fission (see [Figure 4C](#)) over time. Once segregated, healthy  
224 or only mildly damaged mitochondria can fuse to form a robust subnetwork resulting in large  
225 mitochondrial clusters (see [Figure 4D](#)). This suggests that by segregating damaged  
226 subnetworks and fusing together, mitochondria can reduce the spread of dysfunction, thereby  
227 allowing the cell to become more robust to hypoxia. The concept of mitochondria acutely  
228 segregating damaged subnetworks to improve performance has been also observed in the  
229 literature ([Glancy et al., 2017](#)).

230

231 **Bioenergetics are only mildly robust to altered fission and fusion rates in disease states**

232  
233 Finally, we sought to determine if a disease state, such as diabetes, results in enhanced  
234 sensitivity to changes in fission and fusion rates. To answer this, we simulated mitochondrial  
235 dysfunction observed in diabetic cardiomyopathy. More specifically, we leveraged a study by  
236 [Ghosh \(2019\)](#), in which Beard's biophysical model of OXPHOS ([Beard, 2005](#)) is fit to type I  
237 diabetic rat heart data from [Pham et al. \(2014\)](#). That is, to simulate a type I diabetic cell, we  
238 decreased the rate of Complex I and Complex V activity by factors of 0.288 and  $2.72 \times 10^{-4}$   
239 respectively and increased the rate of proton leakage by a factor of 1.75. We then  
240 simultaneously varied our characteristic fission and fusion rates,  $\lambda_{\text{split}}$  and  $\lambda_{\text{fuse}}$ , over a range of  
241 -80% to 200% while simulating a high-intensity workload of  $\text{VO}_2 = 100 \text{ } \mu\text{mol min}^{-1} \text{ g dw}^{-1}$ .

242

243 We found that modifications to the rates of fission and fusion still did not markedly improve  
244 bioenergetics as mediated by mitochondria network morphology (see [Figure 5](#)), despite an  
245 increase in the median cluster size. Our simulations suggest that increases in fission, which  
246 decrease the median cluster size, are compensated for by an increase in the mitochondrial  
247 area fraction. The converse, however, does not appear to be true, i.e., increases in fusion  
248 activity do not decrease the mitochondrial area fraction. Importantly, these two compensatory  
249 changes in network morphology may help regulate bioenergetics in damaged situations:  
250 increased area fractions may enhance bioenergetics by increasing the availability of ATP in the  
251 myofibrils ([Ghosh, 2019](#)); while increased cluster sizes cause increases in connectivity and  
252 thus enhance OXPHOS (as defined in Equation 30, see subsection "ATP production via  
253 OXPHOS" within the Methods section for more details). As a result of these feedback  
254 mechanisms, the cell maintains an average membrane potential that is robust to changes in  
255 fission and fusion. Notably, simulating diabetes does result in the average concentration of Pi  
256 being more sensitive (relative to our simulations at basal conditions) to changes in fission and  
257 fusion. This is a consequence of Pi regulating mitochondrial metabolism to a greater degree  
258 than the ADP/ATP ratio. However, given that type I diabetes is a chronic condition, we accept  
259 that on a longer timescale, promoting elongation via fusion – which in our computational model  
260 lowers energetic stress – may protect the cell from further damage, either as a result of impaired  
261 OXPHOS function or due to external stressors.

262 **DISCUSSION**

263 In this study, we develop a semi-mechanistic model to quantitatively explore the range of fission  
264 and fusion behaviors that may help with ATP distribution. Our modelling reveals that varied  
265 fusion and fission rates do not result in substantial changes to ADP/ATP ratios in  
266 cardiomyocytes in the short term. Furthermore, our modelling shows that changes in  
267 connectivity alone do not have an immediate impact on bioenergetics as has been suggested  
268 in the literature ([Hoitzing et al., 2015](#)).

269

270 **Our model highlights the robustness of bioenergetics to changes in mitochondria**  
271 **OXPHOS and fusion/fission properties:** Scholars debate the link between ATP synthesis  
272 and mitochondrial dynamics. For example, [Cipolat et al. \(2006\)](#); [Frezza et al. \(2006\)](#); [Olichon](#)  
273 [et al. \(2003\)](#); and [Gilkerson et al. \(2003\)](#), propose that mitochondrial networks can increase  
274 ATP production because of changes in membrane shape. [Parra et al. \(2011\)](#) propose that by  
275 more uniformly distributing mitochondrial membrane potentials, increased connectivity may  
276 improve bioenergetics. In direct contrast, [Hoitzing et al. \(2015\)](#) suggests that mitochondrial  
277 dynamics may have no function in relation to increased ATP production. Our model does not  
278 involve parameters for individual mitochondrial morphology but identifies two levels of  
279 metabolic robustness to changes in fission and fusion rates: mitochondria network morphology  
280 and bioenergetics. At the network morphology level, changes in connectivity affect  
281 bioenergetics, which controls energetic stress. These changes in stress then modulate the  
282 rates of fission and fusion, which establishes a feedback loop, resulting in a stable state of  
283 mitochondrial dynamics. At the bioenergetic level, intracellular shuttles such as adenylate  
284 kinase shuttling ([Dzeja and Terzic, 2009](#)) and creatine phosphate shuttling ([Bessman and](#)  
285 [Geiger, 1981](#); [Meyer et al., 1984](#)) mediate energetic buffering. These feedback mechanisms  
286 result in ADP/ATP and PCr/ATP ratios robust to rate changes. Arguably, changing our model  
287 of energetic stress to depend on more dynamic factors may reduce this robustness. Future  
288 work will address this by using network motifs ([Milo et al., 2002](#); [Li et al., 2014](#)), to identify  
289 factors that when incorporated into our stress calculation, would increase the sensitivity of  
290 ADP/ATP to varied fusion and fission rates.

291

292 **Our model is not a complete representation of the cross-talk between energetics and  
293 mitochondrial dynamics:**

294 Our model does not account for the pleiotropic effects of fission and fusion on cellular  
295 architecture, nor does it account for possible changes in signaling pathways, which may serve  
296 as an additional energetic buffer or perhaps even a compensatory role in bioenergetics. For  
297 example, in the context of acute ischemic reperfusion injury, [Hall et al. \(2016\)](#) note that inhibiting  
298 fusion proteins disrupts the tethering between mitochondria and the sarcoplasmic reticulum,  
299 but paradoxically has a cardioprotective effect. Investigating how mitochondrial dynamics  
300 reshape cellular architecture is a key area that we will explore in future work. The model  
301 representation of mitochondrial networks in two-dimensions and the reduced order  
302 representation of individual mitochondrion morphology also remove the possibility to interrogate  
303 the role of mitochondrion size and shape on bioenergetics. Nevertheless, the model provides  
304 insights on the role that mitochondria fusion/fission dynamics may confer based on current  
305 understanding of the relationship between energetics and mitochondrial dynamics.

306

307 It is possible that modulation of fusion and fission may indirectly or directly affect mitochondrial  
308 expression of respiratory complexes, which could then affect ADP/ATP ratios more drastically.  
309 For example, in our model, we assumed that enzyme activity responds linearly to changes in  
310 connectivity. However, the transportation rate of metabolites from the IMM to the IMS via ANT  
311 saturates out for large ATP concentrations ([Beard, 2005](#)), which when coupled with PCr  
312 shuttling, maintains a stable ADP/ATP ratio ([Bessman and Geiger, 1981](#); [Meyer et al., 1984](#)).  
313 Indeed, while dramatic ( $10^2$  to  $10^4$  fold) decreases in enzyme activity, comparable to chronic  
314 disease conditions ([Wu et al., 2007](#)), do increase the average ADP/ATP ratio across the cell in  
315 our simulations, they do not substantially decrease the average mitochondrial membrane  
316 potential (see Table 5). Thus, implementing a larger change in enzyme activity due to  
317 mitochondrial connectivity is unlikely to *qualitatively* change our findings in the present model.

318

319 Finally, previous modelling work conducted by [Dalmasso et al. \(2017\)](#) suggests that  
320 mitochondrial populations establish and maintain homoeostasis, not by fission and fusion, but  
321 rather by mitochondrial motility. However, in cardiomyocytes mitochondria are organized into

322 parallel columns extending along the length of the cell, which impairs motility (Cao and Zheng,  
323 2019). Specifically, Eisner et al. (2017) observe that mitochondria in cardiomyocytes do not  
324 exhibit motility *in vivo*. However, the impact of cross-sectional network morphology on  
325 bioenergetics is still unresolved in our two-dimensional model. Additionally, our model does not  
326 necessarily distinguish between increased connectivity due to increased tethering of inter-  
327 mitochondrial junctions (IMJ) versus increased connectivity due to mitochondrial fusion. For  
328 example, Picard et al. (2013) note that acute exercise increases both IMJ tethering and  
329 mitochondrial mass, without an increase in fission or fusion (as quantified by the expression of  
330 fusion proteins such as Mfn1, Mfn2, and Opa1, and fission proteins such as Drp1 and Fis1).  
331 Future work will address this by using a finite element method to generalize the model to three  
332 dimensions, and then metabolically coupling it to an experimentally validated model of IMJ  
333 coupling that we intend to develop.

334

335 **Our model provides several experimentally testable predictions:** Firstly, our model  
336 simulations predict that changes in fission and fusion activity – at least on a timescale of two  
337 minutes – do not substantially affect bioenergetics, namely,  $\Delta\Psi$ . Additionally, our model  
338 predicts that as workload, quantified via  $VO_2$ , increases, so too does the frequency of fission  
339 and fusion events across the cell, and the median mitochondrial cluster size. While hypoxia is  
340 generally understood to induce mitochondrial fragmentation (when studied *in vitro* on a  
341 timescale of hours), our simulations suggest that this is preceded by a brief moment of  
342 mitochondrial aggregation. In other words, during hypoxia, mitochondria may acutely, i.e., on a  
343 timescale of two minutes, aggregate together before fragmenting or undergoing elongation. We  
344 emphasize that as simulations from a mathematical model, our results are hypothetical and as  
345 such, highlight the need for systematic quantitative measurements of mitochondrial dynamics.  
346 Indeed, we will refine our model along with our model assumptions as more experimental data  
347 becomes available.

348

349 **Our model highlights the need for quantitative, mechanistic understanding of**  
350 **mitochondrial dynamics to identify pathways for novel therapies:** For example, supposing  
351 fission and fusion events do modulate ETC activity in cardiomyocytes at basal conditions, then  
352 how large a change in  $\Delta\Psi$  – either directly or via changes in enzyme activity in the ETC – do

353 we observe? Moreover, our model assumes that we can induce a set change in fission and  
354 fusion activity. Experimentally, however, this is challenging in part because mitochondria can  
355 change their shape without necessarily increasing their expression of fission and fusion  
356 proteins (Picard et al., 2013). This leads to an additional question – can we induce changes in  
357 fission and fusion activity in a manner that is decoupled from inducing deleterious change in  
358 cell function (e.g., hypoxia or UV-induced damage). And finally, when and under what  
359 circumstances do mitochondria “switch” from fusion-dominated dynamics to fission-dominated  
360 dynamics, to minimize cellular stress. These experiments will provide critical insights into how  
361 mitochondrial form and cardiac metabolism are linked, and as a consequence help either  
362 robustly validate or negate our model’s findings with solid quantitative data.

363

364 In conclusion, our model suggests that ATP synthesis is robust to changes in fission and fusion  
365 rates. By combining experimental data with a system of mathematical equations, we developed  
366 a model that accounts for what has been speculated in the literature. We demonstrated that  
367 mitochondria achieve this robust adaptability by dynamically upregulating the number of fission  
368 and fusion events using a simple feedback-feedforward mechanism. Our modelling results  
369 suggest that changes in ATP synthesis might stem from changes to the respiratory-chain  
370 machinery caused by fission or fusion events. Indeed, our study leads to an interesting  
371 question, if in both healthy and damaged cardiomyocytes the ADP/ATP ratio is robust to  
372 changes in fission and fusion, when and under what circumstances are bioenergetics  
373 impacted?

374

## 375 **ACKNOWLEDGMENTS**

376 The authors would like to thank Siavash B. Kalkhoran and Derek J. Hausenloy for their critical  
377 reading of this manuscript. A.K. was supported by an Australian Government Research  
378 Training Program (RTP) Scholarship; P.S.K. was supported by an Australian Research  
379 Council Discovery Project [DP180101512]; and S.G. and V.R. were supported by an  
380 Australian Research Council Discovery Project [DP170101358].

381 **AUTHOR CONTRIBUTIONS**

382 Conceptualization, A.K., S.G., P.S.K., and V.R.; Methodology, A.K., S.G., and V.R.; Software,  
383 A.K. and S.G.; Formal Analysis, A.K. and S.G.; Writing – Original Draft, A.K.; Writing – Review  
384 & Editing, A.K., S.G., P.S.K., and V.R.; Supervision, P.S.K. and V.R.

385 **DECLARATION OF INTERESTS**

386 The authors declare no competing interests.

387 **FIGURE LEGENDS**

388 **Figure 1. Initial model geometry**

389 (A) Initial conditions used for simulations. Colors denote locally unique mitochondrial clusters.  
390 (B) Block diagram depicting the model set up. Starred boxes indicate variables that we vary in  
391 this study.

392

393 **Figure 2. The model reproduces dynamics from the experimental literature**

394 (A) Comparison between our model predictions for spatially averaged ADP vs VO<sub>2</sub> against  
395 data from [Vendelin et al. \(2000\)](#). (B) Comparison between model predictions for radial oxygen  
396 profiles against data from [Takahashi et al. \(1998\)](#).

397

398 **Figure 3. Bioenergetics are robust to rate changes**

399 The average ADP/ATP ratio; average concentration of inorganic phosphate; average  
400 membrane potential; average ATP hydrolysis rate; average mitochondrial area fraction; and  
401 the median mitochondrial cluster size for varied characteristic fission and fusion rates. Colors  
402 denote the percentage deviation from the basal value (no changes to fission and fusion).

403

404 **Figure 4. Mitochondrial dynamics under high workloads and hypoxia**

405 (A) Mitochondrial dynamics for various workloads. (B) Median mitochondrial cluster size and  
406 mean mitochondrial membrane potential for various workloads. (C) Mitochondrial dynamics in  
407 hypoxic conditions. Rates represent averages across the cell for the duration of the  
408 simulation. (D) Visualization of the mitochondrial network in hypoxic conditions.

409

410 **Figure 5. Bioenergetics are mildly robust to rate changes when simulating**  
411 **mitochondrial dysfunction as observed in diabetic cardiomyopathy**  
412 The average ADP/ATP ratio; average concentration of inorganic phosphate; average  
413 membrane potential; average ATP hydrolysis rate; average mitochondrial area fraction; and  
414 the median mitochondrial cluster size for varied characteristic fission and fusion rates. Colors  
415 denote the percentage deviation from the basal value (no changes to fission and fusion).

416 **TABLES**

417 **Table 1. List of variables used in PDE model**

Variable	Description
$t$	Time (seconds)
ATP	ATP concentration ( $\mu$ M)
MgATP	Mg bound ATP concentration ( $\mu$ M)
ADP	ADP concentration ( $\mu$ M)
MgADP	Mg bound ADP concentration ( $\mu$ M)
AMP	AMP concentration ( $\mu$ M)
PCr	Phosphocreatine concentration ( $\mu$ M)
Cr	Creatine concentration ( $\mu$ M)
Pi	Inorganic phosphate concentration ( $\mu$ M)
$O_2$	Oxygen concentration ( $\mu$ M)
$H^+$	$H^+$ concentration (also expressed in pH)
$K^+$	Potassium ion concentration ( $\mu$ M)
$Mg^{2+}$	Free magnesium ion concentration ( $\mu$ M)
NADH	NADH concentration ( $\mu$ M)
NAD	NAD concentration ( $\mu$ M)
Q	Ubiquinone concentration ( $\mu$ M)
$QH_2$	Ubiquinol concentration ( $\mu$ M)
Cred	Cytochrome C (reduced) concentration ( $\mu$ M)
Cox	Cytochrome C (oxidized) concentration ( $\mu$ M)
$\Delta\Psi$	Mitochondrial membrane potential (mV)

418

419 **Table 2. Details of mitochondrial reactions**

Symbol for flux	Description	Source
$v_{\text{ATPase}}$	ATP consumption in myofibrils	<a href="#">Wu et al. (2008)</a>
$v_{\text{CK}}$	Creatine kinase reaction	<a href="#">Vendelin et al. (2000)</a>
$v_{\text{AK}}$	Adenylate kinase reaction	<a href="#">Vendelin et al. (2000)</a>
$v_{\text{DH}}$	Dehydrogenase flux representing the TCA cycle and other NADH-producing reactions	<a href="#">Beard (2005)</a>
$v_{\text{C1}}, v_{\text{C2}}, v_{\text{C3}}, \text{ and } v_{\text{C5}}$	Flux through complex I, complex III, complex IV, and complex V ( $F_0F_1$ - ATP synthase)	<a href="#">Beard (2005)</a>
$v_{\text{leak}}$	Flux of proton leak across the inner membrane	<a href="#">Beard (2005)</a>
$v_{\text{ANT}}$	Rate of exchange of metabolites through adenine nucleotide translocases (ANT)	<a href="#">Beard (2005)</a>
$v_{\text{PiH}}$	Flux through Phosphate Hydrogen co-transporter	<a href="#">Beard (2005)</a>
$v_{\text{KH}}$	Flux through $\text{K}^+/\text{H}^+$ antiporter	<a href="#">Beard (2005)</a>
$v_{\text{mtCK}}$	Flux of mitochondrial creatine kinase reaction	<a href="#">Ghosh (2019)</a>
$v_{\text{MiAK}}$	Flux of mitochondrial adenylate kinase reaction	<a href="#">Vendelin et al. (2000)</a>

420

421 **Table 3. Parameter estimates for the PDE and ABM models**

Parameter	Description	Estimate	Source
$D_{ANP}$	Diffusivity of ATP, ADP, and AMP	30 $\mu\text{m}^2\text{s}^{-1}$	Simson et al. (2016)
$D_{PCr}, D_{Cr}$	Diffusivity of PCr and Cr	260 $\mu\text{m}^2\text{s}^{-1}$	Vendelin et al. (2000)
$D_{Pi}$	Diffusivity of Pi	327 $\mu\text{m}^2\text{s}^{-1}$	Meyer et al. (1984)
$D_{O_2}$	Diffusivity of O <sub>2</sub>	300 $\mu\text{m}^2\text{s}^{-1}$	Rumsey et al. (1990)
$PCr_{\text{total}}$	Total concentration of PCr and Cr in myofibrils and inner membrane space	23 mM	Vendelin et al. (2000)
$ANP_{\text{total}}$	Total concentration of ATP, ANP, and ADP in the cell	10 mM	Vendelin et al. (2000)
$K_{DT}$	Mg <sup>2+</sup> dissociation constant for myofibrillar ATP	24 $\mu\text{M}$	Vendelin et al. (2000)
$K_{DD}$	Mg <sup>2+</sup> dissociation constant for myofibrillar ADP	347 $\mu\text{M}$	Vendelin et al. (2000)
$K_{DTm}$	Mg <sup>2+</sup> dissociation constant for mitochondrial ATP	17 $\mu\text{M}$	Vendelin et al. (2000)
$K_{DDm}$	Mg <sup>2+</sup> dissociation constant for mitochondrial	282 $\mu\text{M}$	Vendelin et al. (2000)

	ADP		
$NAD_{total}$	Total matrix NAD(H) concentration	2970 $\mu\text{M}$	Beard (2005)
$Q_{total}$	Total matrix ubiquinol concentration	1350 $\mu\text{M}$	Beard (2005)
$x_{buff}$	Constant representing the buffering capacity of the matrix space	100 $\text{M}^{-1}$	Beard (2005)
$C_{total}$	Total IMS cytochrome C concentration	2700 $\mu\text{M}$	Beard (2005)
$C_{IMS}$	Capacitance of the inner membrane	1 $\mu\text{M/L/mV}$	Beard (2005)
$R_{exch}$	Coefficient of restricted ATP diffusion	0.01	Aliev and Saks (1997)
$W_{microcomp}$	micro compartment volume per total mitochondrial volume	0.1	Aliev and Saks (1997)
$W_M$	Water volume per total mitochondrial volume	0.72376	Beard (2005)
$W_{IMS}$	IMS water volume per total mitochondrial volume	$0.1W_M$	Beard (2005)

$W_X$	Matrix water volume per total mitochondrial volume	$0.9W_M$	Beard (2005)
$I_{E_0}$	Stress saturation constant	$5 \times 10^{-2}$	Assumption validated against Vendelin et al. (2000)
$\lambda_{\text{fuse}}$	Characteristic mitochondrial fusion rate	$1.67 \times 10^{-2} \text{ s}^{-1}$	Eisner et al. (2017)
$M_{\text{max}}$	Maximum mitochondrial cluster size	$43 \mu\text{m}^2$	Estimate validated against Vendelin et al. (2000) and Takahashi et al. (1998)
$\lambda_{\text{biogen}}$	Characteristic mitochondrial biogenesis rate	$5.77 \times 10^{-4} \text{ s}^{-1}$	Estimate from Dalmasso et al. (2017)
$\lambda_{\text{split}}$	Characteristic mitochondrial fission rate	$2 \times 10^{-3} \text{ s}^{-1}$	Estimate
$d_c$	Intensity of damage-induced fission	250	Estimate
$M_{\text{single}}$	Mass of single mitochondrion	$8.25 \times 10^{-1} \mu\text{m}^2$	Estimate validated against Vendelin et al. (2000) and Takahashi et al. (1998)

$M_0$	Average mass of mitochondrial cluster	22 $\mu\text{m}^2$	Estimate validated against <a href="#">Vendelin et al. (2000)</a> and <a href="#">Takahashi et al. (1998)</a>
$p_d$	Max probability of spontaneous damage	0.01	Estimate
$k_{\text{damage}}$	Progressive damage rate	$10^{-3} \text{ s}^{-1}$	Estimate
$T_d$	Damage threshold	10	Estimate
$p_{\text{death}}$	Probability of a mitochondrial agent “dying”	0.6	Estimate

422

423

**Table 4. Initial conditions for PDE model**

State Variable	Initial value	Compartment
ATP	9638.08 $\mu\text{M}$	Myofibrils
	8964.91 $\mu\text{M}$	IMM
	9610.35 $\mu\text{M}$	IMS
ADP	92.5336 $\mu\text{M}$	Myofibrils
	1035.09 $\mu\text{M}$	IMM
	113.792 $\mu\text{M}$	IMS
ATP <sub>G</sub>	200007.8 $\mu\text{M}$	IMM
AMP	6.61147 $\mu\text{M}$	Myofibrils
	13.0731 $\mu\text{M}$	IMM
	13.0731 $\mu\text{M}$	IMS
PCr	12000.1 $\mu\text{M}$	Myofibrils
	11998.6 $\mu\text{M}$	IMM

	11998.6 $\mu$ M	IMS
Cr	12000 $\mu$ M	Myofibrils
	12001.4 $\mu$ M	IMM
	12001.4 $\mu$ M	IMS
Pi	2341.77 $\mu$ M	Myofibrils
	2058.61 $\mu$ M	IMM
	2306.24 $\mu$ M	IMS
O <sub>2</sub>	47.25 $\mu$ M	Myofibrils
	47.25 $\mu$ M	IMM
	47.25 $\mu$ M	IMS
K <sup>+</sup>	0.15 M	Myofibrils
	137085 $\mu$ M	IMM
	0.15 M	IMS
Mg <sup>2+</sup>	1 mM	Myofibrils
	3.8×10 <sup>2</sup> mM	IMM
	1 mM	IMS
H <sup>+</sup>	0.072530 pH	IMM
NADH	1552.31 $\mu$ M	IMM
QH <sub>2</sub>	597.711 $\mu$ M	IMM
Cred	439.884 $\mu$ M	IMM
ΔΨ	174.993 mV	IMM

424

425 **Table 5. Very large changes to enzyme activity with basal fission and fusion rates**

Fold change in complex I	Fold change in complex III	Fold change in complex IV	Fold change in complex V	ADP/ATP ratio ( $\times 10^3$ )	ΔΨ	Mitochondrial volume fraction	Median cluster size
1	1	1	1	9.1677	163.75	0.37539	9.24
10 <sup>-4</sup>	10 <sup>-3</sup>	10 <sup>-3</sup>	10 <sup>-2</sup>	53.68	125.69	0.41369	19.305
10 <sup>-4</sup>	10 <sup>-2</sup>	10 <sup>-2</sup>	10 <sup>-3</sup>	30.17	135.76	0.40281	16.83
10 <sup>-3</sup>	10 <sup>-4</sup>	10 <sup>-4</sup>	10 <sup>-2</sup>	115.67	119.65	0.43929	22.44
10 <sup>-3</sup>	10 <sup>-2</sup>	10 <sup>-2</sup>	10 <sup>-4</sup>	37.96	147.12	0.40304	16.17

10 <sup>-2</sup>	10 <sup>-4</sup>	10 <sup>-4</sup>	10 <sup>-3</sup>	98.13	118.42	0.42791	19.8
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426

427 **METHODS**

428 **Lead contact and materials availability**

429 Further information and requests for resources and reagents should be directed to and will be  
430 fulfilled by the Lead Contact, Vijay Rajagopal ([vijay.rajagopal@unimelb.edu.au](mailto:vijay.rajagopal@unimelb.edu.au)).

431 **Experimental model and subject details**

432 *Animals*

433 The initial geometry of the model was inspired from an image of a longitudinal section of the  
434 cell. This image was acquired as part of a three-dimensional stack of electron microscopy  
435 images of a block of cardiac tissue from the left ventricular wall of an adult male Sprague  
436 Dawley rat. Details of the tissue preparation and imaging protocol used to collect these  
437 images can be found in [Hussain et al. \(2018\)](#). The longitudinal image was subsequently  
438 processed to identify and demarcate mitochondria boundaries and subsequently used to  
439 initiate the simulations. All animal procedures followed guidelines approved by the University  
440 of Auckland Animal Ethics Committee (for animal procedures conducted in Auckland,  
441 Application Number R826).

442 **Method Details**

443 To quantify the role of mitochondrial network morphology on bioenergetics, we formulate a  
444 hybrid PDE-ABM system. We model biochemical reactions with a system of experimentally  
445 validated reaction-diffusion equations on a rectangular domain  $[0, L] \times [0, H]$ ; and use an  
446 agent-based model to describe changes in mitochondrial network morphology such as fission,  
447 fusion and biogenesis. We assume a constant pH of 7.1 and unless stated otherwise all  
448 fluxes are functions of state variables.

449

450 *Partial differential equation model*

451 **ATP consumption**

452 To model bioenergetics in the myofibrillar region of the cell, we slightly modify the  
453 bioenergetic model of [Ghosh \(2019\)](#) who considers several populations: [ATP], [ADP], and  
454 [AMP], the concentration of adenosine triphosphate, adenosine diphosphate, and adenosine  
455 monophosphate; [Pi], the concentration of inorganic phosphate; [Cr] and [PCr], the  
456 concentration of creatine and phosphocreatine; [O<sub>2</sub>], the concentration of oxygen; and  
457 [MgATP] and [MgADP], the concentration of magnesium-bound ATP and ADP. A table of all  
458 state variables is provided in Table 1. The interactions between these populations are  
459 described with a PDE system:

$$\frac{\partial[\text{ATP}]}{\partial t} = D_{ANP} \nabla^2[\text{ATP}] - v_{CK} + v_{AK} - v_{\text{ATPase}}, \quad (1)$$

$$\frac{\partial[\text{ADP}]}{\partial t} = D_{ANP} \nabla^2[\text{ADP}] + v_{CK} - 2v_{AK} + v_{\text{ATPase}}, \quad (2)$$

$$\frac{\partial[\text{AMP}]}{\partial t} = D_{ANP} \nabla^2[\text{AMP}] + v_{AK}, \quad (3)$$

$$\frac{\partial[\text{PCr}]}{\partial t} = D_{PCr} \nabla^2[\text{PCr}] + v_{CK}, \quad (4)$$

$$\frac{\partial[\text{O}_2]}{\partial t} = D_{O_2} \nabla^2[\text{O}_2], \quad (5)$$

$$\frac{\partial[\text{Pi}]}{\partial t} = D_{Pi} \nabla^2[\text{Pi}] + v_{\text{ATPase}}, \quad (6)$$

$$\frac{\partial[\text{Cr}]}{\partial t} = D_{Cr} \nabla^2[\text{Cr}] - v_{CK}, \quad (7)$$

$$[\text{MgATP}] = \left( \frac{\text{Mg}^{2+}}{K_{DT} + \text{Mg}^{2+}} \right) [\text{ATP}], \quad (8)$$

$$[\text{MgADP}] = \left( \frac{\text{Mg}^{2+}}{K_{DD} + \text{Mg}^{2+}} \right) [\text{ADP}], \quad (9)$$

$$\text{Mg}^{2+} = 1 \text{ mM}, \quad (10)$$

$$\text{K}^+ = 0.15 \text{ M}, \quad (11)$$

460 where the transport of metabolites across the cell is modelled using diffusion. Here,  
461 ATP is hydrolyzed – or consumed – at rate  $v_{\text{ATPase}}$ ; ATP and AMP are catalyzed via  
462 adenylate kinase at rate  $v_{AK}$ ; and creatine is converted into phosphocreatine via the creatine  
463 phosphate shuttle at rate  $v_{CK}$ . Details of these rates are provided in Table 2.

464

465 To approximate the cardiac cycle, we modify the ATP consumption rate used by [Ghosh et al.](#)  
466 (2018)

$$v_{\text{ATPase}} = \frac{X_{\text{ATPase}}}{1 + R \frac{[\text{ATP}][\text{Pi}]}{[\text{ADP}]}} \quad (12)$$

467 by multiplying it with a tent function  $\Lambda(t) = \sum_{j=0}^{\infty} \Lambda_j(t)$ , where

$$\Lambda_j(t) = \begin{cases} \frac{t - 0.18j}{0.03} & 0 \leq t - 0.18j < 0.03, \\ \frac{0.03 - (t - 0.18j)}{0.06} & 0.03 \leq t - 0.18j < 0.06, \\ 0 & 0.06 \leq t - 0.18j < 0.18. \end{cases} \quad (13)$$

468 This function increases linearly from zero to one during the first 30 ms, decreases linearly to  
 469 zero during the next 30 ms, and remains at zero until the end of the cardiac cycle at 180 ms.  
 470 Here,  $X_{\text{ATPase}}$  is a model parameter that quantifies ATP consumption at various workloads  
 471 and  $R$  is a fixed mass-action ratio. Unless stated otherwise, we assume a high-intensity  
 472 workload of  $\text{VO}_2 = 100 \mu\text{mol min}^{-1} \text{ g dw}^{-1}$  corresponding to a value of  $X_{\text{ATPase}} = 5 \times 10^4 \mu\text{M/s}$ .  
 473

#### 474 **ATP production via OXPHOS**

475 In [Ghosh et al. \(2018\)](#), the dynamics inside a mitochondrial matrix are described by two  
 476 separate but metabolically linked PDE systems. One PDE system models the production of  
 477 metabolites via OXPHOS in the inner mitochondrial membrane (IMM), while the other system  
 478 models the transport of these metabolites from the IMM to the inter-membrane space (IMS).  
 479 Once in the IMS, metabolites may diffuse into the myofibrillar region. To link these  
 480 bioenergetic models to the ABM, we modify the OXPHOS model so that ETC enzyme activity  
 481 and proton leakage depend on mitochondrial connectivity.

482  
 483 The production of ATP via OXPHOS in the IMM is described by the following system of  
 484 PDEs:

$$\frac{\partial[\text{NADH}]}{\partial t} = \frac{v_{DH} - v_{C1}}{W_x}, \quad (14)$$

$$\frac{\partial[\text{QH}_2]}{\partial t} = \frac{v_{C1} - v_{C3}}{W_x}, \quad (15)$$

$$\frac{\partial[\text{Cred}]}{\partial t} = 2 \frac{v_{C3} - v_{C4}}{W_{IMS}}, \quad (16)$$

$$\frac{\partial \text{O}_2}{\partial t} = - \frac{v_{C4}}{2W_{IMS}}, \quad (17)$$

$$\frac{\partial [\text{H}^+]}{\partial t} = x_{\text{buff}} \frac{[\text{H}^+]}{W_x} (v_{DH} - 5v_{C1} - 2v_{C3} - 4v_{C4} + (n_A - 1)v_{C5} + 2v_{PiH} + v_{\text{leak}} - v_{KH}), \quad (18)$$

$$\frac{\partial \Delta \Psi}{\partial t} = \frac{4v_{C1} + 2v_{C3} + 4v_{C4} - n_A v_{C5} - v_{ANT} - v_{\text{leak}}}{C_{\text{IMS}}}, \quad (19)$$

$$\frac{\partial [\text{ATP}]}{\partial t} = \frac{v_{C5} - v_{ANT}}{W_X}, \quad (20)$$

$$\frac{\partial [\text{ADP}]}{\partial t} = -\frac{v_{C5} - v_{ANT}}{W_X}, \quad (21)$$

$$\frac{\partial [\text{Pi}]}{\partial t} = \frac{v_{PiH} - v_{C5}}{W_X}, \quad (22)$$

$$\frac{\partial [\text{K}^+]}{\partial t} = \frac{v_{KH}}{W_X}, \quad (23)$$

$$[\text{MgATP}] = \left( \frac{\text{Mg}^{2+}}{K_{DTm} + \text{Mg}^{2+}} \right) [\text{ATP}], \quad (24)$$

$$[\text{MgADP}] = \left( \frac{\text{Mg}^{2+}}{K_{DDm} + \text{Mg}^{2+}} \right) [\text{ADP}], \quad (25)$$

$$[\text{NAD}] = \text{NAD}_{\text{total}} - [\text{NADH}], \quad (26)$$

$$Q = Q_{\text{total}} - [\text{QH}_2], \quad (27)$$

$$[\text{Cox}] = C_{\text{total}} - [\text{Cred}], \quad (28)$$

$$\text{Mg}^{2+} = 3.8 \times 10^2 \text{ } \mu\text{M}, \quad (29)$$

485 Here, ATP is produced via a series of protein complexes: complex I, complex III, complex IV  
 486 and complex V at rates  $v_{C1}, v_{C3}, v_{C4}$ , and  $v_{C5}$ . While it is known that mitochondrial connectivity  
 487 affects the electron transport chain (Fu et al., 2019; Parra et al., 2011; Pernas and Scorrano,  
 488 2016; Youle and van der Bliek, 2012), it is unclear if all or only some complexes are affected.  
 489 Accordingly, we multiply the rates  $v_{C1}, v_{C3}, v_{C4}$ , and  $v_{C5}$  by

$$M_X/M_0, \quad (30)$$

490 where  $M_X$  is the mitochondrial mass of a given matrix and  $M_0$  is the average mass of a  
 491 mitochondrial matrix – thereby assuming that mitochondrial connectivity increases protein  
 492 complex activity.

493  
 494 Mitochondrial connectivity may also modulate proton leakage (Fu et al., 2019; Parra et al.,  
 495 2011; Pernas and Scorrano, 2016; Youle and van der Bliek, 2012). Moreover, mitochondrial  
 496 damage can depolarize membrane potentials via increased proton leakage (Halestrap et al.,

497 2004; Matsuda et al., 2010; Zorov et al., 2014; Zorova et al., 2018; Park et al., 2011). We  
 498 account for these observations by multiplying the rate of proton leakage  $v_{\text{leak}}$ , by

$$d + \frac{2M_0}{M_0 + M_X}. \quad (31)$$

499 The first term in Equation 31 ensures that proton leakage increases with mitochondrial  
 500 damage,  $d$ , while the second term,  $\frac{2M_0}{M_0 + M_X}$ , ensures that mitochondrial connectivity decreases  
 501 proton leakage.

502  
 503 Additionally, ATP is transported to the IMS via adenine nucleotide translocase (ANT) at rate  
 504  $v_{\text{ANT}}$ ; inorganic phosphate is co-transported at rate  $v_{PiH}$ ; and potassium and protons are  
 505 exchanged at rate  $v_{KH}$ . Finally, dehydrogenase flux stemming from the citric acid cycle  
 506 occurs at rate  $v_{DH}$ .

507  
 508 The transport of metabolites in the IMM to the IMS is described by the following system of  
 509 PDEs:

$$\frac{\partial[\text{ATP}]}{\partial t} = D_{\text{ANP}} \nabla^2[\text{ATP}] + v_{\text{MiAK}} - v_{\text{mtCK}} + v_{\text{ANT}}, \quad (32)$$

$$\frac{\partial[\text{ADP}]}{\partial t} = D_{\text{ANP}} \nabla^2[\text{ADP}] + v_{\text{mtCK}} - 2v_{\text{MiAK}} - v_{\text{ANT}}, \quad (33)$$

$$\frac{\partial[\text{AMP}]}{\partial t} = D_{\text{ANP}} \nabla^2[\text{AMP}] + v_{\text{MiAK}}, \quad (34)$$

$$\frac{\partial[\text{PCr}]}{\partial t} = D_{\text{PCr}} \nabla^2[\text{PCr}] + v_{\text{mtCK}}, \quad (35)$$

$$\frac{\partial[\text{Cr}]}{\partial t} = D_{\text{Cr}} \nabla^2[\text{Cr}] - v_{\text{mtCK}}, \quad (36)$$

$$\frac{\partial[\text{O}_2]}{\partial t} = D_{\text{O}_2} \nabla^2[\text{O}_2] - v_{\text{C4}}/2, \quad (37)$$

$$\frac{\partial[\text{Pi}]}{\partial t} = D_{\text{Pi}} \nabla^2[\text{Pi}] - v_{\text{PiH}}, \quad (38)$$

$$[\text{MgATP}] = \left( \frac{\text{Mg}^{2+}}{K_{DTm} + \text{Mg}^{2+}} \right) [\text{ATP}], \quad (39)$$

$$[\text{MgADP}] = \left( \frac{\text{Mg}^{2+}}{K_{DDm} + \text{Mg}^{2+}} \right) [\text{ADP}], \quad (40)$$

$$\text{Mg}^{2+} = 1 \mu\text{M}, \quad (41)$$

$$K^+ = 0.15M. \quad (42)$$

$$\frac{\partial [ATP]_G}{\partial t} = \frac{R_{exch}([ATP] - [ATP]_G) + v_{ANT} - v_{mtCK}}{W_{microcomp}}, \quad (43)$$

510

511 Here, ATP is transported via the protein adenine nucleotide translocase at rate  $v_{ANT}$ ; ATP  
512 and AMP in the IMM are catalyzed via mitochondrial adenylate kinase at rate  $v_{MiAK}$ ; creatine  
513 is converted into phosphocreatine in the IMM at rate  $v_{mtCK}$ ; inorganic phosphate is co-  
514 transported at rate  $v_{PiH}$ ; and oxygen is consumed in the IMM at rate  $v_{C4}/2$ . Equation 43, is a  
515 microcompartment between ANT and mitochondrial CK, that based on previous work by [Aliev](#)  
516 [and Saks \(1997\)](#), models phosphocreatine shuttling. Details of these rates are provided in  
517 Table 2.

518

### 519 **Boundary conditions and initial conditions**

520 As implemented by [Ghosh et al. \(2018\)](#), we impose no-flux boundary conditions (BC) on all  
521 state variables except oxygen, for which we impose a constant Dirichlet BC of 47.25  $\mu\text{M}$  on  
522 the boundary. We use constant initial conditions, with details provided in Table 4.

523

### 524 *Agent based model*

525 Increased ATP demand along with oxidative stress is conducive to mitochondrial fusion and  
526 biogenesis, along with fission of damaged mitochondria ([Dalmasso et al., 2017](#); [Mihaylova](#)  
527 [and Shaw, 2011](#); [Toyama et al., 2016](#); [Egan et al., 2011](#)). To this end, we model changes in  
528 network morphology with an agent-based model.

529

### 530 **Energetic stress**

531 To model biophysical stressors, we introduce the concept of energetic stress

$$E_s = \frac{\langle A_D \rangle}{\langle A \rangle} \frac{1}{M_x} + d. \quad (44)$$

532 The fraction  $\langle A_D \rangle / \langle A \rangle$  is the average ADP to ATP ratio within the mitochondrial matrix, acting  
533 as a measure of biophysical stress. [Parra et al. \(2011\)](#) speculate that increased connectivity  
534 improves bioenergetics by more uniformly distributing the mitochondrial membrane potential.  
535 We account for this by introducing a connectivity penalty to energetic stress of the form  $1/M_x$ ,

536 where  $M_X$  denotes the mitochondrial mass of a given matrix. The final term  $d$  denotes the  
537 level of mitochondrial damage (see [Damage](#))

538

539 **Fusion**

540 We assume that the probability of a mitochondrion undergoing fusion at each time step is  
541 given by

$$p_{\text{fuse}} = 1 - \exp \left[ - \left( \lambda_{\text{fuse}} + \frac{E_s}{I_{E_0} + E_s} \right) \Delta t \right]. \quad (45)$$

542 The size of each time step is denoted by  $\Delta t$ , and the characteristic fusion rate is denoted by  
543  $\lambda_{\text{fuse}}$ . Our characterisation ensures that as energetic stress,  $E_s$ , increases, the probability of  
544 fusion also increases. Here,  $I_{E_0}$  is a stress saturation constant. If a fusion event occurs, our  
545 mitochondrion (hereafter referred to as an agent) will fuse with all adjacent agents, *unless* the  
546 mass of the resultant agent exceeds  $M_{\text{max}}$ .

547

548 **Biogenesis**

549 Similarly, we assume the probability of a mitochondrion undergoing biogenesis at each time  
550 step is given by

$$p_{\text{biogen}} = 1 - \exp \left[ - \left( \lambda_{\text{biogen}} + \frac{E_s}{I_{E_0} + E_s} \right) \Delta t \right]. \quad (46)$$

551 The parameter  $\lambda_{\text{biogen}}$  denotes the characteristic biogenesis rate. If a biogenesis event occurs,  
552 our mitochondrial matrix increases its mass by a single mitochondrion, *unless* the mass of the  
553 resultant matrix exceeds  $M_{\text{max}}$  or if there is no free space. The biogenesis process is  
554 implemented by associating a vacant cell to either the left or right of the current mitochondrial  
555 matrix with our agent. We assume that the contents of the cytosol are pushed away and  
556 distributed equally amongst neighboring cells.

557

558 **Fission**

559 We assume that the probability of a mitochondrion undergoing fission at each time step is  
560 given by

$$p_{\text{split}} = \max \left[ 1 - \exp(-\lambda_{\text{split}} M_x \Delta t), \frac{d}{d_c + d} \right], \quad (47)$$

561 where  $\lambda_{\text{split}}$  is the characteristic fission rate and  $M_x$  the mitochondrial mass of a given matrix.

562 The level of mitochondrial damage is described by  $d \geq 1$ , and the extent to which

563 mitochondrial damage drives fission is described by  $d_c$ .

564

565 Our characterization assumes that in healthy mitochondria, fission occurs independently of

566 energetic stress and is proportional to the number of agents in the given matrix. In [Glancy et](#)

567 [al. \(2017\)](#), the authors note that damaged mitochondria rapidly increase fission to minimize

568 the propagation of mitochondrial dysfunction. We capture this behavior by assuming that the

569 probability of fission in damaged mitochondria is driven by damage according to the term  $\frac{d}{d_c+d}$ .

570 As local damage increases, the probability of fission approaches one. The switch between

571 basal fission and damaged-induced fission occurs when the probability of damaged-induced

572 fission matches the probability of basal-level fission, that is, when  $1 - \exp(-\lambda_{\text{split}} M_x \Delta t) =$

573  $\frac{d}{d_c+d}$ . However, this switching condition is purely phenomenological and as such requires

574 further experimentation to either be phenomenologically refined or replaced with a

575 mechanistic model.

576

577 Suppose our agent has an initial mitochondrial mass of  $M$ . Now let  $M_{\text{single}}$  denote the mass of

578 a single mitochondrion. We assume that fission only occurs if  $M > 2M_{\text{single}}$ . This implies that

579 an agent must consist of more than two linked mitochondria for fission to occur, which is

580 implemented for computational reasons. If fission occurs, the original agent divides into two

581 new agents of mass  $M_1$  and  $M_2$ , where  $M_1$  is chosen as a uniform random variable between

582  $M_{\text{single}}$  and  $M - M_{\text{single}}$  and  $M_2 = M - M_1$ .

583

## 584 **Damage**

585 Mitochondria segregate damaged mitochondria via fission ([Twig et al., 2008](#); [Youle and van](#)

586 [der Bliek, 2012](#); [Glancy et al., 2017](#)). Mitochondrial damage is described by the variable  $d$

587 and is assumed to exist in two states, low and high. We assume that after fission, the

588 probability of the newly separated mitochondria becoming damaged is  $p_{\text{damage}}$ . This probability

589 is defined as

$$p_{\text{damage}} = p_d \frac{E_s}{E_s + I_{E_0}}, \quad (48)$$

590 where the parameter  $p_d$  describes the maximal probability of mitochondrial damage. The  
591 factor  $\frac{E_s}{E_s + I_{E_0}}$  ensures that increased energetic stress results in an increased likelihood of  
592 mitochondrial damage.

593  
594 Damaged mitochondria start from a low-damage state, corresponding to  $d = 1$  and increases  
595 by  $1 - e^{-k_{\text{damage}}\Delta t}$  every time step. Once damage hits a critical threshold  $d = T_d$ , our state  
596 switches from a low-damage state to a high-damage state. Mitochondria that are highly  
597 damaged are assumed to be susceptible to increased turnover ([Hamacher-Brady and Brady, 2016](#)). The probability of mitochondrial turnover is given by  $p_{\text{death}} > 0.5$ . This approach to  
598 modelling mitochondrial damage is similar to the approach utilised by [Dalmasso et al. \(2017\)](#).  
599 Here, turnover is not referring to cell death nor mitophagy *per se*, but rather refers to a  
600 mitochondrial agent dying.

602  
603 Healthy mitochondria are treated as having a zero-damage state, i.e.,  $d = 0$ . If two  
604 mitochondria with damage states of  $d_1$  and  $d_2$  fuse, then the resultant damage is assumed to  
605 be the average of the two, i.e.  $(d_1 + d_2)/2$ . If this value is below one, the mitochondria are no  
606 longer marked as damaged.

607  
608 *Parameter estimates*  
609 A summary of parameter values is provided in Table 3. Where possible, we have used data  
610 from animal models to characterize our estimates; however, some of the available data  
611 comes from *in vitro* models due to the limited availability of animal data. For the PDE model,  
612 we used the flux terms implemented by [Ghosh et al. \(2018\)](#). Details of how these parameters  
613 were estimated are provided therein. We used manual calibration instead of formal parameter  
614 fitting, which in our case is not feasible due to a lack of data directly corresponding to specific  
615 model parameters. We summarize how we obtained these estimates for our ABM parameters  
616 below. Unless stated otherwise, we assume a high-intensity workload of  
617  $\text{VO}_2=100 \text{ }\mu\text{mol min}^{-1} \text{ g dw}^{-1}$ .

618

619 **Agent-based model**

- 620 •  $\lambda_{\text{split}}$ : As a plausible estimate, we use a characteristic fission rate of  $2 \times 10^{-3} \text{ s}^{-1}$ .
- 621 •  $\lambda_{\text{fuse}}$ : Using a murine cardiomyocyte model, [Eisner et al. \(2017\)](#) estimate a  
622 characteristic fusion rate of  $1.67 \times 10^{-2} \text{ s}^{-1}$ .
- 623 •  $\lambda_{\text{biogen}}$ : In [Dalmasso et al. \(2017\)](#) the authors optimize a mathematical model to  
624 estimate a characteristic biogenesis frequency of 28.9 minutes, which equates to  
625  $1/28.9 \text{ min}^{-1} = 5.77 \times 10^{-4} \text{ s}^{-1}$ .
- 626 •  $M_0$ ,  $M_{\text{single}}$ , and  $M_{\text{max}}$ : Based on our image data we estimate  $M_{\text{single}} = 8.25 \times 10^{-1} \mu\text{m}^2$   
627 and  $M_{\text{max}} = 43 \mu\text{m}^2$ . Motivated by this we set  $M_0 = \frac{1}{2}(M_{\text{single}} + M_{\text{max}}) \approx 22 \mu\text{m}^2$ .
- 628 •  $d_c$ ,  $p_d$ ,  $k_{\text{damage}}$ , and  $T_d$ : As plausible estimates, we use  $d_c = 250$ ,  $p_d = 0.01$ ,  
629  $k_{\text{damage}} = 10^{-3} \text{ s}^{-1}$ , and  $T_d = 10$ .
- 630 •  $p_{\text{death}}$ : We arbitrarily set the probability of a mitochondrial agent dying to be 0.6.

631

632 *Simulations*

633 The initial geometry of our mitochondrial network was inspired by tissue samples taken from a  
634 Sprague Dawley rat. Our initial conditions are a simplified 2D electron microscopy  
635 representation of a healthy rat heart and are visualized in [Figure 1A](#). We assume that the  
636 lateral and longitudinal dimensions of our hypothetical cardiomyocyte are 79  $\mu\text{m}$  and  
637 15.75  $\mu\text{m}$  respectively. Our spatial increments,  $\Delta x$  and  $\Delta y$  are taken to be 1  $\mu\text{m}$  and 0.75  $\mu\text{m}$   
638 respectively. These increments coincide with the typical dimensions of a fibre-parallel  
639 mitochondrion, allowing us to model mitochondrial matrices as mesh-points on our domain.  
640 We note that using a different initial condition does not appear to change our results (see  
641 [Figure S1](#)).

642

643 Our PDEs are discretized using the method of lines and solved using semi-implicit Strang  
644 splitting. Our linear and non-linear components are both solved with MATLAB's inbuilt stiff  
645 ODE solver "ode15s" with an absolute error tolerance of  $10^{-6}$ . Our ABM time step,  $\Delta t$ , is set to  
646 be 0.01 s.

647

648 To determine the minimum number of ABM runs, we screen for variance stability (Lee et al.,  
649 2015; Lorscheid et al., 2012), i.e., identify the number of runs required for the coefficient of  
650 variation to be less than some fixed tolerance. We find that with 5 runs, the coefficients of  
651 variation for our ADP/ATP ratios are below  $10^{-2}$ , which is considered acceptable in the  
652 literature (Lee et al., 2015). Thus, the results from our ABM represent an average from 5 runs  
653 unless stated otherwise.

## 654 SUPPLEMENTAL INFORMATION

### 655 **Movie S1. Simulation of the model at baseline conditions.**

656 Inset depicts a moving histogram depicting the distribution of mitochondrial (normalized so  
657 that the total area sums to one). Also depicted are the ADP/ATP ratios; Pi concentrations;  
658 and the  $\Delta\Psi$ 's predicted by the model. These values oscillate for each cardiac cycle. Related  
659 to [Figure 3](#).

### 660 **Movie S2. Simulation of the model under hypoxia.**

661 The value of  $O_2$  on the boundary is 5  $\mu M$  at a workload of  $VO_2 = 100 \mu mol min^{-1} g dw^{-1}$ . Inset  
662 depicts a moving histogram depicting the distribution of mitochondrial (normalized so that the  
663 total area sums to one). Additionally depicted are the ADP/ATP ratios;  $O_2$  concentrations; and  
664 the  $\Delta\Psi$ 's predicted by the model. These values oscillate for each cardiac cycle. Related to  
665 [Figure 4C](#).

### 666 **Figure S1. Sensitivity of ADP/ATP ratio with different initial conditions.**

667 Bioenergetics are robust to changes in the characteristic fission rate,  $\lambda_{split}$ , and fusion rate,  
668  $\lambda_{split}$ , regardless of the initial condition used in the model. Related to [Methods](#).

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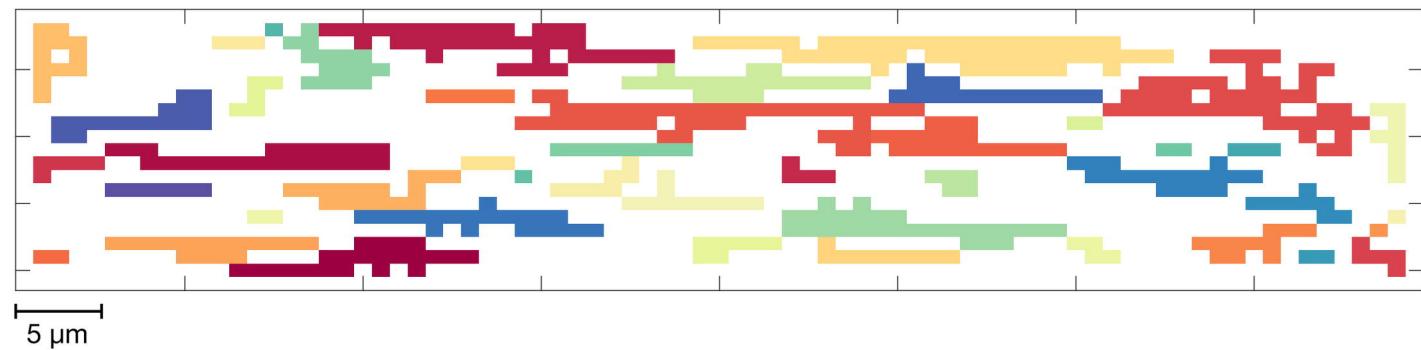
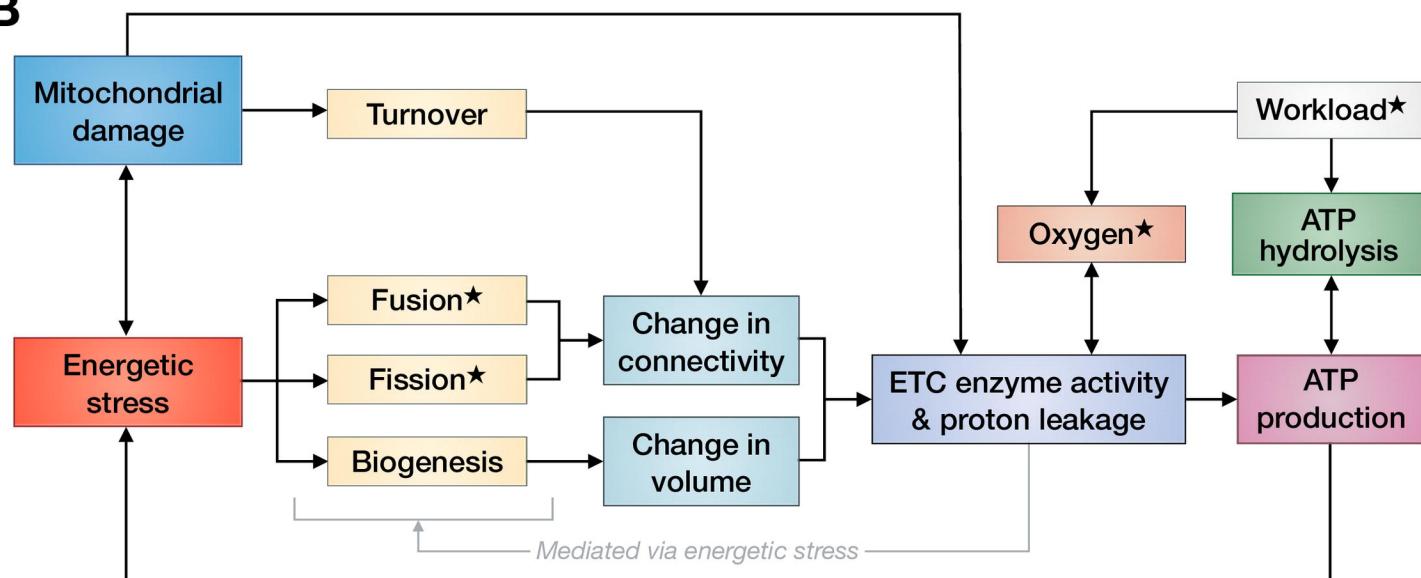
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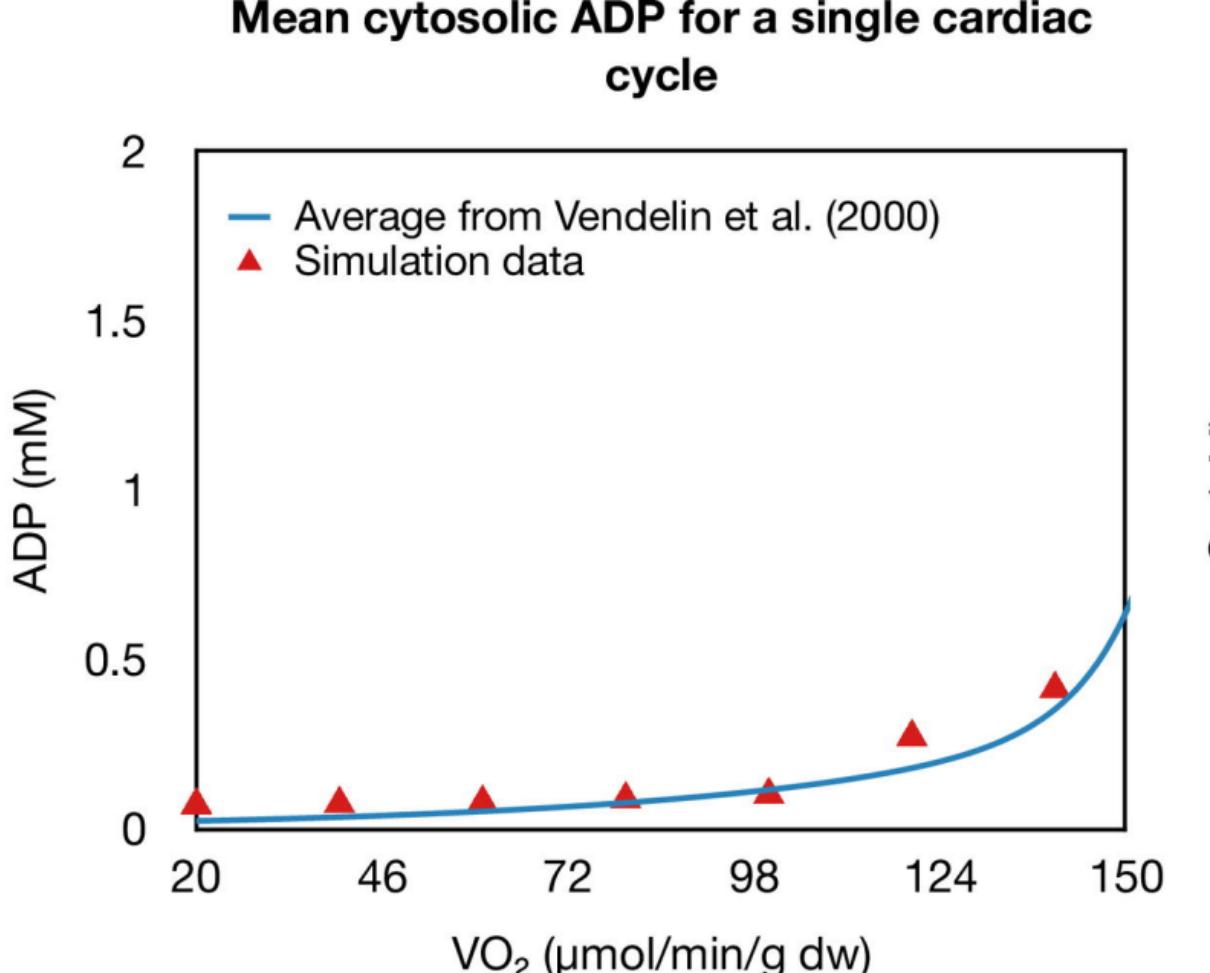
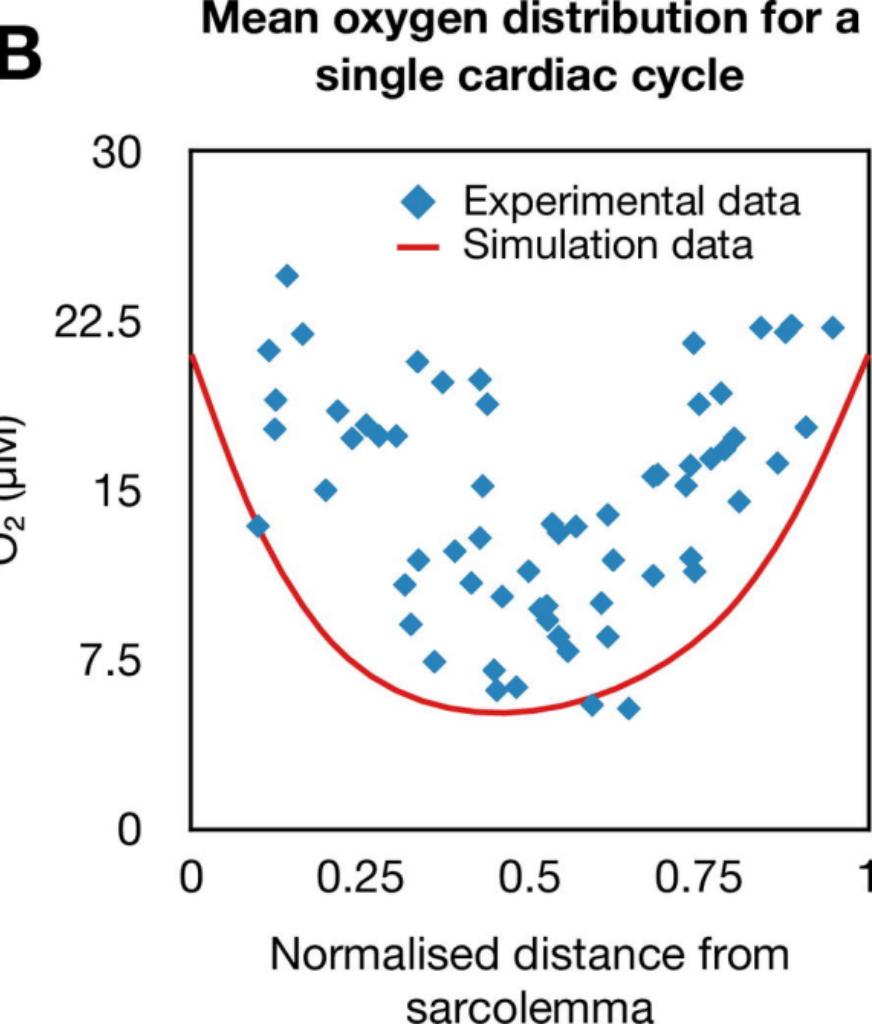
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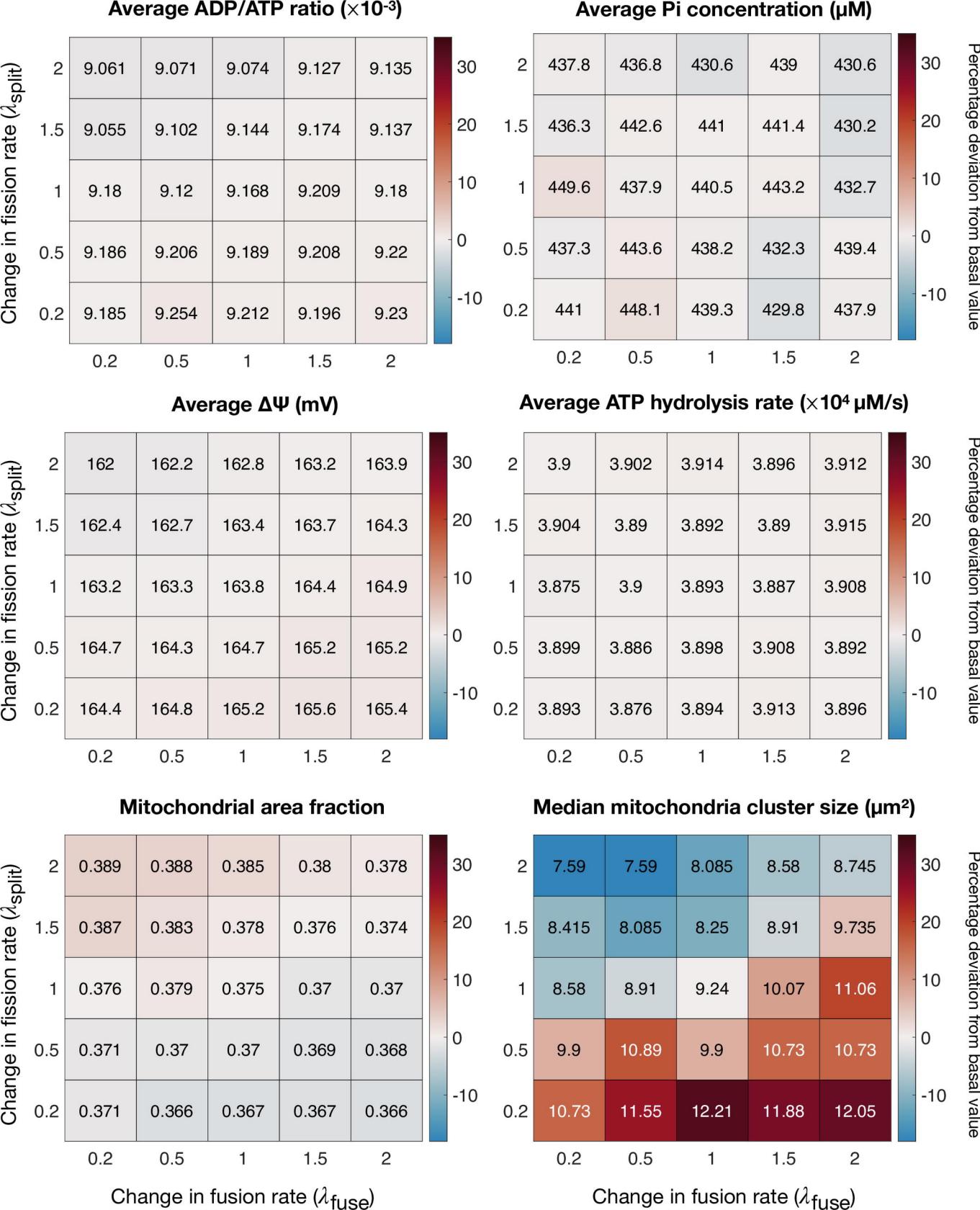
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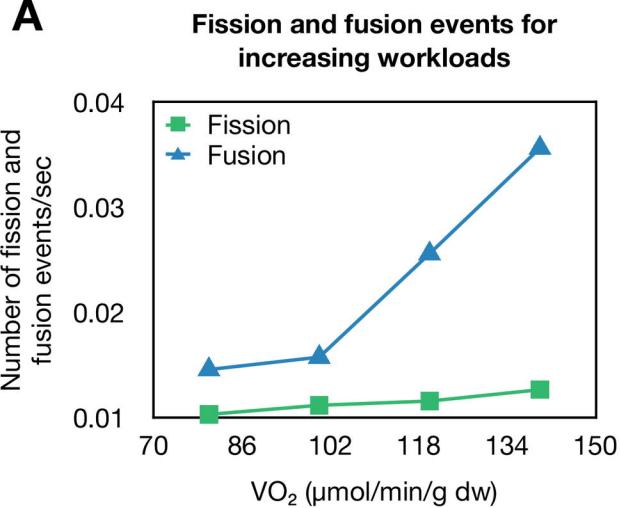
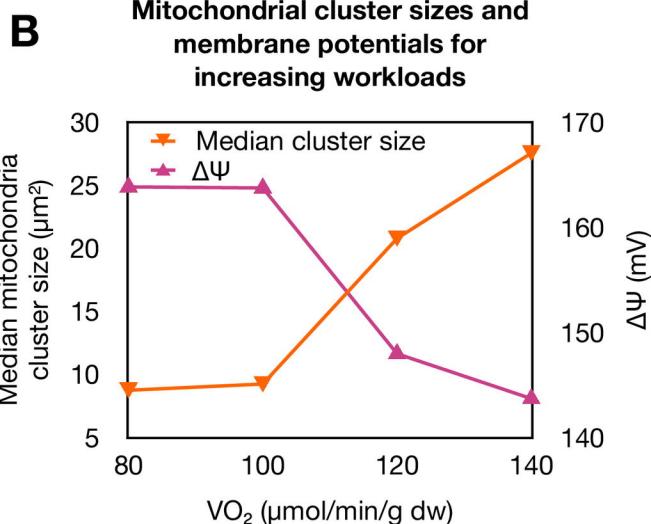
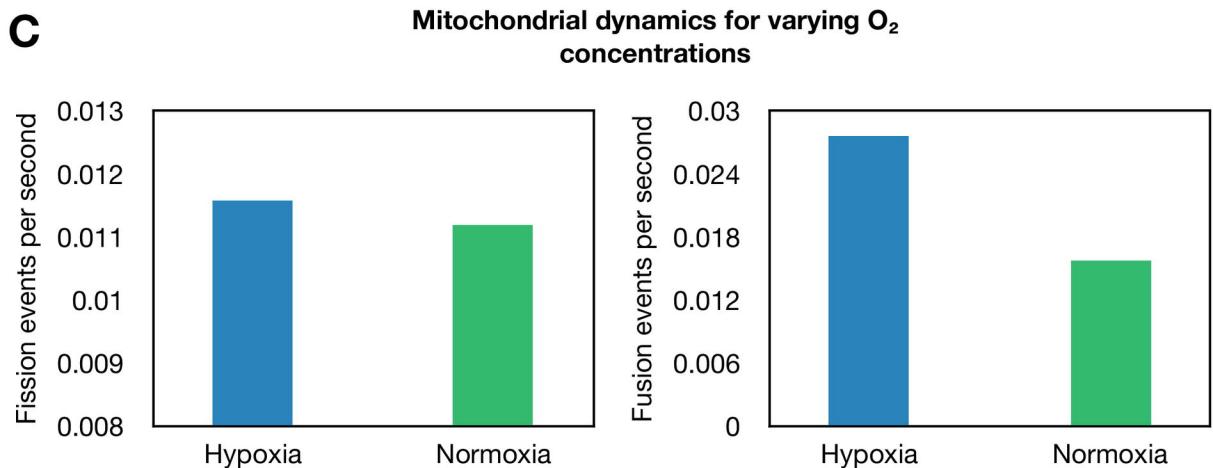
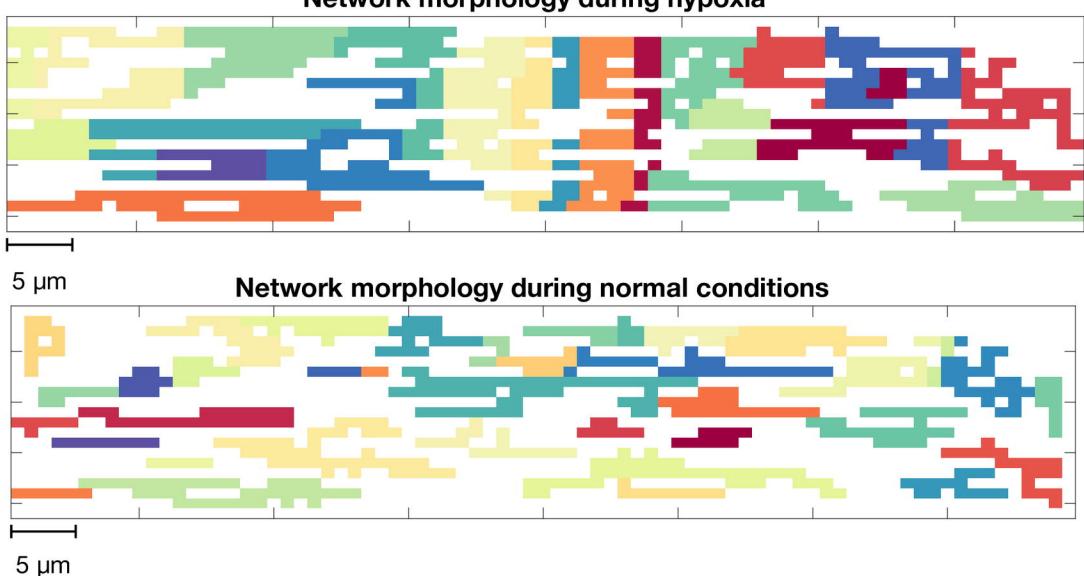
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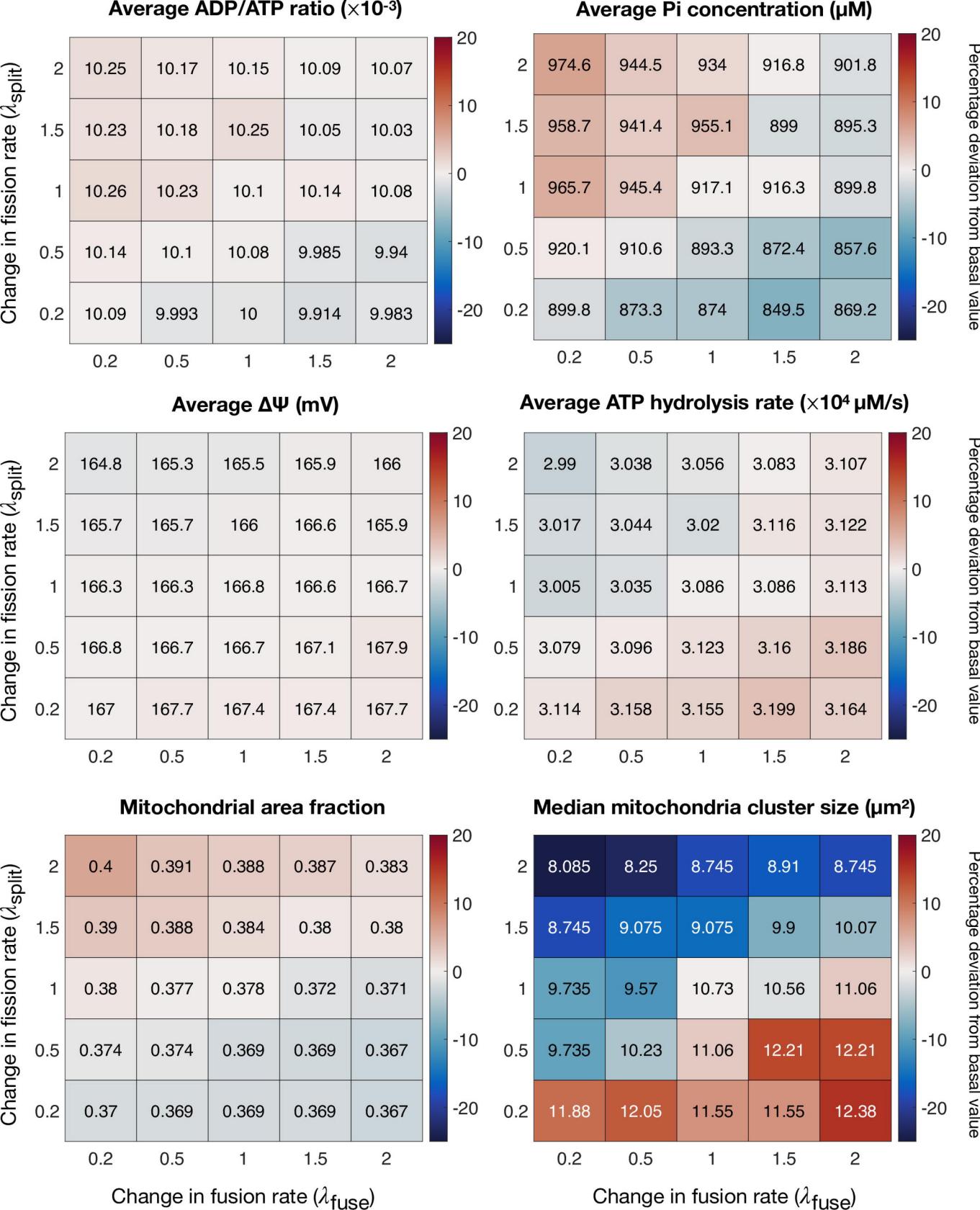
866

**A****B**

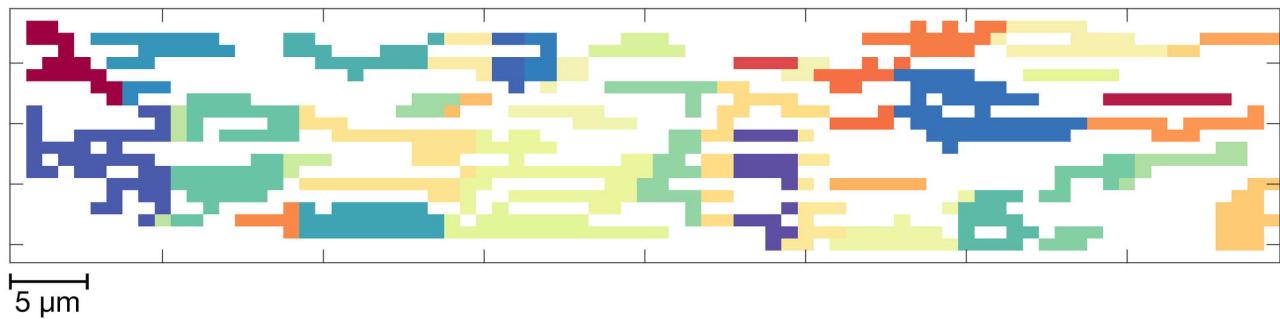
**A****B**



**A****B****C****D**



## New initial condition



## Sensitivity of bioenergetics to characteristic fission and fusion rates

