

# **Increasing the accuracy of exchange parameters reporting on slow dynamics by performing CEST experiments with high $B_1$ fields**

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## Abstract

Over the last decade chemical exchange saturation transfer (CEST) NMR methods have emerged as powerful tools to characterize biomolecular conformational dynamics occurring between a visible major state and ‘invisible’ minor states. The ability of the CEST experiment to detect these minor states, and provide precise exchange parameters, hinges on using appropriate  $B_1$  field strengths during the saturation period. Typically, a pair of  $B_1$  fields with  $\omega_1 (= 2\pi B_1)$  values around the exchange rate  $k_{ex}$  are chosen. Here we show that the transverse relaxation rate of the minor state resonance ( $R_{2,B}$ ) also plays a crucial role in determining the  $B_1$  fields that lead to the most informative datasets. Using  $K = [k_{ex}(k_{ex} + R_{2,B})]^{\frac{1}{2}} \geq k_{ex}$ , to guide the choice of  $B_1$ , instead of  $k_{ex}$ , leads to data wherefrom substantially more accurate exchange parameters can be derived. The need for higher  $B_1$  fields, guided by  $K$ , is demonstrated by studying the conformational exchange in two mutants of the 71 residue FF domain with  $k_{ex} \sim 11 \text{ s}^{-1}$  and  $\sim 72 \text{ s}^{-1}$ , respectively. In both cases analysis of CEST datasets recorded using  $B_1$  field values guided by  $k_{ex}$  lead to imprecise exchange parameters, whereas using  $B_1$  values guided by  $K$  resulted in precise site-specific exchange parameters. The conclusions presented here will be valuable while using CEST to study slow processes at sites with large intrinsic relaxation rates, including carbonyl sites in small to medium sized proteins, amide  $^{15}\text{N}$  sites in large proteins and when the minor state dips are broadened due to exchange among the minor states.

## Keywords

Protein Dynamics; Conformational Exchange; Chemical Exchange; Chemical Exchange Saturation Transfer; FF domain; CEST; NMR

## Introduction

Protein molecules are dynamic entities that at ambient temperature sample various conformational states with differing populations and lifetimes (1, 2). In addition to understanding dynamical processes, such as protein folding/misfolding and aggregation, a knowledge of protein conformational dynamics is often necessary to understand protein function, allostery etc. (2-6). Hence, over the last few decades different classes of NMR experiments have been developed to study protein conformational dynamics occurring on the  $\mu$ s to second time-scale (4, 7-9), including  $R_{1,\rho}$  (10, 11), CPMG (12, 13), CEST (14) and DEST (15). These experiments can detect sparsely populated conformational states that are ‘invisible’ in regular NMR spectra. In all these experiments, the spins are manipulated by pulses whereafter the ‘visible’ major state magnetization is detected and used to reconstruct the spectrum of the ‘invisible’ minor state, which in favorable cases can be used to determine the structures of the minor states (16-20). CEST experiments, originally devised to study slow exchange between visible states (21), are now routinely used to study protein and nucleic acid conformational exchange between a visible major state and invisible minor state(s) occurring over a wide range of time-scales (22-24). CEST methods have been developed to characterize the exchange at various backbone and side-chain sites (25-31) and have been used to study various processes, including protein folding (24, 32), ligand binding (33, 34) and several other processes involving protein and nucleic acid conformational fluctuations (35-37).

In a typical CEST experiment longitudinal magnetization arising from the nucleus of interest is irradiated with a weak  $B_1$  ( $\sim 5$  to  $\sim 300$  Hz) field for a period  $T_{EX}$  of  $\sim 0.25$  to  $\sim 0.6$  s termed the exchange delay, following which the intensity of the visible major-state peak is quantified as a function of the offset at which the  $B_1$  irradiation is applied. When the system of interest consists of a major state, A, in slow exchange with a minor state, B, that is  $A \xrightleftharpoons[k_{AB}]{k_{BA}} B$ , a plot of the normalized intensity ( $I(\omega_{RF})/I_0$ ) versus the offset  $\omega_{RF}$  (ppm) at which the  $B_1$  field is applied will have two dips. These two dips consist of one at the chemical shift (ppm) of the major state ( $\omega_A$ ) and more importantly one at the chemical shift of the minor state ( $\omega_B$ ), which allows one to detect sparsely populated states with fractional populations as low as  $\sim 0.5\%$ .  $I_0$  is the intensity of the major state in the absence of the  $T_{EX}$  delay. The size and width of the minor state dip (largely) depends on the exchange rate ( $k_{ex} = k_{AB} + k_{BA}$ ), the fractional population of the minor state ( $p_B = k_{AB}/k_{ex}$ ), the minor state transverse relaxation rate ( $R_{2,B}$ ) and the value of  $B_1$ . The exchange parameters ( $k_{ex}$ ,  $p_B$ ), the major and minor state chemical shifts, the major ( $R_{2,A}$ ) and the minor-state transverse relaxation rates ( $R_{2,B}$ ), as well as the major state longitudinal relaxation rate ( $R_{1,A}$ ) can all be extracted from the analysis of a pair of CEST profiles recorded with different (suitably chosen)  $B_1$  values (14). For two-state slow exchange ( $k_{ex}/|\Delta\omega_{AB}| < 1$ ) processes considered here, CEST profiles are typically recorded with  $\omega_1$  (rad/s;  $= 2\pi B_1$ ) values guided by  $k_{ex}$ , that is one  $\omega_1$  less than  $k_{ex}$  in the  $0.5k_{ex}$  to  $0.8k_{ex}$  range and one higher than  $k_{ex}$  in the  $1.5k_{ex}$  to  $1.8k_{ex}$  range. Here  $\Delta\omega_{AB} = \omega_B - \omega_A$ , where  $\omega_A$

and  $\omega_B$  are the resonance frequencies (rad/s) of the nucleus of interest in the major and minor states respectively.

The small 71 residue four helix bundle FF domain (38) from human HYPA/FBP11 has served as a model system to understand protein conformational dynamics and folding (39-44). Whilst characterizing the conformational dynamics of the A17G S56P FF domain using methyl  $^{13}\text{C}$  CEST experiments, we found that precise site-specific exchange parameters could not be obtained from the analysis of two CEST datasets recorded with  $\omega_1$  values lower and higher than  $k_{ex}$  ( $\sim 11\text{ s}^{-1}$ ) as described above. We discovered that this problem occurs when the transverse relaxation rate of the minor state,  $R_{2,B}$ , is greater than  $k_{ex}$  and that accurate exchange parameters can be obtained by recording additional CEST datasets with higher  $B_1$  values. We rationalize the benefit of the larger  $B_1$  for deriving accurate exchange parameters by inspecting the equations that govern the shape of the minor state dip in CEST profiles. We conclude that the choice of  $B_1$  values should be informed by  $K = [k_{ex}(k_{ex} + R_{2,B})]^{1/2} (\geq k_{ex})$  rather than  $k_{ex}$ . The efficacy of this strategy is further demonstrated by characterizing the two-state folding reaction ( $k_{ex} \sim 72\text{ s}^{-1}$ ) of the A39G FF domain in which the apparent transverse relaxation rates of several residues in the unfolded (U) state are greater than  $140\text{ s}^{-1}$ . In line with the analysis presented here, choosing  $B_1$  values guided by  $K$  rather than  $k_{ex}$  allows for the accurate determination of the exchange parameters.

## Materials and Methods

### Protein Samples

The A17G S56P FF sample contained of  $\sim 4\text{ mM}$  U- $[^2\text{H}, ^{15}\text{N}]$ , Ile $\delta 1$ - $[^{13}\text{CH}_3]$ , Leu,Val- $[^{13}\text{CH}_3, ^{12}\text{CD}_3]$  labelled protein dissolved in 50 mM acetate, 100 mM NaCl, 30%  $[^2\text{H}]$ -glucose, pH 5.7, 100%  $\text{D}_2\text{O}$  buffer. The A39G FF sample contained  $\sim 4\text{ mM}$  U- $[^{15}\text{N}]$  labelled protein dissolved in 50 mM acetate, 100 mM NaCl, pH 5.7, 10%  $\text{D}_2\text{O}$  buffer. Proteins were overexpressed in *E. coli* BL21(DE3) cells transformed with the appropriate plasmids grown in suitable M9 media (45, 46) and purified as described previously (24, 47, 48).

### NMR Spectroscopy

Methyl  $^{13}\text{C}$  CEST profiles (A17G S56P FF sample,  $7.5^\circ\text{C}$ ) were recorded on a 700 MHz (16.4 T) Bruker Avance III HD spectrometer equipped with a cryogenically cooled triple resonance probe. To accelerate data acquisition,  $^{13}\text{C}$  methyl CEST data was acquired using the DANTE-CEST (D-CEST) sequence (49, 50) that uses the DANTE sequence (51, 52) for RF irradiation during the  $T_{EX}$  period. Amide  $^{15}\text{N}$  CEST profiles (A39G FF sample,  $2.5^\circ\text{C}$ ) were recorded on a 500 MHz (11.7 T) Bruker NEO spectrometer equipped with a room temperature triple resonance probe using the standard amide  $^{15}\text{N}$  CEST sequence (25). During the  $T_{EX}$  delay of both the  $^{13}\text{C}$  and  $^{15}\text{N}$  CEST experiments  $^1\text{H}$  decoupling was carried out using the  $90_x 240_y 90_x$  composite pulse (53) effectively reducing the nucleus of interest (methyl  $^{13}\text{C}$  or the amide

$^{15}\text{N}$ ) to an isolated spin  $\frac{1}{2}$  spin system (25).  $B_1$  fields were calibrated using the nutation method (54). The methyl  $^1\text{H}$ - $^{13}\text{C}$  correlation maps were recorded with 24 complex points (sweep width: 14 ppm) in the indirect ( $^{13}\text{C}$ ) dimension while the amide  $^1\text{H}$ - $^{15}\text{N}$  correlation maps were recorded with 24 complex points (sweep width: 16.9 ppm) in the  $^{15}\text{N}$  dimension. Methyl  $^{13}\text{C}$  CEST data were acquired using 16 scans, whereas 4 scans were used to record the amide  $^{15}\text{N}$  CEST data. Additional details are provided in Table S1.

## Data Analysis

The NMRPipe package (55) was used to process the NMR data, Sparky (56, 57) was used to visualize and label the spectra while the program PINT (58) was used to obtain peak intensities from the spectra. Uncertainties in the peak intensities were estimated based on the scatter in the flat part of the CEST intensity profiles (23). The software package *ChemEx* (59) that numerically integrates (60) the Bloch-McConnell equations (61) was used to both obtain the best fit exchange parameters from the experimental (or synthetic) data and to generate the synthetic CEST profiles (Fig. 2 & S2). The two-state fitting parameters included the major and minor state chemical shifts and transverse relaxation rates, the major state longitudinal relaxation rate ( $R_{1,A}$ ) and the exchange rate and the minor state population. While fitting data from multiple sites to a global two-state process the exchange rate and minor state population were assumed to be the same for all sites. In all the data analysis the longitudinal relaxation rate was assumed to be the same for both states. Unless mentioned uncertainties in the best fit exchange parameters were estimated using a standard Monte Carlo procedure that consisted of 250 trials (62, 63).

## Results and Discussion

### The choice of optimal $B_1$ fields can depend on the minor state transverse relaxation rate in addition to $k_{ex}$

The A17G S56P FF domain exchanges between the folded state (F) and an alternate conformer (I). The (ILV) methyl  $^1\text{H}$ - $^{13}\text{C}$  correlation map of U- $[\text{H},^{15}\text{N}]$ , Ile $\delta$ 1- $[\text{CH}_3]$ , Leu,Val- $[\text{CH}_3,^{12}\text{CD}_3]$  A17G S56P FF is well resolved at 7.5 °C (Fig. 1a) and a minor state dip is clearly visible in the methyl  $^{13}\text{C}$  CEST profiles (Fig. 1b) from six sites (V30 $\gamma$ 2, I43 $\delta$ 1, I44 $\delta$ 1, L52 $\delta$ 2, L55 $\delta$ 1 & L55 $\delta$ 2). Unlike CPMG experiments where precise exchange parameters are often obtained by a global analysis of data recorded from several sites at multiple  $B_0$  fields (64), precise two and even three-state (slow) exchange parameters can be obtained on a per site basis by analyzing CEST data recorded at a single  $B_0$  field, but with multiple  $B_1$  fields instead, allowing one to identify global exchange processes (25, 32). Since the exchange rate was expected to be approximately  $10\text{ s}^{-1}$ , we initially chose  $B_1$  fields of 1.5 and 3.4 Hz ( $\omega_1$  of 9.4 rad/s and 21.4 rad/s). However, when the  $^{13}\text{C}$  CEST profiles ( $B_1 = 1.5$  & 3.4 Hz) from each of the six sites were analyzed independently to obtain site-specific exchange rates  $k_{ex}$  and minor-state fractional populations,  $p_I$  (Fig. 1c,d), the extracted two-state exchange parameters were poorly defined. This was particularly the case for the exchange rates,  $k_{ex}$ , as shown in Fig. 1c, where the  $k_{ex}$  and  $p_I$  values obtained for each of the six sites from a Monte Carlo

procedure with 250 trials are plotted (grey circles) and in Fig. 1d where the distributions of the site specific  $k_{ex}$  and  $p_I$  values are plotted. The best fit  $k_{ex}$  values range from 6.4 to 51.7 s<sup>-1</sup>, whereas the best fit  $p_I$  values range from 5.4 to 12.9 % across the six residues (Table S2). Analysis of <sup>13</sup>C CEST profiles recorded with  $B_I = 9.8$  Hz in addition to the ones recorded with  $B_I = 1.5$  and 3.4 Hz resulted in more precise  $k_{ex}$  and  $p_I$  values (Fig. 1c, blue pluses; compare  $k_{ex}$  distributions in Fig. 1d and 1e) with site-specific best fit  $k_{ex}$  values now varying from 10.1 to 12.4 s<sup>-1</sup> and best fit  $p_I$  values varying from 8.5 to 10.7 % across the six residues (Table S2). The fact that the analysis of CEST profiles from each of the six sites resulted in very similar exchange parameters (Fig. 1c,e) strongly suggests that they are all reporting on the same global exchange process and a global analysis of the <sup>13</sup>C CEST data ( $B_I = 1.5, 3.4$  & 9.8 Hz) from all six sites resulted in good quality fits ( $\chi^2_{red} \sim 1$ ) and  $k_{ex} = 11.2 \pm 0.5$  s<sup>-1</sup> and  $p_I = 9.5 \pm 0.3$  %. Addition of the  $B_I = 9.8$  Hz <sup>13</sup>C CEST dataset into the analysis procedure leads to a narrower minimum especially for  $k_{ex}$  even in the  $\chi^2_{red}$  v.s. ( $k_{ex}, p_B$ ) plots (Fig. 1f v.s. 1g) obtained from a global analysis of CEST data from all six sites.

For a  $k_{ex}$  value of 11.2 s<sup>-1</sup>,  $B_I$  values of 1.5 and 3.4 Hz correspond to  $\omega_1/k_{ex}$  values of 0.8 and 1.9 respectively and this choice of CEST datasets should have sufficed (14) to obtain precise estimates of the exchange parameters unlike what was observed (Fig. 1c,d). To resolve this conundrum, we noted that the fitted  $R_{2,I}$  values ( $\sim 20$  to  $\sim 70$  s<sup>-1</sup>) are all higher than  $k_{ex}$  ( $\sim 11$  s<sup>-1</sup>) for the six sites (Table S2) and we therefore investigated more generally the effect of the minor state transverse relaxation rate on the size of the minor state dip in CEST profiles.

For a two-state ( $A \rightleftharpoons B$ ) reaction ( $k_{ex} = 15$  s<sup>-1</sup>,  $p_B = 7.5\%$ ) the calculated intensity of the minor ( $B$ ) state dip is plotted in Fig. 2a as a function of  $\omega_1$  for different  $R_{2,B}$  values while the inset shows the minor state dip for various  $R_{2,B}$  values when  $\omega_1 = 15$  rad/s ( $= k_{ex}$ ). It is clear that when  $\omega_1$  is fixed to 15 rad/s, the size of minor state dip decreases as the  $R_{2,B}$  values increase (inset Fig. 2a). For example, the minor dip that is prominent when  $R_{2,B} = 5$  s<sup>-1</sup> (black curve Fig. 2a inset) becomes essentially invisible when  $R_{2,B}$  is increased to 125 s<sup>-1</sup> (cyan curve Fig. 2A inset). A physical explanation is that as  $R_{2,B}$  increases, the  $B_I$  field (analogous to  $B_0$  under free precession) becomes less effective at inducing a relative phase change between the magnetization exchanging between states B and A. When  $R_{2,B} = 5$  s<sup>-1</sup>, the intensity of the minor state dip has a distinctive dependence on  $\omega_1$  as  $\omega_1/k_{ex}$  is varied between  $\sim 0.5$  and  $\sim 2$  (black curve in Fig. 2a). However, for large transverse relaxation rates in the minor state, e.g.  $R_{2,B} = 125$  s<sup>-1</sup>, the size of the minor state dip is small and its intensity changes to a lesser degree when  $\omega_1/k_{ex}$  is varied from 0.5 to 2 (cyan curve in Fig. 2a). Thus,  $R_{2,B}$  influences the size of the minor state dips (65, 66) and when  $R_{2,B}$  is substantially larger than  $k_{ex}$ , CEST datasets recorded with  $\omega_1$  values much larger than  $k_{ex}$  will be required to see the minor state dip clearly and to obtain accurate exchange parameters (Fig. 2a).

# **CEST datasets with $B_1$ fields much larger than $k_{ex}/2\pi$ are required to study exchange when $R_{2,B}$ is comparable to or larger than $k_{ex}$**

To understand the differing shapes of the  $I(\omega_B)/I_0$  v.s.  $\omega_1/k_{ex}$  plots in Fig. 2a, we consider a spin  $1/2$  nucleus undergoing conformational exchange in the slow exchange regime with  $\Delta\omega_{AB} \rightarrow \infty$ . Based on the equivalence between CEST and  $R_{1,\rho}$  experiments (35, 65, 67) the decay of the ground state magnetization under weak  $B_1$  irradiation can be described by  $I(\Delta\omega_{RF,A}) \approx I_0 e^{-R_{1,\rho}(\Delta\omega_{RF,A})T_{EX}}$  (65, 68) with,

$$R_{1,\rho}(\Delta\omega_{RF,A}) = R_{eff}(\Delta\omega_{RF,A}) + R_{ex}(\Delta\omega_{RF,A}) \quad \text{Eq 1}$$

Here  $\Delta\omega_{RF,i}$  is the difference (rad/s) between the offset at which the  $B_1$  irradiation is applied and the resonance frequency of the nucleus in state  $i$ .  $R_{eff}(\Delta\omega_{RF,A})$  is the effective relaxation rate of the spin under  $B_1$  irradiation in the absence of exchange and  $R_{ex}(\Delta\omega_{RF,A})$  is the exchange contribution to relaxation. Different expressions have been obtained for  $R_{ex}(\Delta\omega_{RF,A})$  (67, 69). Focusing on the minor state and assuming that the longitudinal relaxation rate is  $0 \text{ s}^{-1}$  the following simple relation (65, 66) for  $R_{ex}(\Delta\omega_{RF,B})$  applies,

$$R_{ex}(\Delta\omega_{RF,B}) = \frac{R_{ex}^{max}(\omega_1)\Gamma^2(\omega_1)}{\Gamma^2(\omega_1) + \Delta\omega_{RF,B}^2} \quad \text{Eq 2}$$

When the RF-irradiation is applied at the offset of the minor state resonance,  $\Delta\omega_{RF,B} = 0$ ,  $R_{ex} = R_{ex}^{max}(\omega_1)$ , which is the maximum value of  $R_{ex}$  for a given  $B_1$ .  $\Gamma$  is the half width at the half maximum of  $R_{ex}$ , (Fig. S1a).  $R_{ex}^{max}$  and  $\Gamma$  are given by,

$$R_{ex}^{max}(\omega_1) = \frac{k_{ex}p_B\omega_1^2}{\omega_1^2 + K^2} \quad \text{Eq 3}$$

$$\Gamma(\omega_1) = \sqrt{\frac{k_{ex} + R_{2,B}}{k_{ex}}} \sqrt{\omega_1^2 + K^2} \quad \text{Eq 4}$$

$$K = \sqrt{k_{ex}(k_{ex} + R_{2,B})} \quad \text{Eq 5}$$

The size of the minor state dip is given by  $(I_{noex}(T_{EX})/I_0)(1 - e^{-R_{ex}^{max}(\omega_1)T_{EX}})$ , while the shape of the minor state dip ( $I(\Delta\omega_{RF,B})/I_0$  v.s.  $\Delta\omega_{RF,B}$ ) is proportional to  $(1 - e^{-R_{ex}(\Delta\omega_{RF,B})T_{EX}})$  (Fig. S1b). Here  $I_{noex}(T_{EX})$  is the intensity measured at  $\omega_B$  in the absence of exchange, or equivalently at an offset far from  $\omega_A$  and  $\omega_B$  in the presence of exchange, and  $I_{noex}(T_{EX})$  essentially accounts for longitudinal relaxation during  $T_{EX}$ . According to Eq 3 the shape of the  $I(\omega_B)/I_0$  v.s.  $\omega_1/k_{ex}$  plots in Fig. 2a is determined by the



ratio of  $\omega_1$  and  $K$ , rather than the ratio of  $\omega_1$  and  $k_{ex}$ , which means that the curves in Fig. 2a should be identical when  $I(\omega_B)/I_0$  is plotted against  $\omega_1/K$ , as can be seen in Fig. 2b. When  $\omega_1 \ll K$  the minor state dip will not be prominent and its size will increase when  $\omega_1$  is increased (Fig. 2b, Eq 3,  $R_{ex}^{max} \approx k_{ex}p_B\omega_1^2/K^2 = p_B\omega_1^2/(k_{ex} + R_{2,B})$ ) while its width will barely increase when  $\omega_1$  is increased (Eq 4;  $\Gamma \approx [(k_{ex} + R_{2,B})/k_{ex}]^{1/2}K = (k_{ex} + R_{2,B})$ ). Note that  $\Gamma$  is the half width at the half maximum of  $R_{ex}(\Delta\omega_{RF,B})$ , Eq 3, 4 and Fig. S1, whereas the width of the dip,  $D_{width} = \Gamma\sqrt{1+\lambda}/\sqrt{1-\lambda}$ , where  $\lambda = \log\left(\cosh\left(\frac{1}{2}R_{ex}^{max}T_{EX}\right)\right)/\left(\frac{1}{2}R_{ex}^{max}T_{EX}\right)$ , which depends on  $\Gamma$ . Provided that the SNR is adequate to see the minor state dip, analysis of CEST profiles recorded with  $\omega_1 \ll K$  can lead to reasonable estimates of  $p_B$  (compare  $R_{ex}^{max}$  and  $\Gamma$ ) but not  $k_{ex}$  as  $k_{ex}$  cannot be separated from  $R_{2,B}$ . On the other hand, when  $\omega_1 \gg K$  the minor state dip will be prominent, but its size will be independent of  $\omega_1$  (Fig. 2b, Eq 3,  $R_{ex}^{max} \approx k_{ex}p_B$ ) while its width will increase when  $\omega_1$  is increased (Eq 4,  $\Gamma \approx [(k_{ex} + R_{2,B})/k_{ex}]^{1/2}\omega_1$ ). Only the forward rate,  $k_{AB} = k_{ex}p_B$  can be estimated by analyzing of CEST profiles recorded with  $\omega_1 \gg K$ . Thus, recording multiple profiles exclusively with  $\omega_1 \gg K$ , or exclusively with  $\omega_1 \ll K$ , will not provide any additional information and will not meaningfully aid in estimating accurate exchange parameters. The intensity (and width, Eq 4) of the minor state dip shows a distinctive dependence on the value of  $\omega_1$ , when  $\omega_1 \sim K$  (Fig. 2b) making it clear that, in order to derive accurate exchange parameters, CEST datasets should be recorded with  $B_1$  values guided by  $K$ . For example, it follows from the above discussion that accurate  $k_{ex}$  and  $p_B$  values can be obtained from a combined analysis of CEST profiles recorded with  $\omega_1 < K$  and  $\omega_1 > K$  because  $p_B$  can effectively be derived from CEST profiles recorded with  $\omega_1 < K$  and  $k_{ex}p_B$  can be estimated by analyzing CEST profiles recorded with  $\omega_1 > K$ .  $K$  is larger than  $k_{ex}$  and it begins to deviate significantly from  $k_{ex}$  as the value of  $R_{2,B}$  becomes greater than  $k_{ex}$ . The above discussion follows the analysis presented previously (14), except for the fact that the effects of  $R_{2,B}$  have been explicitly retained here and as expected when  $R_{2,B} \ll k_{ex}$ ,  $K \sim k_{ex}$  leading to the previous conclusion that to obtain accurate two-state exchange parameters CEST datasets should be recorded using  $B_1$  values informed by  $k_{ex}$ . For a global process, if the value of  $R_{2,B}$  is constant across the molecule *i.e.* same  $K$  for all sites under investigation then a pair of CEST datasets recorded with  $\omega_1$  values in the (0.5-0.8) $K$  and (1.5-1.8) $K$  ranges will suffice to obtain accurate exchange parameters (14). The dataset with  $\omega_1$  in the (0.5-0.8) $K$  range will have small and unbroadened minor state dips, whereas the dataset with  $\omega_1$  in the (1.5-1.8) $K$  range will have prominent but ( $\omega_1$ ) broadened minor state dips. However,  $R_{2,B}$  may not be constant throughout the molecule, as in the cases studied here, and in such cases it may not be possible to obtain precise exchange parameters from just two CEST datasets. Hence it will be useful to record an additional CEST dataset with a relatively high  $B_1$  so that  $\omega_1/K$  is greater than  $\sim 1.8$  for all sites in the molecule to supplement the datasets recorded with lower  $B_1$  values (guided by  $k_{ex}$ ), where  $\omega_1/K$  samples some part of the 0.5 to  $\sim 1$  region for all residues. A  $B_1$  value for which  $2\pi B_1/k_{ex} \sim 4.5$  may serve as a starting  $B_1$  value for the additional (high



$B_I$ ) dataset as this will result in  $\omega_1/K \sim 1.8$  even when  $R_{2,B}$  is relatively high  $\sim 5k_{ex}$ . If an estimate of  $k_{ex}$  is not available, approximate ranges for  $k_{ex}$  and  $R_{2,B}$  can be estimated by analyzing preliminary CEST data that preferably contains a dataset recorded with a relatively high  $B_I$ , for example 50 Hz.  $K$  calculated from these estimates can then be used to guide the choice of  $B_I$  values to record additional CEST datasets. It should be noted that the minor state  $R_2$  values affect the choice of  $B_I$  fields used in DEST experiments where resolving the minor state dip is not a concern (9, 15). The validity of the analysis presented above has been confirmed using Monte Carlo simulations (See supporting text and figure S2).

The above analysis can be used to rationalize the previous observation, that precise exchange parameters could not be extracted for the A17G S56P FF domain  $F \rightleftharpoons I$  process ( $k_{ex} \sim 11 \text{ s}^{-1}$ ) by analyzing CEST datasets recorded with  $B_I = 1.5$  ( $\omega_1/k_{ex} \sim 0.8$ ) and 3.4 Hz ( $\omega_1/k_{ex} \sim 1.9$ ), but accurate exchange parameters could be extracted upon the inclusion of an additional CEST dataset recorded with a relatively high  $B_I = 9.8$  Hz ( $\omega_1/k_{ex} \sim 5.5$ ) in the least-squares fit procedure. As mentioned earlier the fitted  $R_{2,I}$  values at various sites varied from  $\sim 20$  to  $\sim 70 \text{ s}^{-1}$  all of which are substantially higher than  $k_{ex}$ . For a  $R_{2,I}$  of  $25 \text{ s}^{-1}$ ,  $K \sim 20 \text{ s}^{-1}$ , and consequently  $B_I$  values of 1.5 and 3.4 Hz correspond to  $\omega_1/K$  values of  $\sim 0.5$  and  $\sim 1.1$  respectively that are lower than the desired  $B_I$  values required to obtain accurate exchange parameters. A  $B_I$  value of 9.8 Hz corresponds to  $\omega_1/K$  of 3.1 when  $R_{2,I}$  is  $25 \text{ s}^{-1}$  and  $\sim 2$  when  $R_{2,I}$  is  $70 \text{ s}^{-1}$  and therefore including a dataset recorded with  $B_I = 9.8$  Hz in the analysis procedure provides the desired higher  $B_I$  dataset.

To further test the above strategy, we have used amide  $^{15}\text{N}$  CEST experiments to characterize the folding of the A39G FF domain because the minor state dips in the  $^{15}\text{N}$  CEST profiles are severely broadened due to additional exchange. A39G FF folds from the unfolded state (U) to the native state (F) via two intermediates (I1 and I2) at a rate of  $\sim 70 \text{ s}^{-1}$  (3 °C) with  $p_U \sim 1 \%$ ,  $p_{I1} \sim 0.3 \%$  and  $p_{I2} \sim 0.2 \%$  (24). As U and the folding intermediates I1 and I2 rapidly interconvert among each other on the  $\sim 0.1$  to  $\sim 1$  ms timescale, the folding reaction can be treated as a two-state exchange reaction between the native state (F) and a state U'. U' which is a composite of U, I1 and I2 can be described using a combination of the exchange parameters that are used to describe U, I1 and I2 (24). For example,  $p_{U'} \approx p_U + p_{I1} + p_{I2}$  and  $\varpi_{U'} \approx \varpi_U$  is slightly shifted from  $\varpi_U$  towards  $\varpi_{I1}$  because U and I1 are in fast exchange (24). Exchange between U, I1 and I2 severely broadens several U' dips and explicit dips arising from the I1 state are not visible in any of the CEST profiles, whereas the CEST profile of only Ser 56 (that is excluded from the present analysis) contains an explicit dip due to the I2 state (24). The amide  $^{15}\text{N}$ - $^1\text{H}$  correlation map is well resolved (Fig. 3a) and  $^{15}\text{N}$  CEST profiles were obtained (Fig. 3b) for 58 out of 60 (non-proline) ordered (residue 10 to 71) amino acid sites in the molecule. In the discussion that follows we only consider 19 sites with large chemical shift differences ( $|\Delta\varpi_{FU'}| > 3 \text{ ppm}$ ). A global two-state exchange model satisfied ( $\chi^2_{red} \sim 1$ ) the ( $B_I = 6.0$  &  $18.4 \text{ Hz}$ )  $^{15}\text{N}$  CEST data resulting in well-defined exchange parameters with  $k_{ex} = 72 \pm 3 \text{ s}^{-1}$  and  $p_{U'} =$

$1.39 \pm 0.03$  % and the CEST derived  $\Delta\omega_{FU'}$  values are in good agreement with the predicted  $\Delta\omega_{FU}$  values (RMSD 1.8 ppm, Fig. 3c) confirming that the U state is the dominant state among the states that comprise U'. As mentioned above, the exchange between U, I1 and I2 results in some U' dips that are severely broadened, as can be seen in Fig. 3b where the minor state dip of I43 is significantly broader than that of L52 and Q68. The distribution of  $R_{2,U'}$  values obtained from the global two-state analysis of ( $B_I = 6.0$  & 18.4 Hz) amide  $^{15}\text{N}$  CEST profiles is plotted in Fig. 3d (Table S3). The  $R_{2,U'}$  values show a broad distribution with several residues having  $R_{2,U'}$  values above  $2k_{ex}$  ( $\sim 140 \text{ s}^{-1}$ ). For the four residues (L52, L55, K66 & Q68; Table S3) with  $R_{2,U'}$  values less than  $50 \text{ s}^{-1}$  ( $K = 94 \text{ s}^{-1}$  when  $R_{2,U'} = 50 \text{ s}^{-1}$ ) single residue fits of  $^{15}\text{N}$  CEST data ( $B_I = 6.0$  & 18.4 Hz) yielded well defined exchange parameters (Fig. 3e) with  $k_{ex}$  varying from 53 to  $73 \text{ s}^{-1}$  and  $p_{U'}$  varying from 1.4 to 1.7 % across the four different residues (Table S3). Including an additional CEST dataset recorded with  $B_I = 46$  Hz in the fitting procedure only has a small effect on the exchange parameters extracted for these residues (Fig. 3f) with  $k_{ex}$  now varying from 57 to  $64 \text{ s}^{-1}$  and  $p_{U'}$  now varying from 1.4 to 1.6 % across the different residues. In contrast, for the seven residues (T13, K28, R29, M42, I43, I44 & N45; Table S3) with  $R_{2,U'}$  values greater than  $140 \text{ s}^{-1}$  ( $\sim 2k_{ex}$ ), analysis of the  $B_I = 6.0$  and 18.4 Hz  $^{15}\text{N}$  CEST datasets, on a per residue basis, resulted in poorly defined exchange parameters (Fig. 3g) with  $k_{ex}$  varying from 40 to  $146 \text{ s}^{-1}$  and  $p_{U'}$  varying from 1.1 to 1.5 %. This is not surprising as  $K \sim 140 \text{ s}^{-1}$  when  $R_{2,U'} = 200 \text{ s}^{-1}$  resulting in relatively small  $\omega_1/K$  values of 0.27 and 0.82 for  $B_I$  fields of 6.0 and 18.4 Hz, respectively. For these seven residues with large  $R_{2,U'}$  values, significantly more precise exchange parameters were obtained when the CEST data recorded with a  $B_I$  of 46 Hz was also included in the analysis (Fig. 3h), with  $k_{ex}$  now varying from 57 to  $90 \text{ s}^{-1}$  and  $p_{U'}$  now varying from 1.2 to 1.5%. A  $B_I$  field of 46 Hz ( $\omega_1/k_{ex} \sim 4$ ) corresponds to a  $\omega_1/K$  value of 2.1 when  $R_{2,U'} = 200 \text{ s}^{-1}$  and thus including this higher field results in more precise exchange rates when the minor state dips are severely broadened, which is consistent with the theoretical analysis presented above. In a previous study of the A39G FF folding using  $^{15}\text{N}$  CEST experiments precise exchange parameters were obtained because datasets with high  $B_I$  values were inadvertently recorded, while looking for the minor state dips (25).

## Concluding remarks

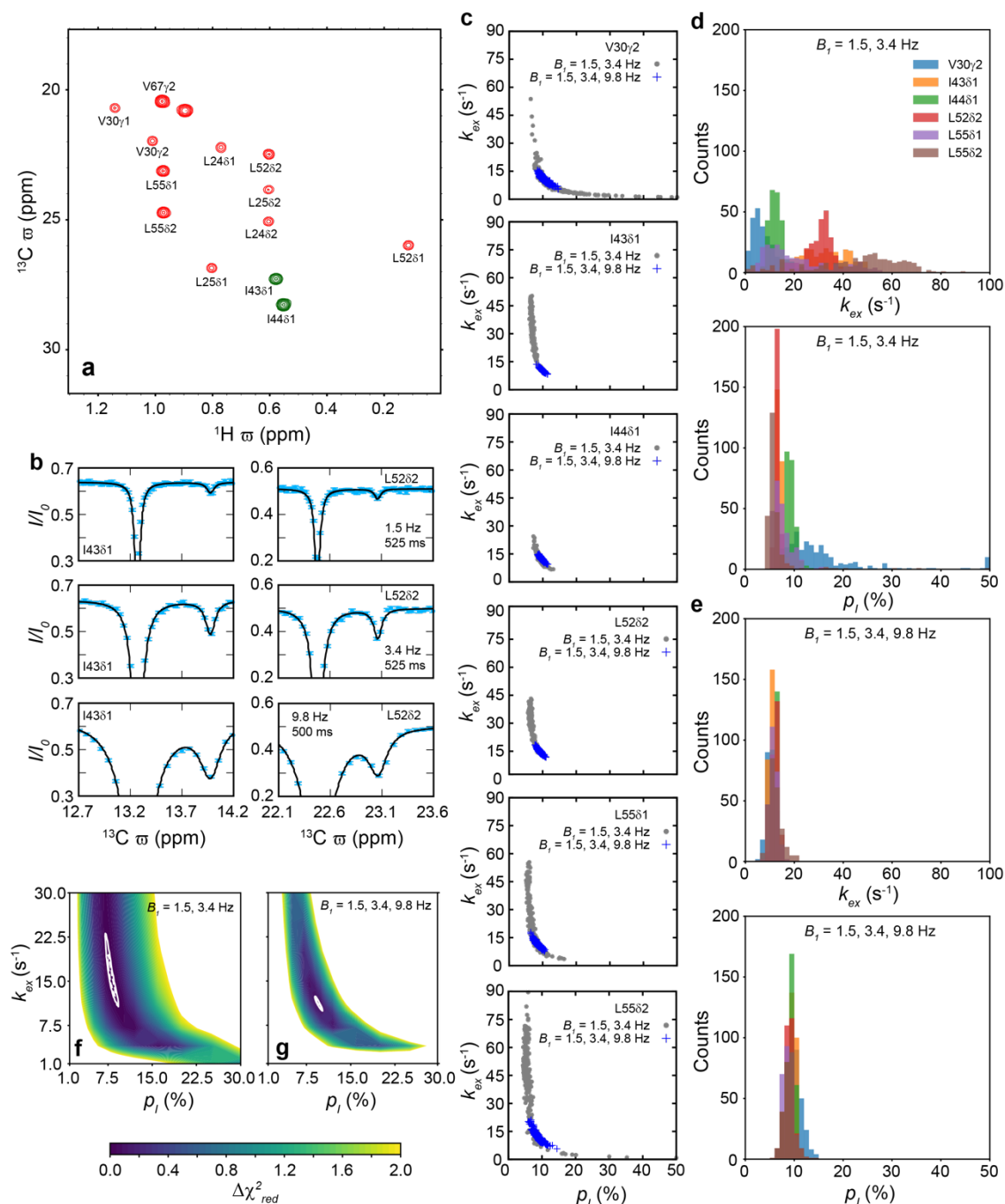
We have shown that the choice of  $B_I$  fields required to characterize chemical exchange using CEST experiments depends on the (apparent) minor state transverse relaxation rate in addition to  $k_{ex}$ . We suggest that the choice of  $B_I$  fields should be governed by  $K (= [k_{ex}(k_{ex} + R_{2,B})]^{1/2})$  as opposed to  $k_{ex}$ . When  $R_{2,B} \ll k_{ex}$ ,  $K \approx k_{ex}$  and the choice of  $B_I$  fields to characterize exchange will be essentially determined by  $k_{ex}$ . However, when  $R_{2,B}$  is substantially greater than  $k_{ex}$ , CEST datasets recorded with higher  $B_I$  fields determined by  $K$ , as opposed to  $k_{ex}$ , are required to obtain accurate exchange parameters. Often this will necessitate recording an additional CEST dataset with a relatively high  $B_I$  value (recommended to be  $\sim 4.5k_{ex}/2\pi$ ) so that  $2\pi B_I/K$  is greater than  $\sim 1.8$  for all the sites in the molecule. Although this strategy

often requires recording additional CEST datasets with higher  $B_1$  values, it should be noted that these datasets can be recorded rapidly compared to datasets with lower  $B_1$  values as the spacing between adjacent offsets, at which  $B_1$  irradiation is carried out, is larger when the  $B_1$  values are higher (14, 70). We expect that the conclusions presented here will be valuable when CEST experiments are used to characterize slow processes ( $k_{ex} \leq \sim 25 \text{ s}^{-1}$ ) in large proteins, processes with  $k_{ex} \leq \sim 10 \text{ s}^{-1}$  in small to medium sized proteins and when the minor state dips are severely exchange broadened due to the presence of other sparsely populated states, as in the case of A39G FF studied here. These results will continue to become more relevant as higher field spectrometers become available because the transverse relaxation rates for several sites in protein molecules will increase with field strength.

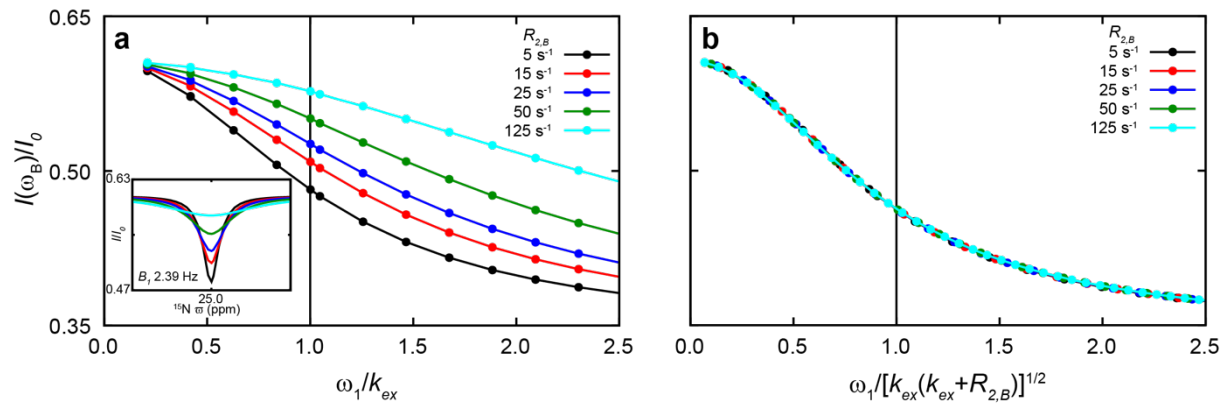
**Supporting Information** Supporting Information is included in this file after the references.

## Acknowledgements

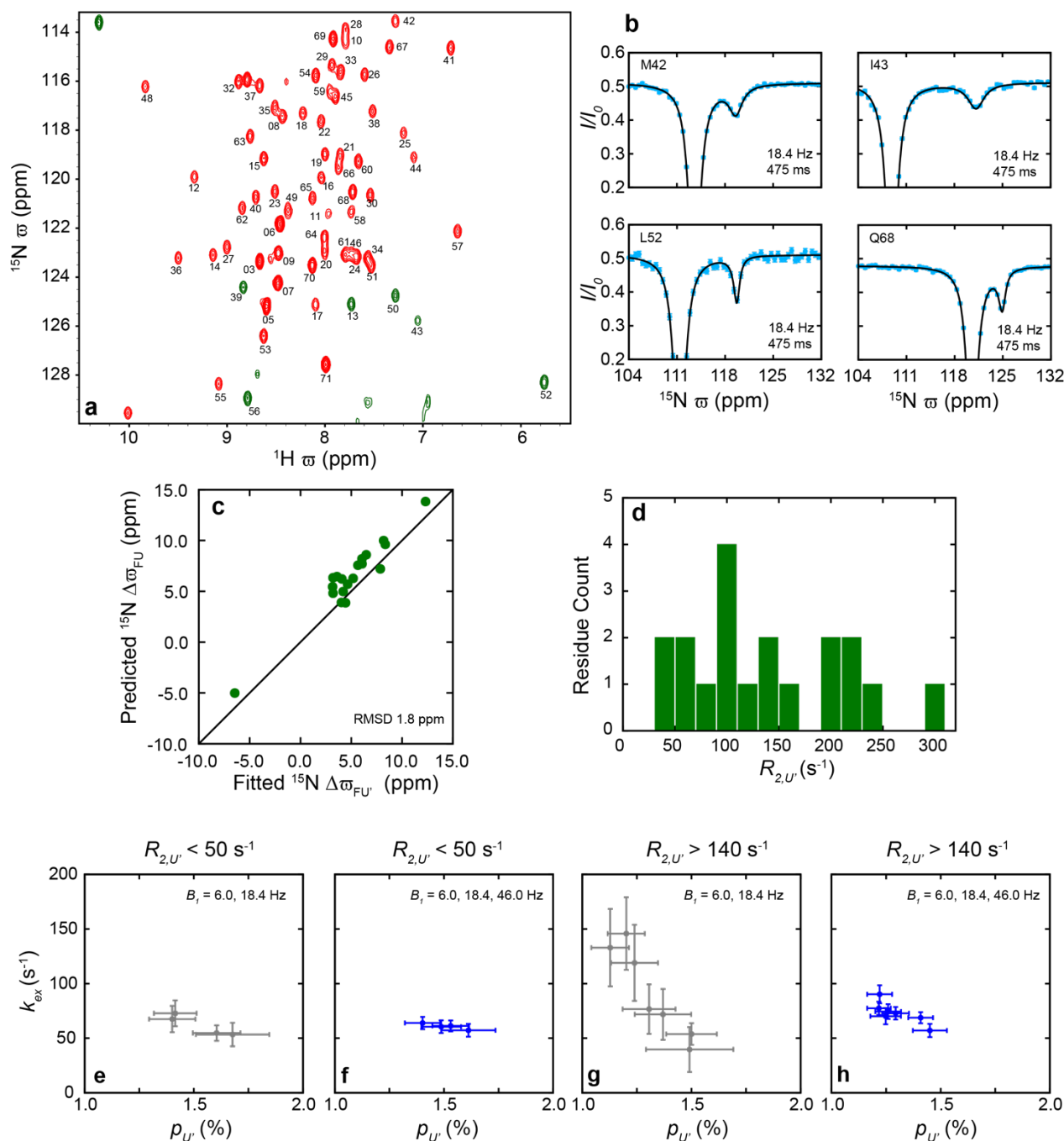
The authors thank Nemika Thapliyal (TIFR Hyderabad) and Dr Debajyoti De (TIFR Hyderabad) for providing the A17G S56P FF sample, Dr. G. Bouvignies (ENS, Paris) for providing the program *ChemEx*, Prof. Lewis Kay (University of Toronto) for providing the CEST and D-CEST pulse sequences, TIFR Hyderabad NMR facility for the generous grant of spectrometer time and Dr Krishna Rao for maintaining the facility. PV acknowledges intramural funding from TIFR Hyderabad (DAE, Government of India, Project No. RTI 4007). DFH is supported by the UKRI and EPSRC (EP/X036782/1). For the purpose of open access, the author has applied a Creative Commons Attribution (CC BY) license to any Author Accepted Manuscript version arising.



**Fig. 1** High  $B_1$  fields lead to precise exchange parameters for the A17G S56P FF  $F \rightleftharpoons I$  reaction ( $k_{\text{ex}} \sim 11.2$  s $^{-1}$ ). (a) Methyl  $^{13}\text{C}$ - $^1\text{H}$  correlation map of the U-[ $^2\text{H}$ ,  $^{15}\text{N}$ ], Ile $\delta$ 1-[ $^{13}\text{CH}_3$ ], Leu,Val-[ $^{13}\text{CH}_3$ ,  $^{12}\text{CD}_3$ ] A17G S56P FF (16.4 T, 7.5  $^\circ\text{C}$ ). Peaks are labelled according to the site from which they arise. Green peaks are aliased in the  $^{13}\text{C}$  dimension. (b) Representative methyl  $^{13}\text{C}$  CEST profiles ( $B_1$  and  $T_{\text{EX}}$  indicated) clearly show the presence of a minor state dip. Cyan circles are used to represent the experimental data while the black line is drawn using the global best fit parameters ( $k_{\text{ex}} = 11.2$  s $^{-1}$ ,  $p_I = 9.52\%$ ; Table S2). (c) Scatter plots showing the distribution of  $k_{\text{ex}}$  and  $p_I$  values obtained using a Monte Carlo procedure with 250 trials. Analysis was carried out separately at each site using two different combinations of CEST datasets:  $B_1$  values of 1.5 and 3.4 Hz (grey circles) and  $B_1$  values of 1.5, 3.4 and 9.8 Hz (blue pluses). (d,e) Histograms showing the distribution of site specific  $k_{\text{ex}}$  and  $p_I$  values from (c).  $\Delta\chi^2_{\text{red}}$  vs ( $k_{\text{ex}}$ ,  $p_I$ ) plots obtained from a global analysis of the methyl  $^{13}\text{C}$  CEST 1.5 and 3.4 Hz (f) and 1.5, 3.4 and 9.8 Hz (g) datasets.  $\Delta\chi^2_{\text{red}}$  is difference between  $\chi^2_{\text{red}}$  and the minimum (best fit) value of  $\chi^2_{\text{red}}$  (lowest value of  $\Delta\chi^2_{\text{red}}$  is 0).  $\Delta\chi^2_{\text{red}}$  values above 2 are in white. In f and g contours corresponding to the 68 and 95% confidence intervals of  $k_{\text{ex}}$  and  $p_I$  based on 10,000 Monte Carlo trials are also shown using dashed and solid white lines respectively.



**Fig. 2** The size of the minor (B) state dip depends on  $R_{2,B}$ . Plot of the normalized minor state dip intensity ( $I(\omega_B)/I_0$ ) as a function of  $\omega_1/k_{ex}$  (a) and  $\omega_1/[k_{ex}(k_{ex} + R_{2,B})]^{1/2}$  (b) for different  $R_{2,B}$  values. The inset in (a) shows the CEST profile ( $B_1 = k_{ex}/2\pi = 2.39$  Hz) around  $\omega_B$  for different  $R_{2,B}$  values. Calculations were performed for a two-state slow exchange reaction ( $k_{ex}/\Delta\omega_{AB} \sim 0$ ) with  $k_{ex} = 15$   $s^{-1}$ ,  $p_B = 7.5$  %,  $R_{1,A} = 1$   $s^{-1}$ ,  $R_{1,B} = 1$   $s^{-1}$ ,  $R_{2,A} = 5$   $s^{-1}$ ,  $\omega_A = 0$  ppm,  $\omega_B = 25$  ppm ( $^{15}N$ , 16.4 T) and  $T_{EX} = 0.5$  s.



**Fig. 3** Folding of the A39G FF domain studied using  $^{15}\text{N}$  CEST experiments. (a) The amide  $^{15}\text{N}$ - $^1\text{H}$  correlation map of U- $^{15}\text{N}$  A39G FF (11.7 T, 2.5  $^{\circ}\text{C}$ ) in which peaks are labelled according to the residue from which they arise. Peaks aliased in the  $^{15}\text{N}$  dimension are shown in green. (b) Representative amide  $^{15}\text{N}$  CEST profiles ( $B_1$  and  $T_{EX}$  indicated) from four different sites in the molecule. Cyan circles represent the experimental data and the black line is drawn according to global best fit parameters ( $k_{ex} = 71.6$  s $^{-1}$ ,  $p_{U'} = 1.39\%$ ; Table S3). (c) Correlation between the predicted  $\Delta\omega_{FU}$  and CEST derived  $\Delta\omega_{FU'}$  shifts.  $\omega_U$  shifts were predicted using the program POTENCI (71). (d) Distribution of the  $R_{2,U'}$  values obtained from a global analysis 6.0 and 18.4 Hz  $^{15}\text{N}$  CEST data. (e) For the four residues with  $R_{2,U'} < 50$  s $^{-1}$  very similar residue specific  $k_{ex}$  and  $p_{U'}$  values are obtained from the analysis of 6.0 and 18.4 Hz  $^{15}\text{N}$  CEST data and the inclusion of 46.0 Hz CEST data does not really have an effect on the distribution of the  $k_{ex}$  and  $p_{U'}$  values (f). (g) For the seven residues with  $R_{2,U'} > 140$  s $^{-1}$  there is a large variation in the residue specific  $k_{ex}$  values obtained from the analysis 6.0 and 18.4 Hz  $^{15}\text{N}$  CEST data and the inclusion of 46.0 Hz CEST data leads to a narrower distribution of  $k_{ex}$  and  $p_{U'}$  values (h).



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## **Supporting Information**

**Increasing the accuracy of exchange parameters reporting on slow dynamics by performing CEST experiments with high  $B_1$  fields**



## Supporting Text

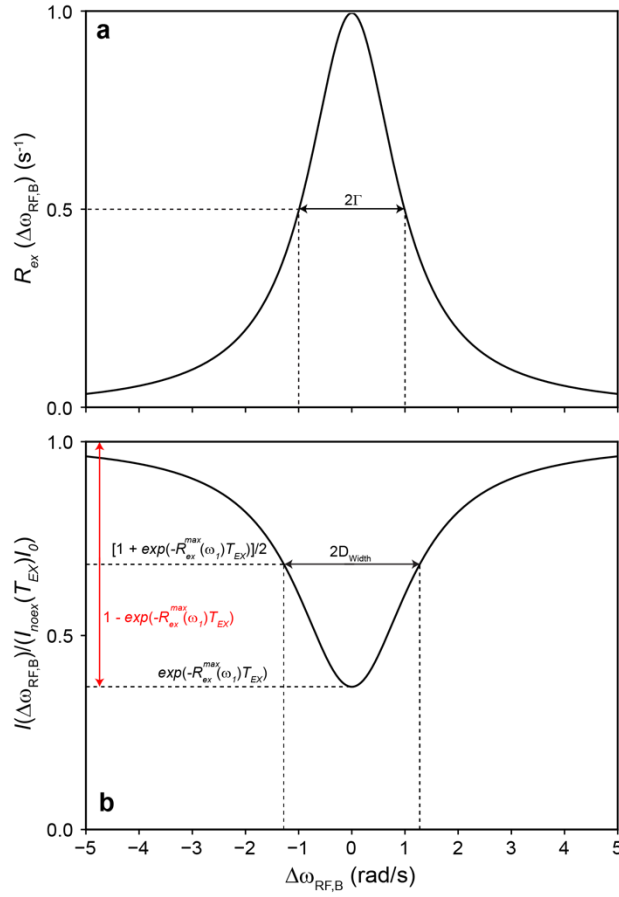
**Calculations confirm that CEST datasets with ‘high’  $B_1$  fields are necessary to obtain precise exchange parameters for slow processes when  $R_{2,B}$  is larger than  $k_{ex}$**

Monte Carlo simulations (1, 2) were used to test the validity of the theoretical analysis presented in the text. CEST profiles ( $^{15}\text{N}$ ; 16.4 T) with  $B_1$  values of 1.7, 4.1, 9.9 and 11.7 Hz were generated for two ‘residues’ with  $k_{ex} = 15 \text{ s}^{-1}$ ,  $p_B = 7.5 \%$  and  $\Delta\omega_{AB} = 5 \text{ ppm}$ . When  $k_{ex}$  is  $15 \text{ s}^{-1}$ ,  $B_1$  values of 1.7, 4.1, 9.9 and 11.7 Hz correspond to  $\omega_1/k_{ex}$  values of 0.7, 1.7, 4.1 and 4.9 respectively.

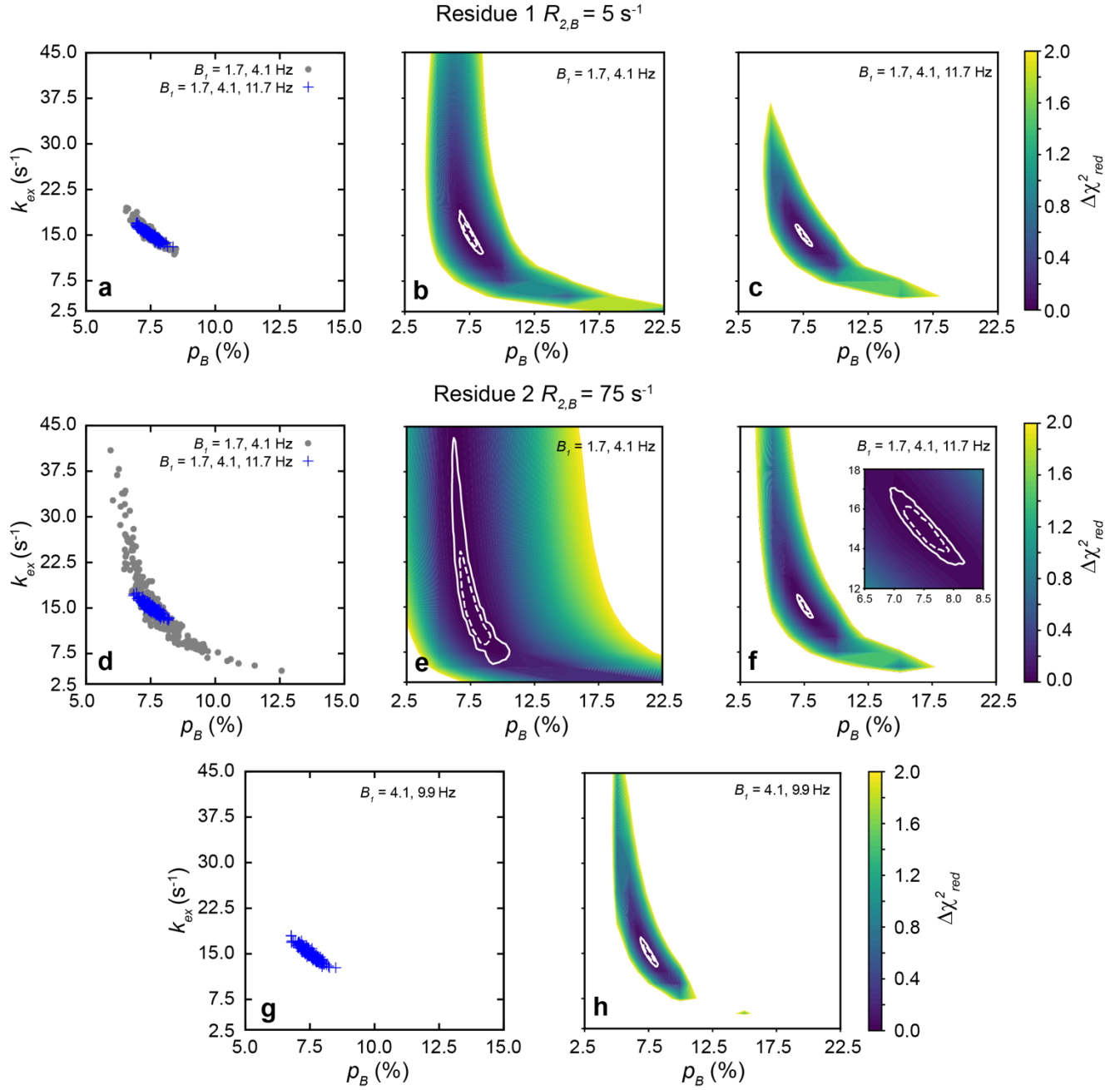
For residue 1  $R_{2,B}$  was set to  $5 \text{ s}^{-1}$  resulting in  $K (= [k_{ex}(k_{ex} + R_{2,B})]^{1/2}) = 17.3 \text{ s}^{-1}$  that is similar to the  $k_{ex}$  value of  $15 \text{ s}^{-1}$ . Fits to the 1.7 and 4.1 Hz CEST profiles results in well-defined exchange parameters ( $k_{ex} = 15 \pm 1.4 \text{ s}^{-1}$ ,  $p_B = 7.5 \pm 0.4 \%$ ; grey circles in Fig. S2a) and a distinct minimum in the  $\chi_{red}^2$  vs  $(k_{ex}, p_B)$  plot (Fig. S2b) because  $B_1$  values of 1.7 and 4.1 Hz correspond to  $\omega_1/K$  values of 0.6 and 1.5 respectively. Including the CEST profile calculated with  $B_1 = 11.7 \text{ Hz}$  only has a small effect on the extracted exchange parameters ( $k_{ex} = 15 \pm 0.7 \text{ s}^{-1}$ ,  $p_B = 7.5 \pm 0.3 \%$ ; blue pluses in Fig. S2a; Fig. S2c).

For residue 2 on the other hand,  $R_{2,B}$  was set to  $75 \text{ s}^{-1}$  resulting in  $K = 36.7 \text{ s}^{-1}$  that is more than twice  $k_{ex}$  and fits to the 1.7 and 4.1 Hz CEST profiles results in poorly defined exchange parameters ( $k_{ex} = 15 \pm 7 \text{ s}^{-1}$ ,  $p_B = 7.5 \pm 1.0 \%$ ; grey circles in Fig. S2d) and a  $\chi_{red}^2$  vs  $(k_{ex}, p_B)$  plot without a sharp minimum (especially along  $k_{ex}$ ) (Fig. S2e) because 1.7 and 4.1 Hz correspond to  $\omega_1/K$  values of 0.3 and 0.7 that are too small for the extraction of accurate exchange parameters. Including the CEST dataset calculated with  $B_1 = 11.7 \text{ Hz}$  in the analysis procedure leads to more precise exchange parameters ( $k_{ex} = 15 \pm 0.8 \text{ s}^{-1}$ ,  $p_B = 7.5 \pm 0.3 \%$ ; blue pluses in Fig. S2d) and a distinct minimum in the  $\chi_{red}^2$  vs  $(k_{ex}, p_B)$  plot (Fig. S2f) because  $B_1 = 11.7 \text{ Hz}$  corresponds to a  $\omega_1/K$  value of 2 for residue 2 which nicely complements the  $B_1 = 1.7$  and 4.1 Hz datasets that correspond to  $\omega_1/K$  values of 0.3 and 0.7 respectively. Finally precise exchange parameters ( $k_{ex} = 15 \pm 1 \text{ s}^{-1}$ ,  $p_B = 7.5 \pm 0.3 \%$ ; Fig. S2g,h) were also obtained by analyzing the 4.1 & 9.9 Hz CEST datasets that correspond to the ‘recommended’  $\omega_1/K$  values of 0.7 and 1.7 (3).

In the above analysis the CEST profiles were generated with no errors but an uncertainty of 0.5 % in the normalized intensities was assumed to carry out the Monte Carlo analysis.



**Fig. S1** Schematic illustration of (a)  $R_{ex}(\Delta\omega_{RF,B})$  ( $R_{ex}^{max}(\omega_1)\Gamma^2(\omega_1)/(\Gamma^2(\omega_1) + \Delta\omega_{RF,B}^2)$ ) v.s  $\Delta\omega_{RF,B}$  and (b)  $I(\Delta\omega_{RF,B})/(I_{noex}(T_{EX})I_0)$  ( $e^{-R_{ex}^{max}(\omega_1)T_{EX}}$ ) v.s  $\Delta\omega_{RF,B}$ . According to equations 1-5 of the text (3-6), the size of the minor state dip is proportional to  $(1 - e^{-R_{ex}^{max}(\omega_1)T_{EX}})$  (red arrow in b) and its shape ( $I(\Delta\omega_{RF,B})/I_0$  v.s  $\Delta\omega_{RF,B}$ ) is proportional to  $(1 - e^{-R_{ex}(\Delta\omega_{RF,B})T_{EX}})$ . The “width” of the dip is  $2D_{width}$  and  $D_{width}$  is the value of  $\Delta\omega_{RF,B}$  at which  $(1 - e^{-R_{ex}(\Delta\omega_{RF,B})T_{EX}}) = (1 - e^{-R_{ex}^{max}(\omega_1)T_{EX}})/2$ . According to equation 2 in the text  $R_{ex}(\Delta\omega_{RF,B}) = R_{ex}^{max}(\omega_1)\Gamma^2(\omega_1)/(\Gamma^2(\omega_1) + \Delta\omega_{RF,B}^2)$  leading to  $D_{width} = \Gamma\sqrt{\frac{1+\lambda}{1-\lambda}}$ , where  $\lambda = \frac{\log(\cosh(\frac{1}{2}R_{ex}^{max}T_{EX}))}{(\frac{1}{2}R_{ex}^{max}T_{EX})}$ . Hence  $D_{width}$  depends on  $\Gamma$  but is not  $\Gamma$  as  $\Gamma$  is the half width at half maximum of  $R_{ex}(\Delta\omega_{RF,B})$ . The plots were made with  $R_{ex}^{max}(\omega_1)$ ,  $\Gamma(\omega_1)$ ,  $T_{EX}$  all set to 1 in their respective units.



**Fig. S2** Simulations confirm that CEST datasets with ‘high’  $B_1$  fields are necessary to obtain precise exchange parameters when  $R_{2,B}$  is high compared to  $k_{ex}$ . Scatter plots of single residue exchange parameters obtained from a Monte Carlo procedure involving 250 trials carried out using calculated  $^{15}\text{N}$  CEST profiles ( $^{15}\text{N}$ ; 16.4 T;  $T_{EX} = 500$  ms) that were generated for two “residues” with  $k_{ex} = 15 \text{ s}^{-1}$ ,  $p_B = 7.5 \%$ ,  $R_{1,A} = R_{1,B} = 1 \text{ s}^{-1}$ ,  $R_{2,A} = 5 \text{ s}^{-1}$ ,  $\varpi_A = 0 \text{ ppm}$ ,  $\varpi_B = 5 \text{ ppm}$ ,  $R_{2,B} = 5 \text{ s}^{-1}$  for residue 1 (a) and  $R_{2,B} = 75 \text{ s}^{-1}$  for residue 2 (d,g).  $\chi^2_{red}$  vs  $(k_{ex}, p_B)$  plots calculated for residue 1 (b,c) and residue 2 (e,f,h) by analyzing CEST datasets calculated using the indicated  $B_1$  values. It is clear that when  $R_{2,B}$  is high compared to  $k_{ex}$  (residue 2, panels d-h) that CEST datasets recorded with higher  $B_1$  values (9.9 or 11.7 Hz) are crucial for obtaining precise exchange parameters. In b,c,e,f and h  $\chi^2_{red}$  values above 2 are in white and contours corresponding to the 68 and 95% confidence intervals of  $k_{ex}$  and  $p_B$  based on 10,000 Monte Carlo trials are shown using dashed and solid white lines respectively. Here  $\Delta\chi^2_{red} = \chi^2_{red}$  because the best fit  $\chi^2_{red} = 0$  as the CEST profiles were generated with no errors.

Sample	Protein	Buffer	Temp (°C)	<sup>13</sup> C D-CEST NMR Experiments						Comments
1	A17G S56P FF	50 mM sodium acetate, 100 mM NaCl, 2 mM NaN <sub>3</sub> , 2 mM EDTA, 30% d7-glucose, 100 % D <sub>2</sub> O (pH 5.7)	7.5	<i>B</i> <sub>0</sub> (T)	<i>B</i> <sub>1</sub> (Hz)	<i>T</i> <sub>EX</sub> (ms)	Centre [ω <sub>cent</sub> ] (ppm)	SW (Hz)	Step Size (Hz)	4 mM U- [ <sup>2</sup> H, <sup>15</sup> N], Ileδ1- [ <sup>13</sup> CH <sub>3</sub> ], Leu,Val- [ <sup>13</sup> CH <sub>3</sub> , <sup>12</sup> CD <sub>3</sub> ] A17G S56P FF
				16.4	1.5	525	19.662	350	3.5	
				16.4	3.4	525	19.662	498	6.0	
				16.4	9.8	500	19.662	500	12.5	
				<sup>15</sup> N CEST NMR Experiments						
2	A39G FF	50 mM sodium acetate, 100 mM NaCl, 2 mM NaN <sub>3</sub> , 2 mM EDTA, 10 % D <sub>2</sub> O (pH 5.7)	2.5	<i>B</i> <sub>0</sub> (T)	<i>B</i> <sub>1</sub> (Hz)	<i>T</i> <sub>EX</sub> (ms)	Centre [ω <sub>cent</sub> ] (ppm)	Range (Hz)	Step Size (Hz)	4 mM U-[ <sup>15</sup> N] A39G FF
				11.7	6.0	525	117.814	± 616	11.0	
				11.7	18.4	475	117.814	± 700	25.0	
				11.7	46.0	475	117.814	± 1000	50.0	

**Table S1** Details of the methyl <sup>13</sup>C D-CEST (7) and the amide <sup>15</sup>N CEST (8) experiments carried out in this study. SW is the sweep width of D-CEST sequence in the CEST dimension.  $B_1$  irradiation was carried out at offsets varying from -SW/2 (-Range) to +SW/2 (+Range) around  $\omega_{cent}$  in steps of 'Step Size'.

Residue	$B_{1s}$ (Hz)	Global Analysis								Residue-Specific Analysis							
		$k_{ex}$ (s <sup>-1</sup> )	$p_I$ (%)	$\Delta\omega$ (ppm)	$R_{1,F}$ (s <sup>-1</sup> )	$R_{2,F}$ (s <sup>-1</sup> )	$R_{2,I}$ (s <sup>-1</sup> )	$\chi^2_{red}$	K (s <sup>-1</sup> )	$k_{ex}$ (s <sup>-1</sup> )	$p_I$ (%)	$\Delta\omega$ (ppm)	$R_{1,F}$ (s <sup>-1</sup> )	$R_{2,F}$ (s <sup>-1</sup> )	$R_{2,I}$ (s <sup>-1</sup> )	$\chi^2_{red}$	K (s <sup>-1</sup> )
V30 <sub>γ</sub> 2	1.5, 3.4 Hz	15.9 ± 2.3	8.00 ± 0.42	-0.30 ± 0.01	1.35 ± 0.01	17.2 ± 0.4	26.8 ± 2.9	0.93	26	6.4 ± 6.8	12.86 ± 9.45	-0.30 ± 0.01	1.35 ± 0.01	18.1 ± 1.1	34.5 ± 7.2	1.01	16
I43δ1				0.71 ± 0.01	0.86 ± 0.01	15.5 ± 0.4	24.6 ± 2.2		25	32.5 ± 9.3	6.82 ± 0.54	0.71 ± 0.01	0.86 ± 0.01	14.3 ± 0.7	13.6 ± 8.0	0.96	39
I44δ1				-0.29 ± 0.01	0.72 ± 0.01	11.8 ± 0.5	15.0 ± 2.9		22	12.1 ± 3.0	9.21 ± 0.95	-0.28 ± 0.01	0.72 ± 0.01	12.1 ± 0.5	18.0 ± 2.3	0.51	19
L52δ2				0.58 ± 0.01	1.30 ± 0.01	16.5 ± 0.7	14.1 ± 4.6		22	33.0 ± 5.7	6.25 ± 0.51	0.59 ± 0.01	1.31 ± 0.01	15.2 ± 0.6	0.0 ± 4.9	1.59	33
L55δ1				0.86 ± 0.01	1.86 ± 0.01	13.6 ± 0.9	31.2 ± 9.5		27	17.3 ± 13.7	7.06 ± 1.92	0.86 ± 0.01	1.86 ± 0.01	13.3 ± 1.1	25.3 ± 10.7	0.77	27
L55δ2				0.50 ± 0.01	2.41 ± 0.01	16.1 ± 0.8	50.3 ± 5.5		32	51.7 ± 18.9	5.37 ± 4.19	0.51 ± 0.01	2.42 ± 0.01	13.9 ± 1.5	7.8 ± 16.4	0.64	55
Residue	$B_{1s}$ (Hz)	Global Analysis								Residue-Specific Analysis							
		$k_{ex}$ (s <sup>-1</sup> )	$p_I$ (%)	$\Delta\omega$ (ppm)	$R_{1,F}$ (s <sup>-1</sup> )	$R_{2,F}$ (s <sup>-1</sup> )	$R_{2,I}$ (s <sup>-1</sup> )	$\chi^2_{red}$	K (s <sup>-1</sup> )	$k_{ex}$ (s <sup>-1</sup> )	$p_I$ (%)	$\Delta\omega$ (ppm)	$R_{1,F}$ (s <sup>-1</sup> )	$R_{2,F}$ (s <sup>-1</sup> )	$R_{2,I}$ (s <sup>-1</sup> )	$\chi^2_{red}$	K (s <sup>-1</sup> )
V30 <sub>γ</sub> 2	1.5, 3.4, 9.8 Hz	11.2 ± 0.5	9.52 ± 0.27	-0.31 ± 0.01	1.36 ± 0.01	19.3 ± 0.3	35.5 ± 1.9	1.00	23	10.1 ± 1.9	10.68 ± 1.21	-0.31 ± 0.01	1.36 ± 0.01	19.2 ± 0.6	36.7 ± 4.9	1.08	22
I43δ1				0.71 ± 0.01	0.87 ± 0.01	17.1 ± 0.3	31.6 ± 1.6		22	10.4 ± 0.9	9.80 ± 0.52	0.71 ± 0.01	0.87 ± 0.01	17.2 ± 0.3	32.0 ± 2.0	0.97	21
I44δ1				-0.28 ± 0.01	0.72 ± 0.01	13.2 ± 0.3	22.0 ± 2.1		19	12.3 ± 1.0	9.64 ± 0.50	-0.28 ± 0.01	0.72 ± 0.01	12.8 ± 0.3	22.5 ± 1.6	0.61	21
L52δ2				0.58 ± 0.01	1.31 ± 0.01	18.0 ± 0.4	23.3 ± 3.0		20	12.4 ± 1.3	9.05 ± 0.60	0.59 ± 0.01	1.31 ± 0.01	17.8 ± 0.4	22.7 ± 2.1	1.66	21
L55δ1				0.86 ± 0.01	1.86 ± 0.01	15.3 ± 0.6	32.8 ± 7.9		22	11.5 ± 1.7	8.48 ± 0.85	0.86 ± 0.01	1.87 ± 0.01	15.1 ± 0.5	33.5 ± 3.7	0.90	23
L55δ2				0.50 ± 0.01	2.41 ± 0.01	16.6 ± 0.5	66.7 ± 4.4		30	11.7 ± 2.7	8.87 ± 1.18	0.50 ± 0.01	2.41 ± 0.01	16.7 ± 0.6	64.7 ± 8.0	0.69	30

**Table S2** Best fit exchange parameters obtained by analyzing two different sets of methyl <sup>13</sup>C D-CEST profiles recorded using the 4 mM U-[<sup>2</sup>H,<sup>15</sup>N], Ileδ1-[<sup>13</sup>CH<sub>3</sub>], Leu,Val-[<sup>13</sup>CH<sub>3</sub>,<sup>12</sup>CD<sub>3</sub>] A17G S56P FF sample at 7.5 °C (16.4 T).

Residue	$B_{1s}$ (Hz)	Global Analysis								Residue-Specific Analysis										
		$k_{ex}$ (s <sup>-1</sup> )	$p_{U'}$ (%)	$\Delta\varpi$ (ppm)	$R_{1,F}$ (s <sup>-1</sup> )	$R_{2,F}$ (s <sup>-1</sup> )	$R_{2,U'}$ (s <sup>-1</sup> )	$\chi^2_{red}$	$K$ (s <sup>-1</sup> )	$k_{ex}$ (s <sup>-1</sup> )	$p_{U'}$ (%)	$\Delta\varpi$ (ppm)	$R_{1,F}$ (s <sup>-1</sup> )	$R_{2,F}$ (s <sup>-1</sup> )	$R_{2,U'}$ (s <sup>-1</sup> )	$\chi^2_{red}$	$K$ (s <sup>-1</sup> )			
T13	6.0, 18.4	71.6 ± 3.1	1.39 ± 0.03	4.07 ± 0.03	1.23 ± 0.01	15.2 ± 0.2	143.6 ± 13.0	0.98	124	53.8 ± 10.6	1.50 ± 0.12	4.05 ± 0.03	1.23 ± 0.01	15.4 ± 0.2	143.8 ± 12.3	0.88	103			
K22																				
K26																				
K28																				
R29																				
N33																				
E37																				
K41																				
M42																				
I43																				
I44																				
N45																				
S50																				
L52																				
K54																				
L55																				
K66																				
V67																				
Q68																				

Residue	$B_{1s}$ (Hz)	Global Analysis								Residue-Specific Analysis									
		$k_{ex}$ (s <sup>-1</sup> )	$p_{U'}$ (%)	$\Delta\varpi$ (ppm)	$R_{1,F}$ (s <sup>-1</sup> )	$R_{2,F}$ (s <sup>-1</sup> )	$R_{2,U'}$ (s <sup>-1</sup> )	$\chi^2_{red}$	$K$ (s <sup>-1</sup> )	$k_{ex}$ (s <sup>-1</sup> )	$p_{U'}$ (%)	$\Delta\varpi$ (ppm)	$R_{1,F}$ (s <sup>-1</sup> )	$R_{2,F}$ (s <sup>-1</sup> )	$R_{2,U'}$ (s <sup>-1</sup> )	$\chi^2_{red}$	$K$ (s <sup>-1</sup> )		
T13	6.0, 18.4, 46.0	69.9 ± 1.4	1.35 ± 0.02	4.07 ± 0.02	1.23 ± 0.01	15.2 ± 0.1	133.4 ± 11.2	1.06	119	57.0 ± 5.3	1.45 ± 0.07	4.06 ± 0.03	1.23 ± 0.01	15.4 ± 0.1	135.3 ± 11.5	0.81	105		
K22																			
K26																			
K28																			
R29																			
N33																			
E37																			
K41																			
M42																			
I43																			
I44																			
N45																			
S50																			
L52																			
K54																			
L55																			
K66																			
V67																			
Q68																			

**Table S3** Best fit exchange parameters extracted from two different sets of <sup>15</sup>N CEST profiles recorded using the 4 mM U-[<sup>15</sup>N] A39G FF sample at 2.5 °C (11.7 T).



## References

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