

High-density volumetric super-resolution microscopy

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1 **Volumetric super-resolution microscopy typically encodes the**
2 **3D position of single-molecule fluorescence into a 2D image**
3 **by changing the shape of the point spread function (PSF) as**
4 **a function of depth. However, the resulting large and complex**
5 **PSF spatial footprints reduce temporal resolution by requir-**
6 **ing lower labelling densities to avoid overlapping fluorescent**
7 **signals. We quantitatively compare the density dependence of**
8 **single-molecule light field microscopy (SMLFM) to other 3D**
9 **PSFs (astigmatism, double helix and tetrapod) showing that**
10 **SMLFM enables an order-of-magnitude speed improvement**
11 **compared to the double helix PSF by resolving overlapping**
12 **emitters through parallax. We then experimentally demon-**
13 **strate the high accuracy ($>99.2 \pm 0.1\%$, $0.1 \text{ locs } \mu\text{m}^{-2}$) and**
14 **sensitivity ($>86.6 \pm 0.9\%$, $0.1 \text{ locs } \mu\text{m}^{-2}$) of SMLFM at point**
15 **detection through whole-cell (scan-free) imaging and tracking**
16 **of single membrane proteins in live primary B cells. We also**
17 **exemplify high density volumetric imaging ($0.15 \text{ locs } \mu\text{m}^{-2}$)**
18 **in dense cytosolic tubulin datasets.**

19 **3D super-resolution microscopy | single-molecule light field microscopy |**
20 **SMLFM | Fluorescence microscopy | Fourier light field | high-throughput**
21 **imaging | PSF engineering |**

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23 Single-molecule localization microscopy (SMLM) is a
24 super-resolution technique that separates the fluorescence
25 emission of individual fluorophores temporally to observe
26 biological systems with sub-diffraction resolution (1–4). Di-
27 rect imaging in three dimensions (3D) enables the study
28 of complex biological morphologies and dynamic processes
29 that would otherwise be underestimated in 2D.

30 In SMLM, the fluorescence from a single emitter is ob-
31 served as a diffraction-limited spot on a detector, known
32 as the *point spread function* (PSF). Generally, 3D-SMLM
33 employs optical elements that transform the standard 2D
34 PSF into spatial distributions that also encode axial position.
35 These 3D PSFs exhibit lateral spatial footprints that are
36 much larger in area than the standard PSF, meaning the pro-
37 jection of a 3D volume onto a 2D detector usually necessi-
38 tates considerably slower acquisition rates (typically 5 to 10-
39 fold) due to a higher likelihood of PSF overlap (5, 6). How-
40 ever, the number of emitters localised per frame governs

41 temporal resolution and therefore dense emitter datasets are
42 desirable. This is exemplified in recent work from Legant
43 *et al.* where impressive super-resolved whole-cell volumes
44 were obtained over very long acquisition times (*i.e.* 3–10
45 days) (7, 8). This extended experimental duration was nec-
46 essary to generate an image with resolution comparable to a
47 corresponding electron microscope experiment (8).

48 Long imaging durations present unrealistic conditions
49 for typical cellular experiments and also reduce the quan-
50 tity of biological repeats that can be performed within ap-
51 propriate timescales. While strategies exist to reduce PSF
52 overlap—such as specialised labelling protocols and post-
53 processing algorithms (9–11)—they are ultimately limited
54 by the decrease in lateral resolution at the expense of a
55 greater depth-of-field (DoF). However, a recent study re-
56 vealed a lack of post-processing solutions specifically for
57 dense 3D datasets (12). Hence, to have broad applicabil-
58 ity to the biological community there is a fundamental need
59 for robust strategies to perform 3D-SMLM at high densities.
60 This will bring 3D-SMLM into line with the timescales and
61 workflows of current 2D cellular experiments, and is another
62 important step toward real-time 3D-SMLM.

63 Sub-diffraction axial precision can be achieved by en-
64 gineering the shape of the PSF to simultaneously encode
65 the lateral and axial position of a single emitter in a 2D
66 image (13, 14). A variety of engineered PSFs have been
67 reported, including astigmatism ($\sim 1 \mu\text{m}$ DoF) (15), a bi-
68 sected pupil ($\sim 1 \mu\text{m}$ DoF) (16), the corkscrew PSF ($\sim 3 \mu\text{m}$
69 DoF) (17), the double helix (DH)PSF ($\sim 4 \mu\text{m}$ DoF) (5,
70 18, 19), and the tetrapod PSF ($6\text{--}20 \mu\text{m}$ DoF) (20, 21).
71 On the other hand, single-molecule light field microscopy
72 (SMLFM) (22) is an SMLM technique that places a refrac-
73 tive microlens array (MLA) in the back focal plane (BFP)
74 of a widefield microscope to encode 3D position into the
75 PSF (see Supplementary Note S1) (23, 24). SMLFM pos-
76 sesses a large tuneable DoF, high photon throughput, (13)
77 the PSF can be fitted with conventional 2D algorithms and is
78 wavelength non-specific. The unique advantage of SMLFM
79 is that it operates through parallax whereby the PSF is
80 comprised of several spatiotemporally correlated perspec-

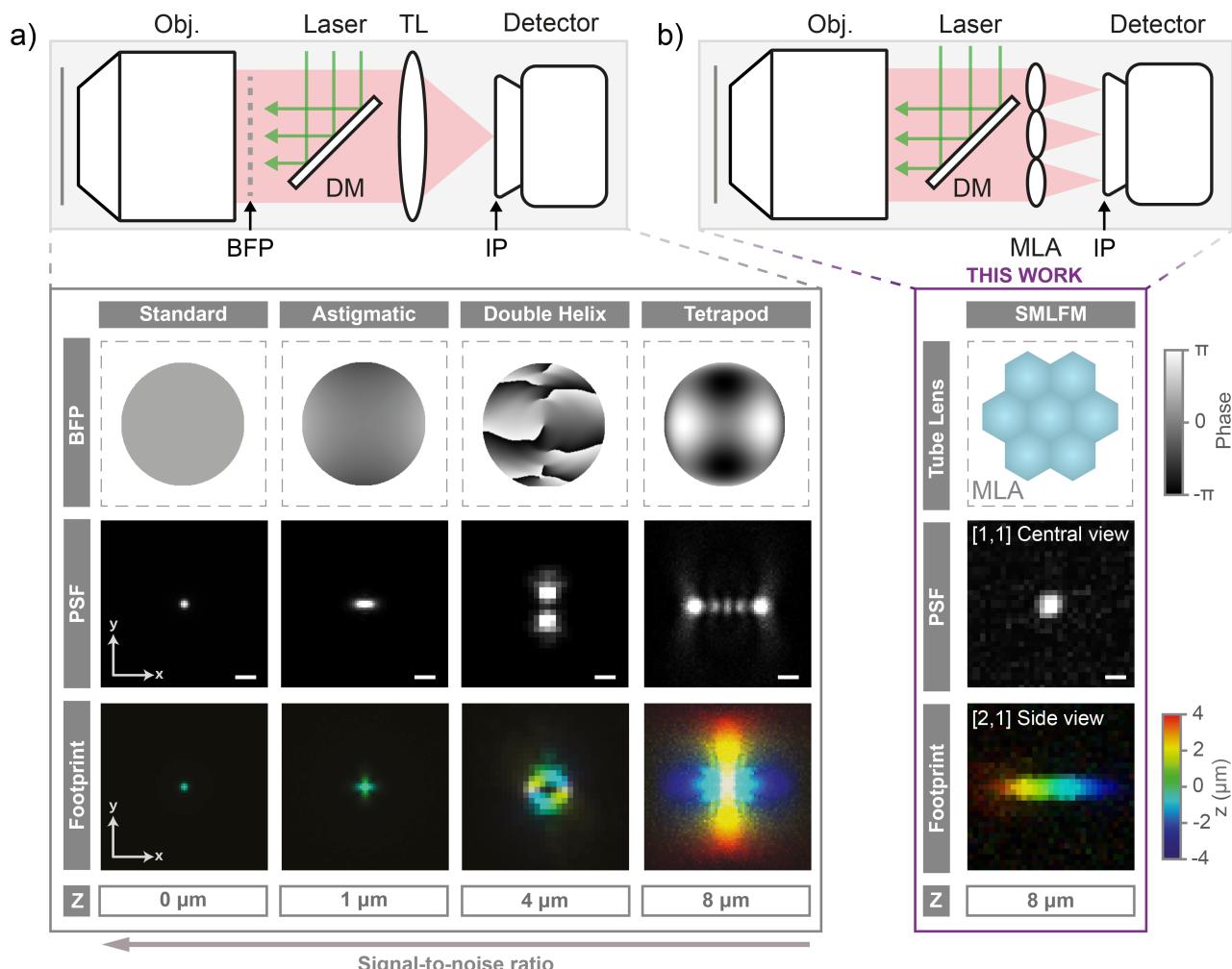


Fig. 1. Encoding the 3D position of single molecule fluorescence into a 2D image. **a)** Optical schematic of a typical widefield microscope, where Obj. = objective lens, TL = tube lens, IP = image plane, BFP = back focal plane, and DM = dichroic mirror. Below are some common 3D point spread functions (PSFs) implemented by phase modulation in the BFP (top row), including the standard PSF, astigmatic PSF, double helix PSF and the tetrapod PSF (middle row) and their associated PSF footprints integrated over their entire axial range (bottom row). Field-of-view is $8 \times 8 \mu\text{m}^{-2}$ and the scale bars are 1 μm. **b)** Optical schematic of a Fourier light field microscope for SMLFM, where MLA = micro-lens array (tube lens), which is placed in the BFP. Below is a schematic of the MLA, the PSF in the central perspective view, and the PSF footprint in the entire 8 μm axial range (see Supplementary Fig. S1 for further details). Pixel size is 110 nm for standard, astigmatic and the tetrapod PSF, and 266 nm for the DHPSF and SMLFM to reflect experimental parameters.

51 **tive views displaced in proportion to the curvature of the**
 52 **wavefront. As such, SMLFM is particularly suited to high**
 53 **spot densities for two key reasons:**

54

- 55 **1. Single emitters that occur at different axial planes (but**
 56 **overlap laterally) are imaged at different locations in**
 57 **different perspective views and can be distinguished.**

 58 - 59 **2. We illustrate a redundancy in that a localisation is not**
 60 **required in every perspective view to be localised in**
 61 **3D.**

62 Multi-focal plane microscopy also segments the BFP
 63 to image two or more focal planes and capture 3D vol-
 64 ume (25–28). However, this work will focus on techniques
 65 that yield sub-diffraction axial precision over extended axial
 66 ranges.

67 In the present work, we report the first hexagonal
 68 SMLFM platform capable of super-resolving single emitters
 69 at very high densities over an 8 μm DoF. We quantitatively

70 **compare the performance of SMLFM to other common 3D**
 71 **PSFs as a function of spot density through simulations. We**
 72 **then apply SMLFM experimentally to the scan-free imag-**
 73 **ing and tracking of individual B-cell receptors and the imag-**
 74 **ing of tubulin in HeLa cells to show that overlapping PSFs**
 75 **minimally affect the localisation precision and that existing**
 76 **labelling strategies can now be directly transferred to 3D**
 77 **imaging pipelines.**

78 **Results and Discussion**

79 **Density dependence of 3D PSFs.** Current state-of-the-
 80 art 3D PSFs are typically created by phase modulation in the
 81 BFP of the objective lens (*i.e.* with a phase mask) as shown
 82 in Fig. 1a. This phase modulation gives rise to spatial dis-
 83 tributions of intensity in the imaging plane that change as a
 84 function of the axial position of the emitter. Alternatively, in
 85 SMLFM the MLA segments the BFP and focuses an array
 86 of spots on the detector as shown in Fig. 1b and Supplemen-

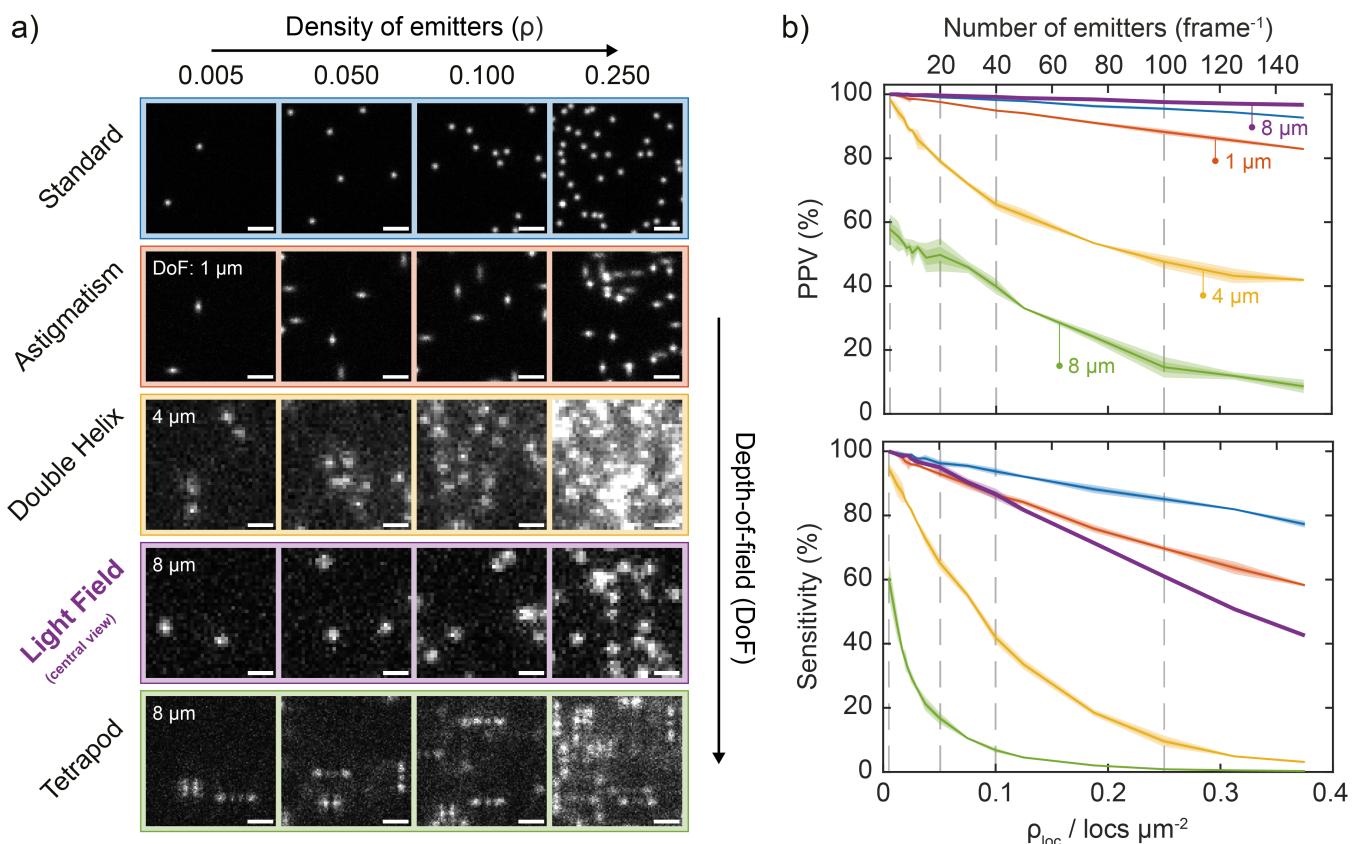


Fig. 2. SMLFM consistently outperforms other 3D-SMLM techniques at correctly identifying and reconstructing single emitters at increasing densities. **a)** Snapshots of simulated raw localization datasets in a $10 \times 10 \mu\text{m}^{-2}$ region for each imaging modality discussed herein (2D, astigmatism, double helix PSF, light field [central view] and tetrapod PSF.) Scale bar represents 2 μm . **b)** Top: Average positive predictive value (PPV) curves for each SMLM technique as a function of emitter density (ρ_{loc}) at 4,000 detected photons, where PPV refers to the number of true positive localizations vs. total number of fitted localizations. Bottom: Average sensitivity curves as a function of ρ_{loc} at 4,000 detected photons, where sensitivity refers to the number of true positive localizations vs. total number of ground truth localizations. Shaded regions represent first and second standard deviation from the mean over three repeats of 100-frame simulated datasets. Example simulated datasets are presented in Supplementary Movie 1.

115 tary Fig. S1. All of these changing PSFs can be understood
 116 in the context of high-density imaging by collapsing the en-
 117 tire PSF onto the 2D detector, which we define as the *PSF*
 118 *footprint*.

119 Raw localization datasets were simulated for the SMLM
 120 modalities presented in Fig. 1 (standard, astigmatic, double
 121 helix, light field and tetrapod) to investigate the effect the
 122 of PSF footprint on the ability to resolve single emitters at
 123 high densities (Fig. 2a). Briefly, the emitter density (ρ_{loc})
 124 of simulated SMLM data was systematically increased from
 125 0.005 μm^{-2} (2 localizations per $20 \mu\text{m} \times 20 \mu\text{m}$ field-of-
 126 view, FoV) to 0.375 μm^{-2} (150 localizations) and subse-
 127 quently processed using conventional fitting algorithms (see
 128 Methods). Each dataset was also simulated for typical pho-
 129 ton values expected for a fluorescent protein (1,000 pho-
 130 tons), an organic dye (4,000 photons) and a next-generation
 131 fluorescent probe (10,000 photons) to reflect different la-
 132 belling scenarios. Computational multi-emitter fitting was
 133 not implemented in the analysis to ensure a fair compari-
 134 son across methods since algorithms are at different lev-
 135 els of technical development for each technique (21, 29),
 136 and single-emitter algorithms have been previously shown

137 to outperform multi-emitter algorithms in high density 3D
 138 SMLM scenarios (12).

139 Direct comparison of PSF footprint with DoF (presented
 140 in Supplementary Fig. S3a) reveals how the area of each
 141 3D PSF changes with axial position. The light field PSF
 142 is highly competitive, achieving an axial range suitable for
 143 imaging entire cells (8 μm) with a PSF area 55% the size of
 144 the DHPSF on average and 18% the area of the tetrapod PSF.
 145 As every perspective view comprises a super-resolvable im-
 146 age of the sample, only the area of the SMLFM PSF in each
 147 perspective view is needed for direct comparison. The sim-
 148 ple and compact PSF footprint of SMLFM is a principle
 149 component in the ability to resolve single emitters at higher
 150 spot densities than the double helix and tetrapod PSFs. In
 151 this work we consider the central view for density studies.

152 Several quality-of-imaging metrics were then computed
 153 for each simulated dataset classifying a localization as ei-
 154 ther a true positive (TP), false positive (FP) or false nega-
 155 tive (FN) with respect to known ground truth (GT) coor-
 156 dinates. The positive predictive value (PPV, also known as
 157 *precision*) describes the fraction of TP localizations relative
 158 to all localizations (TP+FP). Sensitivity (also known as *re-*

call) describes the fraction of accurate localizations that are retrieved (TP/GT). Both PPV and sensitivity are presented as a function of ρ_{loc} in Fig. 2b using datasets simulated at 4,000 detected photons to reflect labelling using an organic dye molecule (see Supplementary Fig. S5 for PPV and sensitivity plots at 1,000, 4,000 and 10,000 detected photons).

Low signal-to-noise ratio (SNR), high background fluorescence and emitter overlap contribute to reconstruction artefacts from the incorrect localization of single emitters. This leads to a decrease in both PPV and sensitivity as a function of ρ_{loc} , in agreement with similar work (6). Unlike for the double helix or tetrapod PSFs, the PPV for SMLFM is linear across the whole ρ_{loc} range with a maximum value of $100.0 \pm 0.0\%$ (mean \pm SD) at $\rho_{loc} = 0.005 \mu\text{m}^{-2}$ and a minimum of $96.7 \pm 0.1\%$ at $\rho_{loc} = 0.375 \mu\text{m}^{-2}$, which can be rationalised by the spatio-temporally correlated PSF filtering out stochastic noise due to the requirement for the same emitter to be localized in each perspective view. PPV for the standard and astigmatic PSFs is also linear across all values of ρ_{loc} as expected from compact (but very low DoF) PSF footprints. Conversely, with an average pixel area of $1.8 \times$ that of the SMLFM PSF, the DHPSF exhibits a non-linear response to ρ_{loc} and a much lower PPV than SMLFM and likewise with the tetrapod PSF. Their weaker resistance to increasing ρ_{loc} can be attributed to their greater size and complexity of photon distributions, for example the tetrapod PSF was specifically designed for optimal Fisher information, and hence low density imaging scenarios (21, 30).

A linear relationship between sensitivity and ρ_{loc} is also observed for SMLFM with a maximum value of $100.0 \pm 0.0\%$ when $\rho_{loc} = 0.005 \mu\text{m}^{-2}$ and a minimum of $42.5 \pm 0.1\%$ when $\rho_{loc} = 0.375 \mu\text{m}^{-2}$. This is a result of distinguishing overlapping emitters through parallax, whereby single-molecule fluorescence is observed at different positions in each perspective view. Resolving overlapping emitters through the double helix and tetrapod PSF shaping methods is either impossible or computationally expensive during post processing. Alternatively, SMLFM facilitates these higher localization rates by resolving emitters through parallax, described herein as optical multi-emitter fitting (distinct from computational multi-emitter fitting). These data combined demonstrate that SMLFM has the capacity to localise $86.6 \pm 0.9\%$ of all emitters at a typical 2D-SMLM localization density of $\sim 0.1 \mu\text{m}^{-2}$ without compromising on the total number of localizations (31). Even at an incredibly high localization density of $0.375 \mu\text{m}^{-2}$ SMLFM is able to recover 43% of all ground truth localisations while this is less than 1% for the double helix and tetrapod PSFs.

By comparing the spot densities at which the DHPSF

and SMLFM achieve equal error rates we determine a maximum speed improvement of $8.95 \times$ for SMLFM at an error rate of 13.5% (at which 86.5% of all localisations are correctly reconstructed in 3D, see Supplementary Note S4 section S4.2 and Supplementary Fig. S4) for 4,000 detected photons. This represents the upper practical limit in what SMLFM can achieve in direct comparison with the DHPSF (the state-of-the-art 3D SMLM modality for DoF and localization precision.) Furthermore, this maximum practical speed improvement was measured to be $25.3 \times$ (error rate of 40.0%) at 1,000 detected photons and $10.6 \times$ (error rate of 11.0%) at 10,000 photons. Therefore, on the basis of speed, SMLFM significantly out-performs the DHPSF at all light levels, particularly at low SNR, aided by the division of background photons across seven lenses.

SMLFM captures the heterogeneity of live B-cell membrane receptors. The density-dependence studies reveal an optical redundancy in SMLFM that would be suited to the high-density volumetric imaging of entire cells through optical multi-emitter fitting. To challenge our method we imaged whole primary mouse B-cell membranes in a scan-free dSTORM modality previously optimised for 2D-SMLM (32–34). The 3D organisation of membrane receptors on immune cells, such as the B-cell receptor (BCR) is of increasing scientific interest to better understand the immune response to infection (35, 36). Single BCR complexes were labelled with a single molecule of Alexa-Fluor 647 (see Methods) and imaged under an inclined illumination angle to improve contrast (see Supplementary Note S1 section S1.3). An average of 40,000 3D localizations were accumulated per cell over an axial range of $\sim 8 \mu\text{m}$ (Fig. 3a-d). Membrane ruffles and microvilli could be observed, consistent with sub-diffraction resolution being obtained (37).

3D localizations were collected over a $\sim 50 \mu\text{m}^2$ circular detector area (image space) with an average localization density of $\sim 0.10 \mu\text{m}^{-2}$ (see Supplementary Fig. S7) corresponding to a PPV of $99.2 \pm 0.1\%$ and sensitivity of $86.6 \pm 0.9\%$. A maximum localization density of $\sim 0.24 \mu\text{m}^{-2}$ was achieved for a small portion of the experiment, which corresponds to a PPV of $97.6 \pm 0.0\%$ and sensitivity of $61.0 \pm 0.2\%$. In comparison, at an average density of $\sim 0.10 \mu\text{m}^{-2}$ the DHPSF would be expected to achieve a PPV of $65.7 \pm 1.4\%$ and a sensitivity of $24.0 \pm 0.5\%$, accurately localizing a quarter of all emitters. Equally, the tetrapod PSF would be expected to correctly localize $39.9 \pm 1.2\%$ of all emitters and recover $6.8 \pm 0.4\%$ of all localizations. The pronounced improvement in performance shown here by SMLFM at high localization densities, in addition to the imaging of complete cellular volumes with-

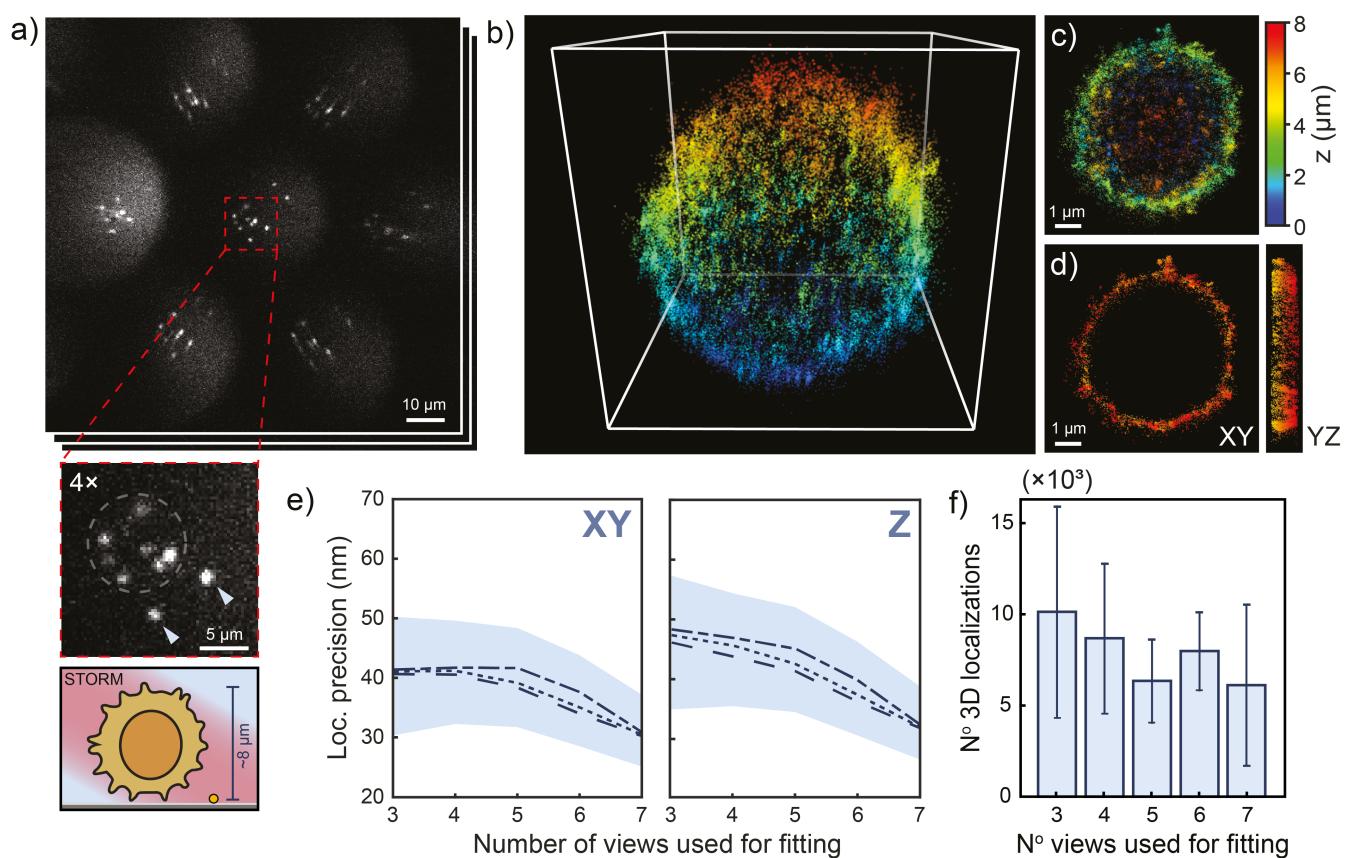


Fig. 3. Scan-free SMLFM-STORM imaging of B cell receptors over whole primary mouse B cells. **a)** Raw SMLFM data of individual membrane receptors comprising 7 perspective views in a hexagonal arrangement. Expanded insert shows seven fluorescent puncta (Alexa Fluor 647) in the central perspective view and two fiducial markers indicated with arrows. Directly below is an illustration of the cell being imaged. **b)** Associated 3D reconstruction of the whole primary mouse B cell (40,000 3D localizations in a 9 μm³ box). **c)** an xy projection and **d)** a 1 μm⁻² thick central clipping to illustrate non-internalisation of dye molecules. **e)** Median lateral and axial fitting error for localizations below 60 nm precision as a function of the number of views used to reconstruct a 3D localization (shading represents interquartile range). **f)** Associated proportion of 3D localizations below 60 nm lateral precision as a function of the number of views used to reconstruct a 3D coordinate.

out scanning, enables a significant improvement in sample throughput in future 3D-SMLM experiments.

SMLFM is advantageous at high densities because single emitters are not required to be isolated in every perspective view to be localized in 3D. We quantified this redundancy that enables optical multi-emitter fitting in these large cellular datasets by considering the localization precision as a function of perspective views used for PSF fitting. Fig. 3e reveals excellent localization precisions of ~40 nm laterally (median) and ~47 nm axially and these values improved to ~30 nm laterally and ~34 nm axially as the number of perspective views for fitting was systematically increased from 3 to 7. This is consistent with a higher effective numerical aperture and better sampling of the PSF position when utilizing a greater number of views. Attempts were made to ensure the distribution of localizations per number of views was equal, see Fig. 3f. Taken together, these data show that optical multi-emitter fitting via parallax is a powerful approach to 3D localizing single molecules at high densities within cells.

Another important application of the high emitter density measurements afforded by SMLFM is single-particle

tracking (SPT) (38–40). 3D-SPT better quantifies diffusive processes than 2D measurements, which tend to underestimate diffusion rates (41, 42). A previous study of membrane protein mobility highlighted the importance of imaging diffusion dynamics away from the glass interface (basal surfaces) (39), which sparked the imaging of apical surfaces in 4 μm optical sections using the DHPSF (5). SMLFM boasts a significant practical advancement over this work, which is two-fold:

1. The larger DoF ensures single proteins can be tracked over entire cell volumes without scanning.
2. Localizing emitters through parallax improves the ability to delineate trajectories that would otherwise be occluded at higher densities.

To demonstrate this we applied SMLFM to the 3D-SPT of BCR complexes found on the surface of live mouse B cells, accumulating hundreds of trajectories in <10,000 frames (~5 minutes) with an average track length of 12.5 points (Fig. 4a). Maximum likelihood estimation of the diffusion coefficient from trajectories over 5 cells (a total of 1806 tracks) yielded a distribution of diffusion coefficients

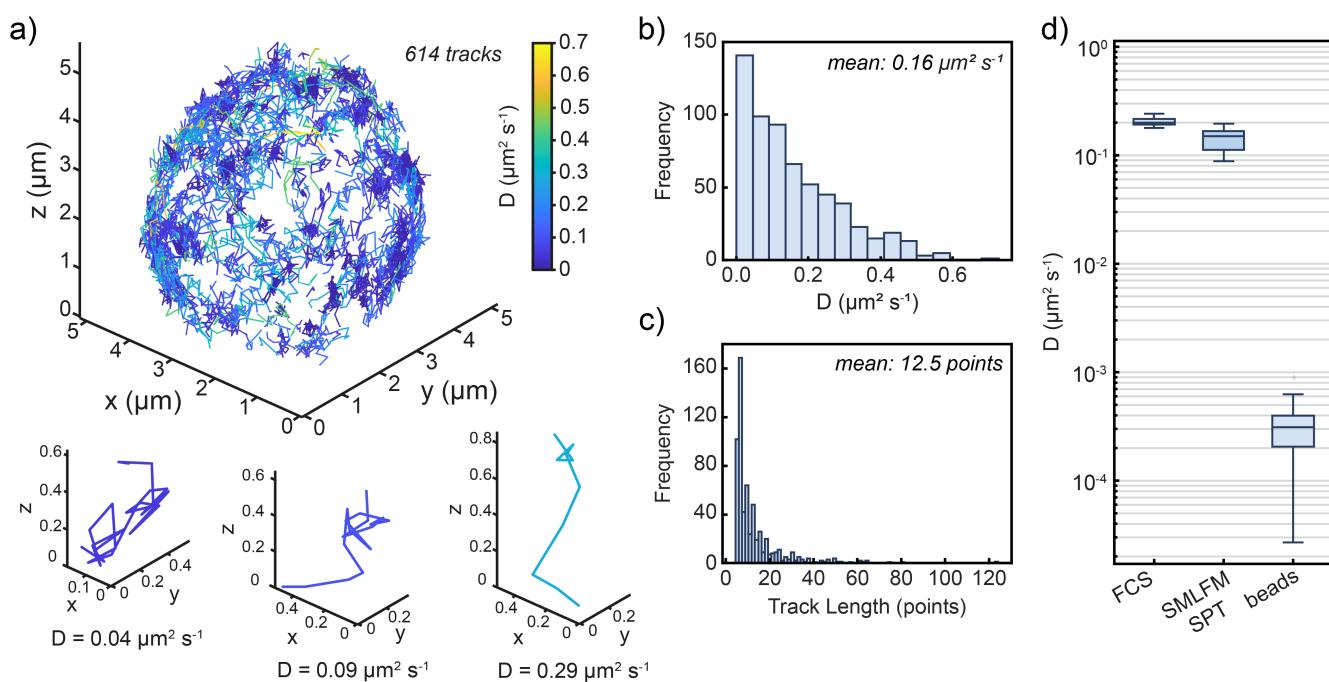


Fig. 4. Scan-free whole-cell 3D SPT of the B cell receptor on primary mouse B-cell membranes using SMLFM. **a)** 3D trajectory map of the BCR over a whole primary mouse B cell totalling 614 tracks color-coded by diffusion coefficient using maximum likelihood estimation. Example isolated trajectories of varying diffusion coefficient are expanded directly below. Histogram of **b)** associated diffusion coefficients and **c)** track lengths (bin widths were determined using Freedman-Diaconis' rule). **d)** Median diffusion coefficient of the BCR measured by FCS ($N = 7$) and SMLFM-SPT ($N = 5$), and the minimum diffusion coefficient measurable by SMLFM-SPT determined with immobilised beads ($N = 9$). SMLFM-SPT comprises a total of 1806 trajectories over 5 cells with a mean track length of 8 points.

(Fig. 4b-d) for individual BCR complexes with a median value of $0.14 \pm 0.08 \mu\text{m}^2 \text{s}^{-1}$, consistent with that observed by Tolar *et al.* on resting murine B cells (43). To confirm this, we measured a median diffusion coefficient of $0.20 \pm 0.01 \mu\text{m}^2 \text{s}^{-1}$ at the apical surface using fluorescence correlation spectroscopy.

SMLFM effectively and accurately captures the heterogeneity of diffusion coefficients of surface receptors and opens up the possibility of the direct observation of dynamic BCR clustering following or proceeding antigen encounters (and in general the clustering of key signalling proteins in other systems), which is of great interest in the study of receptor triggering (44).

High-density SMLFM resolves intracellular structure.

To demonstrate intracellular imaging at very high emitter density, we performed scan-free dSTORM imaging of tubulin in fixed HeLa cells with SMLFM. Fig. 5a contains a snapshot of raw localization data with a cell occupying a $40 \mu\text{m} \times 40 \mu\text{m}$ FoV, which spanned an axial range of $\leq 3 \mu\text{m}$, with the 3D reconstruction shown in Fig. 5b. An expanded region is presented in Fig. 5c. Line profiles (Figure 5d, width 400 nm) were drawn for two ranges to confirm the resolution of individual microtubules. A maximum of 40 3D localizations were detected per image frame, with an average of ~ 22 per frame, totalling 150,000 localizations over ~ 4 minutes (30 ms detector exposure, Fig. 5e). This corresponds to a maximum density of $0.15 \mu\text{m}^{-2}$ and an average

of $0.075 \mu\text{m}^{-2}$, whereupon through simulations SMLFM is shown to retrieve $90.4 \pm 0.6\%$ of all localizations (sensitivity) with an accuracy of $99.4 \pm 0.0\%$ (PPV). A median of 3,900 photons were detected per 3D localization (Fig. 5f) achieving median lateral and axial localization precisions of 42.6 nm and 46.7 nm, respectively (Supplementary Fig. S10). A Fourier Shell Correlation (FSC) of 54 nm resolution at a 1/7 cut-off was calculated from the localizations presented in Figure 5c.

Furthermore, no 3D-specific sample optimisation was undertaken prior to imaging and a dSTORM buffer protocol was implemented that was previously developed for 2D-SMLM (33). As protocol optimisation for SMLM can be time intensive and challenging, this facile translation from 2D to 3D SMLM presents a significant advantage in the future use of 3D-SMLM in biological research (45).

Conclusions

We report the first hexagonal SMLFM platform enabling 3D-SMLM over an $8 \mu\text{m}$ axial range and quantitatively compared its performance to other common 3D PSFs through simulations revealing an order-of-magnitude speed improvement compared to DHPSF microscopy. We attribute this speed improvement to optical multi-emitter fitting through which overlapping emitters in an imaging volume can now be resolved through a redundancy in the number of perspective views required for 3D reconstruction. We

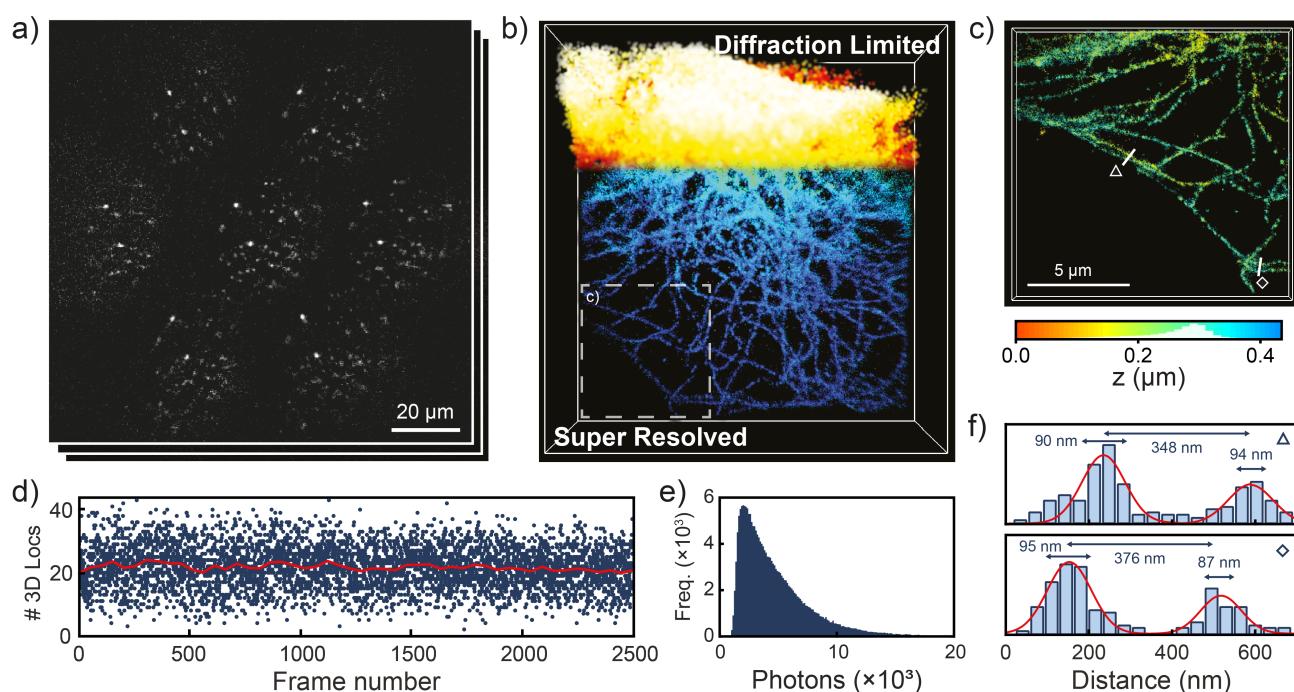


Fig. 5. SMLFM-STORM imaging of Alexa Fluor 647-labelled tubulin in a HeLa cell. **a)** A snapshot of raw localization data in microtubule-stained HeLa cells imaged through the 7-hex SMLFM platform and **b)** the corresponding super-resolved 3D volume. The super-resolved area is $23 \mu\text{m} \times 12 \mu\text{m}$ in size and contains 150,000 3D localizations. **c)** Expansion of the marked region (60,000 3D localizations) in **b)**. **d)** localization rate over the first 2500 frames indicating a mean 3D localization rate (red line, rolling average over 100 frames) of 22 frame^{-1} and an upper limit of $\sim 40 \text{ frame}^{-1}$ corresponding to ~ 0.075 and $\sim 0.15 \text{ locs } \mu\text{m}^{-2}$ respectively. **e)** Histogram of detected photons per 3D localization (a median value of 3,900). **f)** Line plots through pairs of microtubules (width of 400 nm) in panel **d)** showing individual microtubules being resolved.

353 applied SMLFM experimentally to the imaging in both live
354 and fixed whole cells and dense arrays of cytosolic tubu-
355 lin, where it consistently localised single emitters in 3D at
356 high, non-optimised, densities achieving localization preci-
357 sions $\sim 40 \text{ nm}$ laterally and $\sim 50 \text{ nm}$ axially.

358 Future endeavours could couple SMLFM with computa-
359 tional multi-emitter fitting and/or deep learning strategies,
360 high speed detectors, and alternative volumetric labelling
361 strategies to push localization rates even further. We antici-
362 pate the uptake of SMLFM as a powerful tool in improving
363 our understanding of 3D nano-scale architecture, dynamics
364 and will bring robust real-time 3D super-resolution imaging
365 to life scientists.

366 References

1. Rust, M.J., Bates, M. and Zhuang, X. Sub-diffraction-limit imaging by stochastic optical reconstruction microscopy (STORM). *Nature Methods*, 3(10):793–796, October 2006. ISSN 1548-7105. doi: 10.1038/nmeth929. Number: 10 Publisher: Nature Publishing Group.
2. Betzig, E., Patterson, G.H., Sougrat, R., Lindwasser, O.W., Olenych, S., Bonifacino, J.S., Davidson, M.W., Lippincott-Schwartz, J. and Hess, H.F. Imaging Intracellular Fluorescent Proteins at Nanometer Resolution. *Science*, 313(5793):1642–1645, September 2006. ISSN 0036-8075, 1095-9203. doi: 10.1126/science.1127344. Publisher: American Association for the Advancement of Science Section: Report.
3. Hess, S.T., Girirajan, T.P. and Mason, M.D. Ultra-high resolution imaging by fluorescence photoactivation localization microscopy. *Biophysical Journal*, 91(11):4258–4272, December 2006. doi: 10.1529/biophysj.106.091116.
4. Lelek, M., Gyparaki, M.T., Beliu, G., Schueder, F., Griffié, J., Manley, S., Jungmann, R., Sauer, M., Lakadamyali, M. and Zimmer, C. Single-molecule localization microscopy. *Nature Reviews Methods Primers*, 1(1):39, December 2021. ISSN 2662-8449. doi: 10.1038/s43586-021-00038-x.
5. Carr, A.R., Ponjavic, A., Basu, S., McColl, J., Santos, A.M., Davis, S., Laue, E.D., Klennerman, D. and Lee, S.F. Three-Dimensional Super-Resolution in Eukaryotic Cells Using the Double-Helix Point Spread Function. *Biophysical Journal*, 112(7):1444–1454, April 2017. ISSN 00063495. doi: 10.1016/j.bpj.2017.02.023.
6. Nehme, E., Freedman, D., Gordon, R., Ferdinand, B., Weiss, L.E., Alalouf, O., Naor, T., Orange, R., Michaeli, T. and Shechtman, Y. DeepSTORM3d: dense 3d localization microscopy and PSF design by deep learning. *Nature Methods*, 17(7):734–740, June 2020. doi: 10.1038/s41592-020-0853-5.
7. Chen, B.C., Legant, W.R., Wang, K., Shao, L., Milkie, D.E., Davidson, M.W., Janetopoulos, C., Wu, X.S., Hammer, J.A., Liu, Z., English, B.P., Mimori-Kiyosue, Y., Romero, D.P., Ritter, A.T., Lippincott-Schwartz, J., Fritz-Laylin, L., Mullins, R.D., Mitchell, D.M., Bembeneck, J.N., Reymann, A.C. et al. Lattice light-sheet microscopy: Imaging molecules to embryos at high spatiotemporal resolution. *Science*, 346(6208), October 2014. doi: 10.1126/science.1257998.
8. Legant, W.R., Shao, L., Grimm, J.B., Brown, T.A., Milkie, D.E., Avants, B.B., Lavis, L.D. and Betzig, E. High-density three-dimensional localization microscopy across large volumes. *Nature Methods*, 13(4):359–365, March 2016. doi: 10.1038/nmeth.3797.
9. Holden, S.J., Uphoff, S. and Kapanidis, A.N. DAOSTORM: an algorithm for high-density super-resolution microscopy. *Nature Methods*, 8(4):279–280, April 2011. ISSN 1548-7091, 1548-7105. doi: 10.1038/nmeth.0411-279.
10. Gu, L., Sheng, Y., Chen, Y., Chang, H., Zhang, Y., Lv, P., Ji, W. and Xu, T. High-Density 3D Single Molecular Analysis Based on Compressed Sensing. *Biophysical Journal*, 106(11):2443–2449, June 2014. ISSN 0006-3495. doi: 10.1016/j.bpj.2014.04.021.
11. Speiser, A., Müller, L.R., Hoess, P., Matti, U., Obara, C.J., Legant, W.R., Kreshuk, A., Macke, J.H., Ries, J. and Turaga, S.C. Deep learning enables fast and dense single-molecule localization with high accuracy. *Nature Methods*, 18(9):1082–1090, September 2021. ISSN 1548-7105. doi: 10.1038/s41592-021-01236-x. Number: 9 Publisher: Nature Publishing Group.
12. Sage, D., Pham, T.A., Babcock, H., Lukes, T., Pengo, T., Chao, J., Velmurugan, R., Herbert, A., Agrawal, A., Colabrese, S., Wheeler, A., Archetti, A., Rieger, B., Ober, R., Hagen, G.M., Sibarita, J.B., Ries, J., Henriques, R., Unser, M. and Holden, S. Super-resolution fight club: assessment of 2D and 3D single-molecule localization microscopy software. *Nature Methods*, 16(5):387–395, May 2019. ISSN 1548-7091, 1548-7105. doi: 10.1038/s41592-019-0364-4.
13. Grover, G., Quirin, S., Fiedler, C. and Piestun, R. Photon efficient double-helix PSF microscopy with application to 3D photo-activation localization imaging. *Biomedical Optics Express*, 2(11):3010–3020, November 2011. ISSN 2156-7085. doi: 10.1364/BOE.2.003010. Publisher: Optical Society of America.
14. Rehman, S.A., Carr, A.R., Lenz, M.O., Lee, S.F. and O'Holleran, K. Maximizing the field of view and accuracy in 3D Single Molecule Localization Microscopy. *Optics Express*, 26(4):4631–4637, February 2018. ISSN 1094-4087. doi: 10.1364/OE.26.004631. Publisher: Optical Society of America.
15. Huang, B., Wang, W., Bates, M. and Zhuang, X. Three-dimensional Super-resolution

549 Methods

550 **Optical setup.** The SMLFM platform described in this work was
551 constructed using an epi-fluorescence microscope (Eclipse Ti-U, Nikon)
552 housing a 1.27 NA water immersion objective lens (Plan Apo VC 60×,
553 Nikon, Tokyo, Japan) for imaging above the coverslip. The z-position of
554 the objective was controlled with a scanning piezo (P-726 PIFOC, PI, Karl-
555 ruhe, Germany). The Fourier lens ($f = 175$ mm, ThorLabs) was placed in
556 a $4f$ configuration with the tube lens ($f = 200$ mm, Nikon) to relay the back
557 focal plane (BFP) outside of the microscope body (see Supplementary Fig.
558 S1). A hexagonal microlens array ($f = 175$ mm, pitch = 2.39 mm, custom-
559 made by CAIRN) was placed in the BFP to relay the image plane onto an
560 EMCCD (Evolve Delta 512, Photometrics, Tucson, AZ, 16 μm pixel size).
561 Excitation was achieved using a 640 nm (~ 10 kW cm^{-2} power density,
562 150 mW, iBeam Smart-S 640-S, Toptica, Munich, Germany) and activation
563 by a 405 nm (~ 0.04 kW cm^{-2} power density, 120 mW, iBeam Smart-
564 S 405-S, Toptica, Munich, Germany) laser, that were circularly polarised,
565 collimated and focused on to the BFP of the objective to create an evanes-
566 cent excitation wave. Unless stated otherwise, samples were excited with
567 a highly inclined and laminated optical sheet (HILO) which was achieved
568 by laterally displacing the excitation beam towards the edge of the BFP
569 of the objective (see Supplementary Note S1 section S1.3). Fluorescence
570 was collected by the same objective and separated from the excitation beam
571 using a quad-band dichroic mirror (Di01-R405/488/561/635-25x36, Sem-
572 rock, Rochester, NY). Long-pass (BLP02-640R-25, Semrock) and band-
573 pass (FF01-680/X-25, Semrock) emission filters were placed immediately
574 before the detector to isolate fluorescence emission. The pixel size in image
575 space was measured at 266.

576 **3D reconstruction of SMLFM data.** All experimental data were
577 recorded as .tif stacks. 2D gaussian fitting of all emitter positions in all
578 perspective views was carried out in Fiji using PeakFit (GDSC SMLM 2.0)
579 to yield a set of 2D localizations for each raw frame. Given this initial set
580 of 2D localizations, individual emitters were localized in 3D using custom
581 Matlab scripts as outlined in (22). Briefly, the most likely subset of 2D
582 localizations in different perspective views corresponding to a unique emit-
583 ter were identified. Provided that this set of localizations contained more
584 than 3 elements, the 3D location of this emitter was calculated as the least-
585 squares estimate to an optical model relating axial emitter position to the
586 parallax between perspective views. If the residual light field fit error was
587 below 200 nm, the fit was accepted and the subset of 2D localizations was
588 removed. This procedure was repeated for each individual emitter. Drift
589 correction was performed by localizing the position of a fiducial marker in
590 each frame and subtracting the resulting 3D fiducial points from all local-
591 izations of the corresponding frame. System and sample aberrations were
592 corrected for by subtracting the residual disparity (calculated for data ac-
593 quired for all emitters localized during the first 1,000 frames) from all 2D
594 localizations prior to calculating the 3D light field fit. For full details of the
595 light field localization fitting procedure refer to the Supplementary Infor-
596 mation of (22). 3D visualisation was carried out in ViSP (46).

597 **Optical 3D calibration.** Fluorescent beads (200 nm, Deep Red Flu-
598 oSpheres, ThermoFisher, Waltham, MA) were immobilised on a glass slide
599 and imaged to calibrate for deviations in experimental and calculated the
600 disparity from the SMLFM optical model. Glass slides were cleaned under
601 argon plasma (PDC-002, Harrick Plasma, Ithaca, NY) for 1 hour and incu-
602 bated with poly-L-lysine (PLL, 50 μL , 0.1% w/v, Sigma-Aldrich, P820) for
603 10 minutes. Glass slides were washed with PBS (3 \times 50 μL) and incubated
604 with fluorescent beads (50 μL , ca. 3.6×10^8 particles/mL) incubated for 3
605 minutes before washing further with PBS (3 \times 50 μL). The piezo stage (P-
606 726 PIFOC, PI, Karlsruhe, Germany) was used to scan the objective lens
607 axially over 8 μm recording 10 frames at 30 ms exposure per 60 nm in-
608 crement. The data was reconstructed in 3D and plotted against the known
609 movement of the piezo stage. A linear fit was applied to the calibration

610 curve, the gradient of which was a correction factor subsequently applied
611 to all reconstructed data presented in this work.

612 **SPT analysis.** Following 3D reconstruction of SMLFM data a
613 custom-written MATLAB code was implemented to temporally group lo-
614 calizations into single trajectories. Some parameters were chosen by the
615 user, including number of dark frames, linking distance, and minimum
616 track length. The diffusion coefficient was then calculated from each tra-
617 jectory using maximum likelihood estimation, which has previously been
618 shown to yield statistically robust measurements of the diffusion coeffi-
619 cient (47).

620 To determine the minimum observable diffusion coefficient, fluorescent
621 beads (200 nm, Deep Red FluoSpheres, ThermoFisher, Waltham, MA)
622 were immobilised on a glass slide and imaged under conditions (641 nm
623 excitation at ~ 2 mW cm^{-2} power density, 20 ms exposure time) that arti-
624 ficially reproduce the same photon intensities as PA-JF646 used for SPT
625 experiments. The raw data was reconstructed in 3D and trajectories anal-
626 ysed as described previously (22) to yield the smallest resolvable diffusion
627 coefficient.

628 **Analysis of simulated data.** 2D and astigmatic datasets were fit-
629 ted in PeakFit (GDSC SMLM 2.0, Fiji plug-in) using a circular and astig-
630 matic Gaussian PSF, respectively. DHPSF datasets were initially fitted in
631 PeakFit using a circular Gaussian PSF before 3D reconstruction using DH-
632 PSFU (<https://github.com/TheLauLab/DHPSFU>). SMLFM datasets were
633 initially fitted in PeakFit using a circular Gaussian PSF before 3D recon-
634 struction using a custom MATLAB code described previously. Tetrapod
635 PSF data was fitted using Zola (Fiji plug-in) for 3D reconstruction (48).

636 A custom MATLAB code was written to compare the fitted (3D, 2D for
637 the standard PSF) point data to the ground truth coordinates. Specifically,
638 the root mean square distance matrix is calculated between all ground truth
639 coordinates and all reconstructed data points on a frame-by-frame basis
640 and counted as either a true positive, false positive or false negative given
641 a user-specified distance tolerance. The tolerance applied was different for
642 each technique and dictated by the precision and thresholds (determined by
643 the fitting error) were applied to determine true positive and false positives.

644 **Preparation of coverslips for B-cell imaging.** Glass slides (24
645 \times 50 mm borosilicate, thickness No. 1, Brand, Wertheim, Germany) were
646 washed with propan-2-ol and water, dried under nitrogen and cleaned un-
647 der argon plasma (PDC-002, Harrick Plasma, Ithaca, NY) for 1 hour. Glass
648 slides were then incubated with poly-L-lysine (PLL, 50 μL , 0.1% w/v,
649 Sigma-Aldrich, P820) for 1 hour and washed with filtered (0.02 μm syringe
650 filter, Whatman, 6809-1102) PBS (3 \times 50 μL) before incubation with gold
651 nanoparticles (5 μL , 0.1 μm , Merck) for 20 minutes.

652 For fixed cell imaging, glass slides were then washed with filtered PBS
653 (3 \times 50 μL). 1×10^5 fixed labelled B cells were washed in dSTORM buffer
654 (50 mM Tris-HCl, 10 mM NaCl, 10% glucose, 10 mM MEA, 84 $\mu\text{g}/\text{mL}$
655 catalase, 0.2 mg/mL GLOX, adjusted to pH 8), plated in 20 μL dSTORM
656 buffer and left to settle for > 20 minutes. Prior to imaging, the sample was
657 washed into fresh buffer dSTORM buffer.

658 For live cell tracking, PLL-coated glass slides were prepared as above and
659 placed in filtered PBS. For SPT, the surface was incubated with gold beads
660 as above, washed 3 \times in filtered PBS, and cells labelled with PA-JF646-
661 conjugated Fab-Halo were allowed to settle onto the surface for 5–10 min-
662 utes prior to imaging.

663 For point fluorescence correlation microscopy (pFCS), cells labelled with
664 AF647-conjugated Fab-HaloTag were incubated onto the PLL surface for
665 5–10 minutes and imaged using an Zeiss LSM780 inverted confocal micro-
666 scope using a 40 \times water objective, with the sample excited using a 633 nm
667 He-Ne laser. The confocal volume was placed on the apical surface of the
668 cell membrane and five repeated measurements were taken per cell. Data
669 was analysed using PyCorrFit and the diffusion coefficient calculated from

670 the average transit time (τ), using the confocal beam width as calculated
671 using a solution of 100 nM AF647 HaloTag ligand solution.

672 **B-cell culture and fluorescent labelling.** Primary murine B
673 cells were isolated from the spleens of male C57BL/6J mice aged between
674 8 and 12 weeks. Splenocytes were isolated by mechanical disruption of the
675 spleen, and incubated with ACK lysing buffer (Lonza, LZ10-548E) for 2
676 minutes at room temperature to lyse erythrocytes. The cells were washed in
677 RPMI-1640 (Gibco) medium supplemented with 10% foetal bovine serum
678 (FBS) and B cells were isolated using a B Cell Isolation Kit, mouse (Mil-
679 tenyi Biotec, 130-090-862) according to the manufacturer's instructions.
680 Purified murine B cells were either resuspended in PBS for dSTORM la-
681 belling or frozen in FBS supplemented with 10% DMSO to later culture for
682 live cell imaging (SPT and FCS).

683 BCR complexes were labelled using a recombinant protein based on the Fab
684 fragment of the anti-murine CD79b antibody HM79-16. A self-labelling
685 HaloTag domain was introduced to the C-terminus of the Fab heavy chain
686 to ensure single-dye labelling of the probe. Fab-Halo protein was labelled
687 with HaloTag ligand dyes by incubation with 2-fold molar excess of dye for
688 90 minutes at room temperature, with free dye removed using a Bio-Spin
689 P-6 gel column (BioRad, 7326227) according to manufacturer's instruc-
690 tions. Labelled protein was aliquoted and stored at -80°C . For dSTORM
691 imaging, freshly isolated C57BL/6J B cells were labelled at 4°C with re-
692 recombinant Alexa Fluor 647 Fab-Halo protein. 2×10^6 cells were washed
693 in $0.22 \mu\text{m}^{-2}$ filtered PBS and incubated in $2.5 \mu\text{M}$ Fab-Halo (AF647) for
694 45 minutes at 4°C . Cells were washed twice in cold filtered PBS, fixed in
695 1% paraformaldehyde (Sigma, 28906) for 30 minutes at 4°C , and placed
696 in filtered PBS at a final density of 4×10^7 cells/mL.

697 For live cell imaging, as conducted for SPT and FCS, cells were thawed
698 from frozen stocks and cultured in primary B cell medium (RPMI-1640
699 supplemented with 10% FBS, 2mM L-Glutamine, 10 mM HEPES, 1 mM
700 sodium pyruvate, 50 μM 2-mercaptoethanol, 50 U/mL penicillin and 50
701 $\mu\text{g}/\text{mL}$ streptomycin), supplemented with 10 $\mu\text{g}/\text{mL}$ anti-mouse CD40
702 (clone 1C10, Biolegend 102812) and 10 ng/mL murine IL-4 (Peprotech,
703 214-14). For live cell imaging, 2×10^5 cells were washed in filtered PBS
704 and incubated with 1 μM fluorescent Fab-Halo for 15 minutes at room tem-
705 perature, and washed twice in PBS prior to incubation with the coverslip.

706 **HeLa cell culture and fluorescent labelling.** HeLa TDS cells
707 were cultured at 37°C and 5% CO_2 in DMEM (Gibco, Invitrogen) sup-
708 plemented with 10% FBS (Life Technologies), 1% penicillin/streptomycin
709 (Life Technologies), and 1% glutamine (Life Technologies). Cells were
710 passaged every three days and were regularly tested for mycoplasma. One
711 day prior to fixation, cells were seeded on high-precision 1.5 glass cover-
712 slips (MatTek, P35G-0.170-14-C) for imaging.

713 Cells were fixed and permeabilised simultaneously for 6 minutes in Cy-
714 toskeleton Buffer with Sucrose (CBS, 10 mM MES, 138 mM KCl, 3
715 mM MgCl_2 , 2 mM EGTA, and 4.5% sucrose w/v, pH 7.4) containing 4%
716 methanol-free formaldehyde (FA, Fisher Scientific) and 0.2% Triton, fol-
717 lowed by a second fixation for 14 minutes in CBS + 4% methanol-free
718 formaldehyde at 37°C and 5% CO_2 . Post fixation, cells were washed three
719 times in PBS + 0.1% Tween (PBST), and further permeabilised in PBS +
720 0.5% Triton for 5 minutes. Cells were then washed in PBST three times
721 and blocked in 5% BSA (in PBS) for 1 hour at RT. Samples were further
722 washed three times in PBST, after which samples were incubated with an
723 anti- α -tubulin antibody (ab7291, clone DM1A, at $2.5 \mu\text{g mL}^{-1}$ in 5%
724 non-fat milk) overnight at 4°C . Cells were then washed six times in PBST
725 after which a Donkey anti-Mouse IgG (H+L) Highly Cross-Adsorbed Sec-
726 ondary Antibody AlexaFluor 647 (Invitrogen, A-31571, at $2.0 \mu\text{g mL}^{-1}$
727 in 5% non-fat milk) was added to the sample for 1 hour at 4°C . The cells
728 were then washed six times in PBS and the sample flooded with STORM
729 imaging buffer prepared as described previously.

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735 Halo-Tag used for 3D-SPT.

736 **AUTHOR CONTRIBUTIONS**

737 SD, AH and SFL conceived the project. SFL supervised the research. SD and
738 AH built the optical set-up and performed baseline experiments. SD performed all
739 SMLFM imaging and data analysis. JFF prepared all labelled B-cell samples and
740 conducted FCS experiments. EB programmed and performed the simulations. RP
741 and BZ prepared labelled tubulin samples. ES provided samples for early tests.
742 RRS, KOH, SB and BZ wrote and maintained the 3D reconstruction code. JSB
743 provided diffusion analysis code. SD and SFL wrote the manuscript with input
744 from all authors.

745 **COMPETING FINANCIAL INTERESTS**

746 CAIRN research has a co-development agreement with SFL and KOH at the Uni-
747 versity of Cambridge.